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Coupled equations for Kähler metrics and Yang-Mills connections

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Introduction

This thesis is concerned with coupled equations for a Kähler metric on a compact complex manifold and a connection on a holomorphic principal bundle over it. In the framework of Kähler geometry it is natural to consider equations with symplectic interpretation, appearing as the vanishing of a moment map. Many partial differential equations in geometry and mathematical physics appear in this way. Fundamental examples include the Hermite–Yang–Mills equation for connections and the constant scalar curvature equation for Kähler metrics. The moment map interpretation of these two equations provides a conceptual framework and a guideline to understand many phenomena in Yang–Mills theory and Kähler geometry. Our equations have a moment map interpretation, that intertwines the scalar curvature of a Kähler metric with the curvature of a Hermite–Yang–Mills connection. The coupling occurs as a direct consequence of the structure of the group of symmetries. We expect that the symplectic reduction process will provide moduli spaces of solutions endowed with a Kähler structure.

A fundamental result in complex geometry is the Hitchin–Kobayashi correspondence [21, 79, 62]. It relates the existence of solutions to the Hermite–Yang–Mills (HYM) equation for a connection on a holomorphic principal bundle over a compact Kähler manifold with the Mumford stability of the bundle. As pointed out first by Atiyah and Bott for the case of Riemann surfaces [6] and generalised by S.K. Donaldson to arbitrary compact Kähler manifolds, solutions to the HYM equation arise as points of a symplectic reduction for the action of the gauge group \mathcal{G} on the space of connections. This moment map interpretation, leads one to think of the Hitchin–Kobayashi correspondence as an infinite dimensional version of the Kempf–Ness Theorem relating symplectic reduction with GIT quotients. A similar result for constant scalar curvature Kähler (cscK) metrics on a compact complex manifold has been conjectured by S.T. Yau, G. Tian and S.K. Donaldson. As observed by Fujiki [28] and Donaldson [23], the cscK equation has a moment map interpretation where the group of symmetries is the group \mathcal{H} of Hamiltonian symplectomorphisms of a fixed symplectic structure. This gives a fundamental evidence of the *Yau–Tian–Donaldson conjecture*, that relates the existence of solutions of the cscK equation with the stability of the complex structure on the manifold. The notion of K-stability, introduced by G. Tian in [74] and generalised by S.K. Donaldson in [24], plays a key role in this context. The search of such a correspondence for the cscK equation is, today, one of the central problems in complex and algebraic geometry. In §1.3 we give a brief survey of the state of the art of the cscK problem, centering attention on those aspects relevant to the contents of this thesis. General overviews of the Kempf–Ness Theorem and the Hitchin–Kobayashi correspondence for the HYM equation are provided in §1.1 and §1.2.

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The moment map interpretations of the cscK equation and the HYM equation provide a guide that we follow, in order to define our system of equations for Kähler metrics and connections. The moment map interpretation for the cscK equation arises when we fix a symplectic form ω on a smooth compact Kählerian manifold X , i.e. a manifold that admits a Kähler structure. Let G be a real compact Lie group and E a smooth principal G -bundle over X . The phase space for our problem is provided by a suitable subspace of the product

$$(0.0.1) \quad \mathcal{P} \subset \mathcal{J} \times \mathcal{A},$$

where \mathcal{J} is the space of ω -compatible complex structures on the base X and \mathcal{A} is the space of connections on the bundle E . Let G^c be the complexification of the group G . The space \mathcal{P} parameterizes pairs consisting of a Kähler structure on the base X , with Kähler form ω , together with a structure of holomorphic principal bundle on the associated principal G^c -bundle

$$E^c = E \times_G G^c$$

over the corresponding complex manifold. At this point there are mainly two issues that will lead us to a ‘good’ system of equations: the identification of the group of symmetries and the choice of a symplectic form on the space of parameters. For the first issue, let us consider the group $\text{Aut}(E)$ of G -equivariant diffeomorphisms of $\text{Tot}(E)$, the total space of the bundle E . Note that any element of $\text{Aut}(E)$ covers a unique diffeomorphism on the base. The group of symmetries is taken to be the subgroup $\tilde{\mathcal{G}} \subset \text{Aut}(E)$ covering \mathcal{H} , the group Hamiltonian symplectomorphisms of (X, ω) . The *extended gauge group* $\tilde{\mathcal{G}}$ is an extension of \mathcal{H} by the gauge group \mathcal{G} of the bundle E (see §2.2), i.e. there is an extension of groups

$$1 \rightarrow \mathcal{G} \rightarrow \tilde{\mathcal{G}} \rightarrow \mathcal{H} \rightarrow 1.$$

Recall that the groups \mathcal{G} and \mathcal{H} are, respectively, the symmetry groups for the HYM problem and the cscK problem. For the choice of symplectic form we rely on the symplectic forms used in these problems. Each factor in the product (0.0.1) has a natural symplectic structure provided, on \mathcal{J} and on \mathcal{A} , by the Fujiki–Donaldson symplectic structure $\omega_{\mathcal{J}}$ [28, 22] and by the Atiyah–Bott symplectic structure $\omega_{\mathcal{A}}$ [6]. Our choice of symplectic structure on \mathcal{P} is the simplest possible: we take a pair of non-zero coupling constants $\alpha = (\alpha_0, \alpha_1) \in \mathbb{R}^2$ and consider the 2-form

$$(0.0.2) \quad \omega_{\alpha} = \alpha_0 \cdot \omega_{\mathcal{J}} + \frac{4\alpha_1}{(n-1)!} \cdot \omega_{\mathcal{A}},$$

where $2n$ is the real dimension of the manifold X . The previous normalization of α_1 is convenient to simplify the shape of the equations. When $\frac{\alpha_1}{\alpha_0} > 0$, the symplectic form ω_{α} is Kähler with respect to the structure of complex manifold on \mathcal{P} (see (2.3.4) and Proposition 2.3.2). Thus, a moduli construction using the symplectic reduction process on $(\mathcal{P}, \omega_{\alpha})$ will lead to a Kähler moduli space. The remaining ingredient is the fact that the extended gauge group $\tilde{\mathcal{G}}$ acts on the space \mathcal{P} by holomorphic isometries and there exists a $\tilde{\mathcal{G}}$ -equivariant moment map for any value of the coupling constants (see Proposition 2.3.1)

$$\mu_{\alpha}: \mathcal{P} \rightarrow \text{Lie } \tilde{\mathcal{G}}^*,$$

where $\text{Lie } \tilde{\mathcal{G}}$ is the Lie algebra of $\tilde{\mathcal{G}}$. Our coupled system of equations arise as the zero locus of this moment map. A priori it may seem strange that a single moment map gives rise to a system of equations. The reason for this fact is that any pair $(J, A) \in \mathcal{P}$ determines a vector space splitting $\text{Lie } \tilde{\mathcal{G}}^* \cong \text{Lie } \mathcal{G}^* \oplus \text{Lie } \mathcal{H}^*$ provided by the connection A , and we use this to split in two terms the condition

$$\mu_\alpha(J, A) = 0.$$

The details of the previous construction are provided in Chapter 2.

Let us now introduce the coupled equations that arise when we split the condition above. The unknown variables are a Kähler structure on the base X and a connection on the bundle. Let (g, ω, J) be a Kähler structure on the smooth compact manifold X , where g , ω and J are respectively the metric, the symplectic form and the complex structure. We denote by \mathfrak{g} the Lie algebra of the real compact group G . Recall that E is a smooth principal G -bundle over X . We fix a positive definite inner product (\cdot, \cdot) on \mathfrak{g} invariant under the adjoint representation. Considering the space $\Omega^k(\text{ad } E)$ of smooth k -forms on X with values in the adjoint bundle $\text{ad } E$, such an inner product induces a pairing

$$\Omega^p(\text{ad } E) \times \Omega^q(\text{ad } E) \rightarrow \Omega^{p+q},$$

that we write as $(a^p \wedge a^q)$ for any $a^j \in \Omega^j(\text{ad } E)$, $j = p, q$. We will say that a Kähler structure (g, ω, J) on X and a connection A on E satisfy the *coupled equations* with coupling constants $\alpha_0, \alpha_1 \in \mathbb{R}$ if

$$(0.0.3) \quad \left. \begin{aligned} \Lambda F_A &= z \\ \alpha_0 S_g + \alpha_1 \Lambda^2(F_A \wedge F_A) &= c \end{aligned} \right\}.$$

Here S_g is the scalar curvature of the Kähler metric, F_A is the curvature of the connection, z is an element in the center \mathfrak{z} of the Lie algebra \mathfrak{g} and $c \in \mathbb{R}$ is a topological constant. The map $\Lambda: \Omega^{p,q}(\text{ad } E) \rightarrow \Omega^{p-1,q-1}(\text{ad } E)$ is the contraction operator acting on (p, q) -type valued forms determined by the Kähler structure. The link with complex geometry is provided by the additional compatibility condition

$$(0.0.4) \quad F_A^{0,2} = 0,$$

between the complex structure J on the base and the connection A . Here $F_A^{0,2}$ denotes the $(0, 2)$ part of the curvature, regarded as an $\text{ad } E$ valued smooth form on X . Let G^c be the complexification of the group G . When (0.0.4) holds, the pair (J, A) endows the associated principal G^c -bundle $E^c = E \times_G G^c$ with a structure of holomorphic principal bundle over the Kähler manifold X . In the language of the moment map interpretation of the coupled equations (0.0.3) already explained, the compatibility condition (0.0.4) is encoded in the definition of the space \mathcal{P} (0.0.1).

Similar coupled equations, for pseudo-Riemannian metrics and connections on smooth bundles, have been studied for a long time by physicists in the context of unification theory. In this context, the Kähler structure and the connection can be thought of, respectively, as a gravitational field and a gauge field whose interaction is governed by the system of equations. The coupling of these two variables, provided by the equations, reflects their mutual interaction. It is natural to ask if the coupled equations have a variational interpretation, i.e. if their solutions arise as critical points or absolute minima of a suitable action, since many equations in physics follow this

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variational principle. A fundamental example is provided by the Einstein–Yang–Mills (EYM) equations [10], for which the action is obtained by adding the Yang–Mills functional, considered as a functional in both the metric and the connection, to the Hilbert–Einstein action. The coupled equations (0.0.3) also have a variational interpretation as the absolute minima of a suitable functional that we define in §2.4. This functional is purely Riemannian, i.e. it can be considered as a functional for a Riemannian metric g on the smooth compact manifold X and for a connection A on the G -bundle E . Given such a pair (g, A) , the functional is defined by

$$(0.0.5) \quad CYM(g, A) = \int_X (\alpha_0 S_g - 2\alpha_1 |F_A|^2)^2 \cdot \text{vol}_g + 2\alpha_1 \cdot \|F_A\|^2,$$

where vol_g is the volume form of the metric g and the coupling constants α_0 and α_1 are as in (0.0.3). The term $\|F_A\|^2$ is the Yang–Mills functional, considered as a functional in both the metric and the connection, and $|F_A|^2 \in C^\infty(X)$ is the corresponding density. In Proposition 2.4.1 we prove that, when we restrict the functional to Kähler metrics g compatible with a fixed symplectic structure, the solutions to (0.0.3) satisfying the compatibility condition (0.0.4) arise as absolute minima (after suitable re-scaling of the coupling constants). In order to understand this functional better, in §2.4 we think of the data (g, A) as defining a G -invariant metric \hat{g} on $\text{Tot}(E)$. This interpretation shows that the first summand in the definition of the functional (0.0.3) can be thought of, up to scaling, as the L^2 norm of the scalar curvature of a suitable G -invariant metric \hat{g} on E associated to (g, A) . To construct the metric \hat{g} we have to assume the positivity of the ratio of the coupling constants

$$(0.0.6) \quad \frac{\alpha_1}{\alpha_0} > 0.$$

This ratio measures, roughly speaking, the volume of the fibers of the bundle equipped with \hat{g} with respect to the projection $E \rightarrow X$. An unexpected fact here is that (0.0.6) is precisely the condition for the symplectic form ω_α (0.0.2) on \mathcal{P} (0.0.1) to be Kähler. Thus, one can expect that a moduli construction for such metrics using the symplectic reduction process leads to a Kähler moduli space.

This positivity of the ratio (0.0.6) is also a necessary assumption for the main result in §3 (Theorem 3.2.4). In this chapter, we fix a holomorphic principal G^c bundle (E^c, I) over a compact complex manifold (X, J) , where I denotes the holomorphic structure on E^c . We consider the space

$$\tilde{\mathcal{K}} = \mathcal{K} \times \mathcal{R},$$

where \mathcal{K} and \mathcal{R} denote, respectively, the space of Kähler metrics in a fixed Kähler class and the space of smooth G -reductions \mathcal{R} on E^c . The *coupled equations* are then equations for a pair $(\omega, H) \in \tilde{\mathcal{K}}$, where the role of the connection in (0.0.3) is played by the associated Chern connection. In §3.1 we discuss the general framework provided by the Kempf–Ness theorem where our problem fits and establish the link with the moment map interpretation in §2.3. The formalism introduced provides the preliminary material necessary for the rest of this chapter in a way that is completely independent of the choice of the symplectic form in (0.0.1). Using the previous formalism, in §3.2 we prove an existence criterion for the existence of solutions to the coupled system (0.0.3) for small values of the ratio of the coupling constants $\frac{\alpha_1}{\alpha_0} > 0$. This criterion is

formulated in terms of a real subgroup $\tilde{\mathcal{G}}_I$ of the automorphism group $\text{Aut}(E^c, I)$ of the holomorphic bundle. This lead us to the following general result (see Theorem 3.2.4) that we apply in §5.1 to provide non trivial examples to the existence of solutions.

Theorem 0.0.1. *Let (X, L) be a compact polarised manifold, G^c be a complex reductive Lie group and E^c be a holomorphic G^c -bundle over X . If there exists a cscK metric $\omega \in c_1(L)$, X has finite automorphism group and E^c is stable with respect to L then, given a pair of positive real constants $\alpha_0, \alpha_1 > 0$ with small ratio $0 < \frac{\alpha_1}{\alpha_0} \ll 1$, there exists a solution $(\omega_\alpha, H_\alpha)$ to (0.0.3) with these coupling constants and $\omega_\alpha \in c_1(L)$.*

In the following sections we consider the coupled equations (0.0.3) with fixed positive coupling constants $\alpha_0, \alpha_1 > 0$. In §3.3, we give an obstruction coming from a character

$$\mathcal{F}_\alpha: \text{Lie Aut}(E^c, I) \rightarrow \mathbb{C}$$

that generalizes the usual Futaki invariant for the cscK problem [29] and also other invariants defined by A. Futaki in [30]. This character can be considered, roughly speaking, as a \mathbb{C} -linear extension of the moment μ_α whose zero locus correspond to the solutions to the coupled equations. In §3.4 we construct the integral of the moment map

$$\mathcal{M}_\kappa: \tilde{\mathcal{K}} \rightarrow \mathbb{R}$$

and determine a sufficient condition for its convexity and existence of lower bounds in terms of the existence of suitable smooth curves on $\tilde{\mathcal{K}}$ joining any two points. These curves play a role analogue of the one played by the geodesics in \mathcal{K} for the cscK problem. In §3.5 we restrict to the case in which the base is a compact complex surface. Given a pair $(\omega, H) \in \tilde{\mathcal{K}}$ we consider the functional $\mathcal{M}_H = \mathcal{M}_\kappa(\cdot, H): \mathcal{K} \rightarrow \mathbb{R}$ and prove convexity properties of \mathcal{M}_H along smooth geodesics on \mathcal{K} . We use this to prove a uniqueness result for the coupled equations (0.0.3) in the case of toric complex surfaces (see Proposition 3.5.2). We obtain lower bounds of \mathcal{M}_H supposing the existence of solutions of (0.0.3) (see Proposition 3.5.5), when $c_1(X) \leq 0$. For the proof we adapt the argument of X. X. Chen in [17] concerning the Mabuchi K-energy and ϵ -approximate geodesics. We extend the previous result for Kähler surfaces with nonnegative bisectional curvature, making use of the weak solutions of the geodesic equation in \mathcal{K} found by X. X. Chen and G. Tian in [19] and adapting their argument concerning the Mabuchi K-energy (see Proposition 3.5.5).

In §4 we define an algebraic notion of stability for triples (X, L, E) (see Definition 4.1.2) consisting of a smooth projective variety X , an ample line bundle L and a holomorphic vector bundle E over X of rank r and zero degree. Our stability notion depends on a positive real parameter α and relies on Donaldson's definition of K -stability for polarized varieties (see §1.3.6). As in [24], it is formulated in terms of flat degenerations of the triple over \mathbb{C} , called *test configurations*. The central fiber of each test configuration is endowed with a \mathbb{C}^* -action and the stability of the triple, that we call α -K-stability, is measured in terms of the non negativity of a suitable weight F_α (see (4.1.2)) associated to this action. In §4.5 we make the following conjecture (see 4.5.1), where the constant $\alpha = \frac{r\pi^2\alpha_1}{\alpha_0}$ is a normalized ratio of the coupling constants in (0.0.3).

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Conjecture 0.0.2. *If there exists a Kähler metric $\omega \in c_1(L)$ and a Hermitian metric H on E satisfying the coupled equations (0.0.3) with coupling constants α_0, α_1 then the triple (X, L, E) is α -K-polystable.*

The rest of this chapter is devoted to give evidence that the above conjecture holds. In §4.2, Proposition 4.2.1, we prove that the weight F_α satisfies

$$F_\alpha = \frac{-1}{4 \operatorname{vol}(X)\alpha_0} \mathcal{F}_\alpha(\zeta)$$

for degenerations with smooth central fiber, where \mathcal{F}_α is the character defined in §3.3 and ζ is the generator of the $S^1 \subset \mathbb{C}^*$ action on the central fiber. Thus, up to multiplicative constant factor, F_α equals the value of the moment map μ_α (2.3.3) at ζ . Recall that the zero locus of μ_α corresponds to the solutions of the coupled equations. Once more this fits in the general picture of the Kempf–Ness Theorem explained in §1.1 if we view a test configuration as a 1-parameter subgroup (see §4.3) and the associated α -invariant as the corresponding Mumford weight (see §1.1.3 and cf. (1.1.9)). This fact motivates the definition of α -K-stability (see Definition 4.1.2) and leads to think that this stability notion should play an important role in the existence problem for the coupled equations (3.0.7). Moreover, this fact gives the following more direct evidence.

1) *If there exists a solution to the coupled equations (0.0.3) with positive coupling constants $\alpha_0, \alpha_1 \in \mathbb{R}$ then $F_\alpha = 0$ for any product test configuration (see Corollary 4.2.2).*

In §4.3 we interpret test configurations for a triple as 1-parameter subgroups of suitable product group G_k acting on a relative Quot scheme parameterizing triples (X, L, E) . This gives the following evidence.

2) *In the limit case $\alpha = 0$ the notion of α -K-stability is equivalent to the notion of algebraic K-stability of the underlying polarized manifold (see Proposition 4.3.3).*

In §4.4 we characterize the notion of Mumford–Takemoto stability of the bundle in terms of *base preserving test configurations* (see Proposition 4.4.1). This leads to the following evidence.

3) *If there exists a solution to the coupled equations (0.0.3) with positive coupling constants $\alpha_0, \alpha_1 \in \mathbb{R}$, then $F_\alpha \geq 0$ for any base preserving test configuration, with equality only if the test configuration is a product (see Corollary 4.4.3).*

We hasten to add that Conjecture 0.0.2 is difficult to prove even in complex dimension 1. Note that in this case the coupled equation are a system in separated variables independent of the coupling constants. In fact, they reduce to the HYM equation for the bundle due to the uniformisation theorem, that asserts that any Riemann surface has a unique cscK metric if we prescribe the total area. When $\dim_{\mathbb{C}} X = 1$, Conjecture 0.0.2 reduces to prove that Mumford–Takemoto stability of the bundle implies α -K-stability of the triple for any $\alpha > 0$, due to the Hitchin–Kobayashi correspondence. The algebro-geometric approach to the proof of this fact seems to be difficult in the light of previous work by R. Thomas and J. Ross in [64], where the algebraic K-stability of Riemann surfaces is analysed with purely algebro-geometric methods. A more likely approach, valid for arbitrary dimension of X , seems to be provided by Donaldson’s methods in [25] or Mabuchi’s methods in [49]. The main issue to adapt

the arguments in [25] or [49] is to prove that α -K-stability arises as a suitable asymptotic version of a Chow-type stability notion in the relative quot scheme considered in §4.3. The verification of this stability issue would provide a strong evidence of the fact that (0.0.3) are gauge-theoretic equations related to the algebro-geometric moduli problem for triples (X, L, E) . The moduli problem for such triples has been considered by R. Pandharipande in [60] when the base X is a Riemann surface. In [60], Geometric Invariant Theory is used in order to compactify the moduli space of pairs consisting of a (10-canonically polarised) smooth curve of genus greater than one and a semistable vector bundle over the curve. For arbitrary dimension of the base very little has been done. In [66], G. Schumacher and M. Toma construct a moduli space of (non-uniruled) polarized Kähler manifolds equipped with isomorphism classes of stable vector bundles via the method of versal deformations. In this work the authors were able to endow the moduli space with a Kähler metric, under suitable cohomological restrictions (see §5.1), when the base is Kähler–Einstein and the bundle is projectively flat. As we see in §5.1, the gauge-theoretical equations corresponding to the moduli construction in [60] and the Kähler moduli in [66] provide particular solutions of the coupled equations (0.0.3).

Finally, in §4.5 we analyze the converse of the statement in Conjecture 0.0.2 in the cases $\dim_{\mathbb{C}} X = 1$, E arbitrary and $E \cong \oplus^r \mathcal{O}_X$, X arbitrary. We respectively check that it holds and that it reduces to the Yau–Tian–Donaldson Conjecture for the cscK problem. The latter gives evidence of the fact that, as a recent example in [2] suggests for the cscK problem, the notion of α -K-stability introduced in Definition 4.1.2 may not be sufficient to guarantee the existence of solutions of (0.0.3) and so, to formulate a converse of 0.0.2, a stronger notion may be required. We indicate how to strengthen this condition following [63].

Chapter §5 contains some examples of solutions to (0.0.3). We also discuss how the existence of solutions in the limit case $\alpha_0 = 0$ for $G^c = SL(r, \mathbb{C})$ can be applied in certain cases, using results of Y. J. Hong in [36], to obtain a cscK metric in the projectivization of the associated vector bundle. In §5.1 we deal with the case of vector bundles over Riemann surfaces and projectively flat bundles over Kähler manifolds satisfying a topological constraint. In both situations, the coupled system (0.0.3) reduces to the limit case $\alpha_1 = 0$ (cscK equation and HYM equation). In §5.2 we consider homogeneous Hermitian holomorphic bundles E over homogeneous Kähler manifolds and show the existence of invariant solutions to (0.0.3) in any Kähler class provided that the representation inducing the homogeneous bundle E is irreducible. The invariant solutions arise as (simultaneous) solutions to the limit systems $\alpha_0 = 0$ and $\alpha_1 = 0$ in (0.0.3). When the complex dimension of the base is 2 the invariant solutions satisfy in addition the Kähler analogue of the Einstein–Yang–Mills equation (see equation (2.4.6)) considered in §2.4. In §5.3 we discuss some (well known) examples of stable bundles over Kähler–Einstein manifolds without infinitesimal automorphisms where Theorem 0.0.1 applies. Using a bubbling argument for Anti-Self-Dual connections, this setting provides our first genuine examples in which the metric on the base is not cscK. In §5.4 we discuss the relation of the limit $\alpha_0 = 0$ in (0.0.3) with the cscK equation on ruled manifolds as mentioned above.

INTRODUCTION

In §6 we consider the problem of finding solutions to (0.0.3) on the non-compact (standard) complex manifold \mathbb{C}^2 endowed with the trivial smooth principal bundle $E = \mathbb{C}^2 \times SU(2)$. We prove that there is no obstruction to finding solutions for arbitrary coupling constants with fixed behaviour at infinity (see Theorem 6.3.5). The fixed behavior at infinity is the analogue in this non-compact case of fixing the Kähler class and the topological type of the bundle. To prove Theorem 6.3.5, we use a deformation argument and a suitable re-scaling of the solutions. As initial data we take the standard Euclidean metric and a classical solution of 't Hooft to the instanton equation. This procedure gives solutions for which the Kähler metric is not cscK, making explicit the ‘bubbling’ argument in §5.3.

List of symbols

Lie groups, Lie algebras, manifolds and bundles

G^c	Complex reductive Lie group
\mathfrak{g}^c	Lie algebra of G^c
$G \subset G^c$	Maximal compact subgroup of G^c
\mathfrak{g}	Lie algebra of G
$\mathfrak{z} \subset \mathfrak{g}$	centre of \mathfrak{g}
X	Compact differentiable manifold (projective scheme in Chapter §4)
Z	Projective variety
E	Smooth principal G -bundle over X (coherent sheaf over X in Chapter §4)
$\text{ad } E$	Adjoint bundle of E
E^c	Smooth principal G^c -bundle over X
L	Line bundle over X (ample holomorphic in Chapter §4)
$\text{Tot}(E)$	Total space of the bundle E
$\text{Tot}(E^c)$	Total space of the bundle E^c
$\text{Tot}(L)$	Total space of the bundle L
(X, L, E)	Triple consisting of a projective scheme X , an ample invertible sheaf L and a coherent sheaf E over X

Geometric structures

ω	Symplectic structure on X
J	Almost complex structure on X
g	Riemannian metric; also elements of the groups $\text{Aut } E$ or $\text{Aut } E^c$ (below)
ω_ϕ	$\omega + dd^c\phi$, for a Kähler potential ϕ on (X, J, ω)
S_g, S_J	Scalar curvature of g , Hermitian scalar curvature of J
A	Connection on E
θ_A	Vertical projection determined by A
θ_A^\perp	Horizontal lift determined by A
F_A	Curvature of A
I	G^c -invariant almost complex structure on E^c
H	G -reduction on E^c
A_H	Chern connection of H , when E^c is holomorphic
F_H	Curvature of A_H

Constants and parameters

$2n$	Real dimension of X
r	Rank of of a vector bundle E
$\hat{c} \in \mathbb{R}$	Normalized integral of the $2n$ -form $\Lambda^2(F_A \wedge F_A) \wedge \omega^{n-2}$
$\hat{S} \in \mathbb{R}$	Normalized integral of the scalar curvature
$\alpha_0, \alpha_1 \in \mathbb{R}$	Coupling constants
$\alpha \in \mathbb{R}$	Normalized ratio of the coupling constants

$c \in \mathbb{R}$	weighted sum of \hat{c} and \hat{S} with weights α_0 and α_1 , respectively
$z \in \mathfrak{z}$	Element in the center of \mathfrak{g}

Spaces of geometric structures

$C^\infty(X)$	Smooth functions on X
$C_0^\infty(X)$	Smooth functions with zero integral on (X, ω)
$\Omega^k(F)$	Space of F -valued smooth k -forms, F being a vector bundle
\mathcal{J}	Space of compatible almost complex structure on (X, ω)
$\Omega_J^{p,q}$	Space of smooth (p, q) -type forms determined by $J \in \mathcal{J}$
$\mathcal{J}^i \subset \mathcal{J}$	Space of integrable elements of \mathcal{J}
\mathcal{A}	Space of connections on E
$\mathcal{A}_J^{1,1} \subset \mathcal{A}$	Connections of $(1, 1)$ -type with respect to $J \in \mathcal{J}$
\mathcal{P}	Space of pairs $p = (J, A)$, with $J \in \mathcal{J}^i$ and $A \in \mathcal{A}_J^{1,1}$
\mathcal{K}	Space of Kähler metrics on X , on a fixed Kähler class.
\mathcal{R}	Space of G -reductions of E^c
$\tilde{\mathcal{K}}$	$= \mathcal{K} \times \mathcal{R}$

More (infinite dimensional) Lie groups

\mathcal{H}	Group of Hamiltonian symplectomorphisms of (X, ω)
$\text{Diff } X$	Group of diffeomorphisms of X
$\text{Diff}_0 X$	Identity component of $\text{Diff } X$
$\text{Aut } (X, J)$	Group of biholomorphisms of a complex manifold
$\text{Aut } (X, J, \omega)$	Group of holomorphic isometries of a Kähler manifold
\mathcal{G}	Gauge group of E
$\text{Aut } E$	Group of G -equivariant diffeomorphisms of E
$\tilde{\mathcal{G}} \subset \text{Aut } E$	Extended gauge group of E and (X, ω)
\mathcal{G}^c	Gauge group of E^c
$\text{Aut } E^c$	Group of G -equivariant diffeomorphisms of E^c

Vector fields and functions

$\phi \in C^\infty(X)$	Smooth function on X . Also Kähler potential on (X, J, ω)
y	Vector field on X or E
η_ϕ	Hamiltonian vector field of $\phi \in C^\infty(X)$. Also element of Lie \mathcal{H}
η	Element of Lie \mathcal{H}
ζ	Element of Lie $\tilde{\mathcal{G}}$
ξ	Element of Lie \mathcal{G}
Y_ζ	Infinitesimal action of ζ on \mathcal{P}

CHAPTER 1

Preliminaries

1.1. The Kempf-Ness Theorem

We start recalling some well known facts about symplectic and algebraic quotients that apply formally to our problem. We refer the reader to [55] for the material concerning symplectic quotients and to [54] for that concerning Geometric Invariant Theory (GIT). For the proof of the Kempf-Ness Theorem we have mainly used the theory developed in [57] but also the classical references [39] and [54]. For further results on this topic for the case of non-projective Kähler manifolds we refer to [58].

Let V be a complex Hermitian vector space and let $Z \subseteq \mathbb{P}(V)$ be a compact smooth projective variety. We consider Z endowed with the Kähler metric ω given by restricting the Fubini-Studi metric of $\mathbb{P}(V)$. Suppose that a reductive complex algebraic group G^c acts on Z via an algebraic representation $\rho : G^c \rightarrow SL(V)$, thus the G^c -action is naturally linearized with respect to the line bundle $L = \mathcal{O}_Z(1)$. Let $G \subset G^c$ be a maximal compact subgroup that maps onto $SU(V)$. Then G acts on Z by holomorphic isometries and the linearization provides a G -equivariant moment map $\mu : Z \rightarrow \mathfrak{g}^*$ (see e.g. [27, §6.5.1]). In this situation there are two quotients, provided by symplectic geometry and algebraic geometry, and a well known correspondence between them that will be explained along this section.

On the other hand, the symplectic quotient of (Z, ω) by G is constructed via the moment map μ . Recall that a map $\mu : Z \rightarrow \mathfrak{g}^*$ is called a G -equivariant moment map if it satisfies

$$\langle d\mu, \zeta \rangle = i_{Y_\zeta} \omega, \quad \mu(gz) = \text{Ad}(g^{-1})^* \mu(z),$$

for all $z \in Z$, $g \in G$ and $\zeta \in \mathfrak{g}$. Here the vector field Y_ζ denotes the infinitesimal action of ζ on Z . If $0 \in \mathfrak{g}^*$ is a regular value of μ then $\mu^{-1}(0)$ is a smooth submanifold of Z preserved by G and the quotient $\mu^{-1}(0)/G$ is a compact manifold (orbifold if the G -action on $\mu^{-1}(0)$ is not free) that inherits a natural symplectic structure, called the symplectic quotient of (Z, ω) by G .

On the other hand, the (algebraic) GIT quotient of Z by G^c with respect to the linearization L is constructed as follows. Let

$$R = \bigoplus_{k \geq 0} H^0(X, L^k)$$

be the graded ring given by tensorization, where $H^0(X, L^k)$ denotes the space of global holomorphic sections of L^k . The GIT quotient of Z by G^c is the projective variety

$$Z//G^c := \text{Proj } R^{G^c},$$

where R^{G^c} is the (graded) ring of G^c -invariant elements of R . In more down-to-earth terms, there is an isomorphism of graded rings $R \cong \mathbb{C}[\hat{Z}]$, where $\hat{Z} \subset V$ is the affine cone of Z and $\mathbb{C}[\hat{Z}]$ denotes its ring of regular functions. Since G^c is a reductive complex algebraic group we can take a finite set of generators $\{f_j\}_{j=0}^m$ of the ring R^{G^c} . Let $\mathbb{C}[z_0, \dots, z_m]$ be the ring of polynomials in $m+1$ free variables and consider the (graded) ring morphism defined by

$$(1.1.1) \quad \mathbb{C}[z_0, \dots, z_m] \rightarrow R^{G^c} : z_j \mapsto f_j.$$

The kernel of (1.1.1) is an homogeneous ideal of $\mathbb{C}[z_0, \dots, z_m]$ cutting a projective variety on \mathbb{P}^m that turns out to be isomorphic to $Z//G^c$.

The link between the symplectic and GIT quotients is provided by the Kempf-Ness Theorem that states that there exists an homeomorphism between them (see [38, 39] or [57, Theorem 5.4]).

Theorem 1.1.1 (Kempf-Ness).

$$(1.1.2) \quad \mu^{-1}(0)/G \cong Z//G^c.$$

While the points in the symplectic quotient have a clear geometric meaning, they represent G -orbits on $\mu^{-1}(0)$, it is not clear a priori which G^c -orbits on Z are represented in the GIT quotient. In order to make the homeomorphism explicit we shall clarify this point, for what we need to introduce the notion of stability.

1.1.1. GIT Stability. Note that the inclusion $R^{G^c} \hookrightarrow R$ induces a G^c -invariant rational map

$$(1.1.3) \quad Z \dashrightarrow Z//G^c.$$

In fact it is not always possible to make this map every where well-defined; if all the G^c -invariant functions of \hat{Z} vanish at some point $z \in Z$ then clearly (1.1.3) is not well defined at z . This leads to the fundamental definition of GIT.

Definition 1.1.2. A point $z \in Z$ is said to be *semistable* if there exists a nonconstant function $f \in R^{G^c}$ such that $f(z) \neq 0$. We say z is *polystable* if it is semistable and the orbit $G^c \cdot \hat{z} \subset V$ of any lift $\hat{z} \in \hat{Z}$ of z is closed in V . We say z is *stable* if it is polystable and z has finite stabilizer on G^c .

The locus of semistable, polystable and stable points will be denoted respectively by Z^{ss} , Z^{ps} and Z^s . It is clear that the notion of stability do not depend on the representant of any G^c orbit, so it makes sense to talk about stability of an orbit of the action. The semistable orbits are precisely those that appear in the GIT quotient, i.e. Z^{ss} is the domain of definition of the rational map (1.1.3). Any semistable orbit is glued in the quotient to the unique polystable orbit on its closure and, only when restricted to Z^s , the map (1.1.3) parameterizes G^c -orbits. This justifies the equality

$$Z//G^c = Z^{ss}//G^c.$$

A point will be referred to as *unstable* if it is not semistable. Unstable points are characterized by the following property: $z \in Z$ is unstable if and only if the orbit of any lift $\hat{z} \in \hat{Z}$ contains $0 \in V$ on its closure.

The main result that makes workable Definition 1.1.2 is a numerical criterion, due to David Hilbert and David Mumford, that provides a systematic procedure to check which are the semistable points. It is stated in terms of the restriction of the action to *1-parameter subgroups*. Let $\lambda : \mathbb{C}^* \rightarrow G^c$ be a 1-parameter subgroup, i.e. a morphism of algebraic groups from \mathbb{C}^* to G^c . Since Z is projective for any $z \in Z$ there exists a limit point

$$z_0 = \lim_{t \rightarrow 0} \lambda(t) \cdot z,$$

that is fixed by the \mathbb{C}^* -action. The *maximal weight* $w(z, \lambda)$ associated to λ at z is the weight of the induced \mathbb{C}^* -action on the fiber L_{z_0} . Note that the morphism λ induces a representation of \mathbb{C}^* on V that, in suitable coordinates $\hat{z} = (z_0, \dots, z_n)$, reads

$$\lambda(t) \cdot \hat{z} = (t^{q_0} z_0, \dots, t^{q_n} z_n),$$

and so the maximal weight equals

$$w(z, \lambda) = \max\{-q_i : x_i \neq 0\}.$$

If $w(\lambda, z) < 0$ then, taking $t \rightarrow 0$, the origin is contained in the closure of the orbit of \hat{z} in V so the point z is unstable. The Hilbert-Mumford theorem says that the converse is true, i.e. that if z is unstable then $w(\lambda, z)$ is positive for some 1-parameter subgroup λ . Moreover,

Theorem 1.1.3 (Hilbert-Mumford Criterion). *A point $z \in Z$ is semistable if and only if $w(z, \lambda) \geq 0$ for all 1-parameter subgroups. A point $z \in Z$ is polystable if and only if $w(z, \lambda) \geq 0$ for all 1-parameter subgroups with equality only if λ fixes z . A point $z \in Z$ is stable if and only if $w(z, \lambda) > 0$ for all non-trivial 1-parameter subgroups.*

1.1.2. The character of a complex orbit. As we have seen, the GIT quotient does not represent all the G^c -orbits in the Kähler manifold (Z, J, ω) so, in order to prove Theorem 1.1.1, one has to identify which G^c -orbits cut the zero locus of the moment map. Before going into the proof of the theorem we introduce a well known invariant of the G^c -orbit of a point $z \in Z$ that gives a necessary condition for this fact. Our approach differs slightly from the ones in [31] and [81], since we make an explicit use of an exact 1-form on G^c defined in [57]. Let $\sigma_z \in \Omega^1(G^c)$ be given by

$$(1.1.4) \quad \sigma_z(v_g) = \langle \mu(gz), -\mathbf{i}\pi_1(R_{g^{-1}})_*v_g \rangle$$

for any $v_g \in T_g G^c$, where $R_{g^{-1}}$ denotes right translation by g^{-1} and $\pi_1 : \mathfrak{g}^c \rightarrow \mathbf{i}\mathfrak{g}$ is the standard projection. This form is exact, invariant by the left G action on G^c and satisfies

$$(1.1.5) \quad R_g^* \sigma_z = \sigma_{gz}$$

for any $g \in G^c$ (see [57, Lemma 3.1]). Since G^c retracts to G and the restriction of σ_z to G is zero, this 1-form is exact if and only if it is closed. We sketch a proof of the latter for completeness.

Proposition 1.1.4 (Mundet i Riera [56]). *The 1-form σ_z is closed.*

PROOF. By property (1.1.5) it is enough to check it at the tangent space of the identity element on G^c . Given $\zeta_j \in \mathfrak{g}^c$, $j = 1, 2$, the symbol ζ_j^r means the right invariant vector field on G^c generated by ζ_j . We have to check that

$$(1.1.6) \quad d\sigma_z(\zeta_1^r, \zeta_2^r) = \zeta_1^r(\sigma_z(\zeta_2^r)) - \zeta_2^r(\sigma_z(\zeta_1^r)) - \sigma_z([\zeta_1^r, \zeta_2^r]) = 0.$$

Since σ_z vanishes on G we can skip the case $\zeta_1, \zeta_2 \in \mathfrak{g}$. By G -equivariance of the moment map it follows that (1.1.6) holds if $\zeta_1 \in \mathfrak{g}$ and $\zeta_2 \in i\mathfrak{g}$. Finally, for the remaining case we take $i\zeta_j \in i\mathfrak{g}$ and so

$$d\sigma_z(i\zeta_1^r, i\zeta_2^r) = \omega_z(Y_{\zeta_2}, JY_{\zeta_1}) - \omega_z(Y_{\zeta_1}, JY_{\zeta_2}) = 0. \quad \square$$

Let $(G^c)_z$ be the isotropy group of z in G^c , with Lie algebra $(\mathfrak{g}^c)_z$ and consider the \mathbb{C} -linear extension of the moment map $\mu : Z \rightarrow (\mathfrak{g}^c)^*$. We define a complex valued map $f_{z,g} : (\mathfrak{g}^c)_z \rightarrow \mathbb{C}$ by

$$(1.1.7) \quad \begin{aligned} f_{z,g}(\zeta) &= \langle \mu(gz), \text{Ad}(g)\zeta \rangle \\ &= i\sigma_z(\zeta^l)|_g + \sigma_z((i\zeta)^l)|_g, \end{aligned}$$

where $\pi_0 : \mathfrak{g}^c \rightarrow \mathfrak{g}$ is the standard projection and ζ^l is the left invariant vector field on G^c determined by any $\zeta \in \mathfrak{g}^c$. Then we obtain the desired obstruction.

Proposition 1.1.5 ([81]). *The map $f_z : (\mathfrak{g}^c)_z \rightarrow \mathbb{C}$ given by (1.1.7) is independent of $g \in G^c$. It defines a character of $(\mathfrak{g}^c)_z$ that vanishes if $G^c \cdot z$ contains a zero of the moment map.*

PROOF. For the first part of the statement it is enough to check that $\sigma_z(\zeta^l)$ is constant on G^c for any $\zeta \in (\mathfrak{g}^c)_z$ and using (1.1.5) we just have to check it at 1. Thus, given any $\eta \in \mathfrak{g}^c$,

$$\frac{d}{dt}|_{t=0} \sigma_z(\zeta^l)|_{e^{t\eta}} = d(\sigma_z(\zeta^l)) \cdot \eta = (-i\zeta^l d\sigma_z + L_{\zeta^l} \sigma_z) \cdot \eta = L_{\zeta^l} \sigma_z \cdot \eta = 0,$$

where the last equality follows from (1.1.5). Combining $d\sigma_z(\zeta^l) = 0$ with (1.1.6), we obtain

$$f_{z,g}([\zeta_1, \zeta_2]) = i\sigma_z([\zeta_1^l, \zeta_2^l]) + \sigma_z([(i\zeta_1)^l, \zeta_2^l]) = 0,$$

for any pair $\zeta_1, \zeta_2 \in (\mathfrak{g}^c)_z$ as claimed. \square

Remark 1.1.6. The complex group G_z^c , the complexification of G_z , is a complex subgroup of $(G^c)_z$ but they may not coincide for every $z \in Z$. However, if $z \in Z$ contains a zero of the moment map in the orbit $G^c \cdot z$ then $G_z^c = (G^c)_z$ (see [31, Proposition 2.4]).

1.1.3. The proof of Kempf-Ness Theorem. The key idea to prove the existence of the homeomorphism (1.1.2), due to G. Kempf and L. Ness in [38], is to construct a convex functional on each G^c orbit on V whose critical points coincide with the lifts of zeros of the moment map μ . Using the Hilbert-Mumford criterion it follows that the G^c orbits of polystable points are precisely those that cut the zero locus of μ . Here we adapt the proof of I. Mundet i Riera in [57] to the polystable case, since it provides a general argument for non projective Kähler manifolds that will be useful later in this thesis.

First we shall relate the tools used to construct the symplectic and GIT quotients, i.e. the moment map and the linearization. As pointed out before, the representation $\rho : G^c \rightarrow SL(V)$ gives a lift of the G^c -action to the line bundle L that determines a moment map

$$\langle \mu(z), \zeta \rangle = \frac{\mathbf{i}}{2\pi} \frac{h(\hat{z}, \rho_*(\zeta)\hat{z})}{|\hat{z}|^2}$$

for the G action ([27, §6.5.1]), where $\zeta \in \mathfrak{g}$, $\hat{z} \in V$ is any vector lying over $z \in Z$ and the squared norm $|\hat{z}|^2$ is taken with respect to the Hermitian metric h on V . The moment map and the linearization are matched together via an infinitesimal identity on $\text{Tot}(L)$. Indeed, if $\zeta \in \mathfrak{g}$ and we denote respectively by Y_ζ and \hat{Y}_ζ the infinitesimal actions on Z and L , we have

$$(1.1.8) \quad \hat{Y}_\zeta = \theta_A^\perp(Y_\zeta) + 2\pi \mathbf{i} \langle \mu, \zeta \rangle \text{Id}.$$

Here A is the Chern connection of the Hermitian metric in L (induced by h) and θ_A^\perp denotes the horizontal lift of vector fields on Z with respect to A . By Id we denote the canonical vertical vector field on L induced by the identity endomorphism. The identity (1.1.8) shows that if $\zeta \in \mathfrak{g}$ generates a 1-parameter subgroup λ_ζ the corresponding maximal weight satisfies

$$(1.1.9) \quad w(z_0, \lambda_\zeta) = 2\pi \langle \mu(z_0), \zeta \rangle.$$

The right hand side of this expression can be rewritten using that (cf. [57, Definition 2.3])

$$(1.1.10) \quad \langle \mu(z_0), \zeta \rangle = \langle \mu(\lim_{t \rightarrow 0} \lambda(t)z), \zeta \rangle = \langle \mu(\lim_{t \rightarrow \infty} e^{\mathbf{i}t\zeta} z), \zeta \rangle = \lim_{t \rightarrow \infty} \langle \mu(e^{\mathbf{i}t\zeta} z), \zeta \rangle.$$

We now recall from [57] the construction of a convex functional, very often referred as the *integral of the moment map*. Given $z \in Z$ the integral of the moment map $\mathcal{M}_z : G \setminus G^c \rightarrow \mathbb{R}$ is defined as the unique functional that satisfies the properties

$$(1.1.11) \quad d\mathcal{M}_z = \sigma_z \quad \text{and} \quad \mathcal{M}_z([1]) = 0,$$

where σ_z denotes now the 1-form on $G \setminus G^c$ induced by the G -invariant form (1.1.4). Property (1.1.5) of σ_z is translated into the following cocycle condition

$$(1.1.12) \quad \mathcal{M}_z(gh) = \mathcal{M}_{hz}(g) + \mathcal{M}_z(h).$$

We will now use \mathcal{M}_z to give a prove of the fundamental lemma from which Kempf-Ness Theorem follows. This lemma says that the polystability of the point z is equivalent to the fact that \mathcal{M}_z attains a minimum, the latter being equivalent to $G^c z \cap \mu^{-1}(0) \neq \emptyset$ due to (1.1.11). To give a clue of why this is true recall that $G \setminus G^c$ is a symmetric space of noncompact type that can be endowed with a preferred negatively curved Riemannian metric. The geodesics corresponding to this metric are given by maps $t \in \mathbb{R} \rightarrow [e^{\mathbf{i}t\zeta} g] \in G \setminus G^c$ for any $\zeta \in \mathfrak{g}$ and $g \in G^c$. Along a geodesic, the integral of the moment has a nice explicit expression:

$$(1.1.13) \quad \mathcal{M}_z([e^{\mathbf{i}t\zeta} g]) = \mathcal{M}_z([g]) + \int_0^t \langle \mu(e^{\mathbf{i}s\zeta} gz), \zeta \rangle ds.$$

We conclude from (1.1.9) and (1.1.10) that the maximal weights of the 1-parameter subgroups in G^c control the asymptotic behavior of \mathcal{M}_z along geodesics (a priori

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only of those arising from 1-parameter subgroups, but in fact of all of them by [56, Lemma 2.4.8]). We go now into the proof of the lemma.

Lemma 1.1.7 (Kempf-Ness). *A point $z \in Z$ is polystable if and only if $G^c \cdot z$ contains a zero of the moment map. There is at most one G -orbit inside each G^c -orbit on which the moment map vanishes.*

PROOF. Suppose first that $G^z \cdot z$ contains a zero of the moment map. We will show that z is polystable using a refinement of the Hilbert-Mumford criterion due to F. Kirwan (see [39] pg. 107), that shows that it is sufficient to consider 1-parameter subgroups $\lambda : \mathbb{C}^* \rightarrow G^c$ with $\lambda(S^1) \subset G$. Let λ be such a 1-parameter subgroup and denote by $\zeta \in \mathfrak{g}$ the generator of the S^1 action. Then, from (1.1.9), (1.1.10) and $\mu(z) = 0$, we have $w(z, \lambda) \geq 0$, since

$$\frac{d}{dt} \langle \mu(e^{it\zeta} z), \zeta \rangle = \omega(Y_\zeta, JY_\zeta) = \|Y_\zeta\|^2.$$

Moreover, $w(z, \lambda) = 0$ only if $Y_{\zeta|e^{it\zeta}} = 0$ for all $t > 0$. In particular $Y_{\zeta|z} = 0$ which implies that λ fixes z . Suppose now that $z \in Z$ is polystable. We claim that \mathcal{M}_z attains a minimum on $G \setminus G^c$. Using the polar decomposition of G^c any element $g \in G^c$ can be written as $g = ke^{i\zeta}$, with $k \in G$ and $\zeta \in \mathfrak{g}$. Combined with (1.1.13), we thus obtain the equality

$$(1.1.14) \quad \begin{aligned} \mathcal{M}_z([g]) &= \int_0^1 \langle \mu(e^{it\zeta} z), \zeta \rangle dt = \int_0^1 \frac{i}{2\pi} \frac{h(e^{it\zeta} \hat{z}, \rho_*(\zeta) e^{it\zeta} \hat{z})}{|e^{it\zeta} \hat{z}|^2} dt \\ &= \frac{1}{\pi} \log(|e^{it\zeta} \hat{z}|^2)|_0^1 = \frac{2}{\pi} (\log |g\hat{z}| - \log |\hat{z}|). \end{aligned}$$

Let $\{[g_j]\}_{j=0}^\infty$ be a minimizing sequence for \mathcal{M}_z on $G \setminus G^c$. We take representatives $e^{i\zeta_j}$ of $[g_j]$, with $\zeta_j \in \mathfrak{g}$. By (1.1.14), the sequence $e^{i\zeta_j} \hat{z}$ is bounded on V , so we can suppose that it converges to \hat{z}_0 (by taking a convergent subsequence). Since z is polystable \hat{z}_0 is in the G^c -orbit of \hat{z} , i.e. there exists $g \in G^c$ such that $[g]$ attains the minimum as claimed. By property (1.1.11) and G -equivariance of the moment map we conclude that $G \cdot gz \subset \mu^{-1}(0)$. Finally, let $z, z' \in G^c$ such that $\mu(z) = \mu(z') = 0$. Then $z' = gz$ for some $g \in G^c$ and $[1], [g] \in G \setminus G^c$ are critical points of \mathcal{M}_z . Let $g = ke^{i\zeta}$ be the polar decomposition of g . The second derivative of \mathcal{M}_z along $t \rightarrow [e^{it\zeta}]$ must vanish since

$$\frac{d^2}{dt^2} \mathcal{M}_z([e^{it\zeta}]) = \langle \mu(e^{it\zeta} z), \zeta \rangle = \omega(Y_\zeta, JY_\zeta) = \|Y_\zeta\|^2.$$

This implies that $e^{i\zeta} \in G_z^c$, so $z' = gz = ke^{i\zeta} z = kz \in G \cdot z$ as claimed. \square

The proof of the Kempf-Ness theorem is now straightforward. By Lemma 1.1.7 we have an inclusion $\mu^{-1}(0) \subset Z^{ps} \subset Z^{ss}$ that induces a continuous map

$$i : \mu^{-1}(0)/G \rightarrow Z^{ss}/G^c.$$

Since any semistable orbit in the GIT quotient is glued to the unique polystable orbit on its closure, this map is surjective (again by Lemma 1.1.7). If $i([z]) = i([z'])$ then $z' = gz$ for some $g \in G^c$, since both points are polystable and by Lemma 1.1.7 their G -orbits must coincide. Therefore, i is a continuous bijection from a compact space to a Hausdorff space and hence it is an homeomorphism.

1.1.4. The non projective case. The identification between the symplectic and algebraic quotients given by Kempf-Ness Theorem can be extended to an arbitrary Kähler manifold (neither necessarily compact, nor projective) in a suitable sense. We follow [58] for this topic (see also references therein). It is in this setting where the correspondence concerning the stable G^c orbits and the zeros locus of the moment map (as in Lemma (1.1.7)) provides a general framework that formally applies in many situations where the objects involved are infinite dimensional.

Let (Z, ω, J) be a Kähler manifold endowed with a Hamiltonian left action of a compact group G . Let us fix a moment map μ and suppose that G acts by holomorphic isometries. By a theorem of Guillemin and Sternberg (see [33]), the action of G extends to a unique holomorphic action of G^c on Z . The action of G^c , however, no longer preserves the symplectic form of Z nor does it preserve the zero level set $\mu^{-1}(0)$. The correspondence (1.1.2) in the projective case motivates the search for a result analogue to Lemma (1.1.7) in this general setup. A partial answer to this problem was given in [57], where a suitable notion of stability, called analytic stability, was defined using the right hand side of formula (1.1.10) as the definition of the maximal weight. A point $z \in Z$ is then called *analytically stable* if the maximal weight

$$(1.1.15) \quad w(z, \zeta) = \lim_{t \rightarrow \infty} \langle \mu(e^{it\zeta} z), \zeta \rangle$$

is strictly positive for any non-zero $\zeta \in \mathfrak{g}$. The notion of analytically stable is a priori stronger than GIT stable in the projective case but they are in fact equivalent (a proof can be found in [56, Lemma 2.4.8]). In general it is not straightforward that analytic stability is an invariant of the G^c -orbit of a point. This is proved in [57, Theorem 5.4] as a consequence of the characterization of analytic stability in terms of the so-called linear properness of the integral of the moment map \mathcal{M}_z , which is constructed exactly as in the projective case. Recall that \mathcal{M}_z is said to be linearly proper if

$$(1.1.16) \quad |\zeta| \leq C_1 \mathcal{M}_z([e^{i\zeta}]) + C_2,$$

for all $\zeta \in \mathfrak{g}$, suitable norm $|\cdot|$ in \mathfrak{g} , and positive real constants $0 < C_1, C_2 \in \mathbb{R}$. The general result concerning this notion of stability can be stated as follows (see [57, Theorem 5.4]).

Theorem 1.1.8 ([57]). *A point $z \in Z$ is analytically stable if and only if the stabiliser G_z^c of z in G^c is finite and $G^c \cdot z \cap \mu^{-1}(0) \neq \emptyset$. Furthermore, there is at most one G -orbit inside each G^c -orbit on which the moment map vanishes.*

The proof follows as the one given for Lemma 1.1.7 in the previous section. The only point that slightly differs is the “only if” part, but it can be easily adapted using the linear properness of \mathcal{M}_z (compare (1.1.14) with (1.1.16)). Generalizing previous work of A. Teleman, M. Riera [58] has extended the notion of analytic stability to the polystable case. He imposes some mild restrictions on the growth of μ and the G -action. The maximal weights are viewed as defining a collection of maps, for each $z \in Z$, $\lambda_z : \partial(G \setminus G^c) \rightarrow \mathbb{R}$, where $\partial(G \setminus G^c)$ is the boundary at infinity of the symmetric space $G \setminus G^c$. A suitable result, similar to Lemma 1.1.7, is proved for this notion of stability (see [58, Theorem 1.2]). It is very likely that this new notion of analytic polystability will be relevant for our problem and also for the cscK problem, where no

notion of analytic polystability is known (analytic (semi)stability was introduced by X. X. Chen in [18] with the name of *geodesic (semi)stability*).

1.2. The Hitchin-Kobayashi correspondence for the HYM equation

In this section we recall the moment map interpretation of the HYM equation and explain how the general framework provided by §1.1 applies, leading to a Hitchin-Kobayashi correspondence. First we give a short historical review of the problem concerning the HYM equation.

1.2.1. Overview. In 1965 Narasimhan and Shesadri proved that holomorphic bundles over a compact Riemann surface X are stable if and only if they arise from irreducible unitary representations of $\pi_1(X)$. This connection between holomorphic and unitary structures, that in the classical theory of line bundles is essentially equivalent to the identification between holomorphic and harmonic 1-forms, was studied by Atiyah and Bott in [6] from a gauge-theoretic point of view. They identified the stable bundles over X with the irreducible, projectively flat unitary connections and pointed out that the curvature of a connection can be viewed as a moment map for the action of the gauge group. This interpretation led them to consider the Yang-Mills functional as the squared norm of this moment map and the correspondence provided by the Narashiman-Seshadri Theorem as an infinite dimensional version of Kempf-Ness theorem relating symplectic reduction with GIT quotients. The proper generalization of this fact for holomorphic vector bundles over a Kähler manifold of arbitrary dimension was suggested independently by Hitchin and Kobayashi. They conjectured that the stability of a holomorphic bundle should be related to the existence of a unitary Hermitian-Yang-Mills (HYM) connection (equivalently Hermite-Einstein metric). The Hitchin-Kobayashi correspondence in the case of the HYM equation for $\mathrm{GL}(r, \mathbb{C})$ -bundles is a theorem of Donaldson (for complex projective varieties) and Uhlenbeck and Yau (for compact Kähler manifolds) [21, 79], generalized by Ramanathan and Subramanian ([62]) in the algebraic case for principal bundles with arbitrary complex reductive Lie group.

1.2.2. Moment map interpretation. In this section we review the notion of Hermitian-Yang-Mills connection, explaining its moment map interpretation. Further details on this topic are given in sections §2.2 and §2.3.

Let X be a compact symplectic manifold of dimension $2n$, with symplectic form ω and volume form $\mathrm{vol}_\omega = \frac{\omega^n}{n!}$. We fix a real compact Lie group G with Lie algebra \mathfrak{g} and a smooth principal G -bundle E over X . We also fix a positive definite inner product

$$(\cdot, \cdot) : \mathfrak{g} \times \mathfrak{g} \rightarrow \mathbb{R}$$

on \mathfrak{g} invariant under the adjoint action. Recall from the introduction that this metric gives a pairing $\Omega^p(\mathrm{ad} E) \times \Omega^q(\mathrm{ad} E) \rightarrow \Omega^{p+q} : (a, b) \mapsto (a \wedge b)$. A connection on E is a G -equivariant splitting $A : TE \rightarrow VE$ of the short exact sequence

$$0 \rightarrow VE \longrightarrow TE \longrightarrow \pi^* TX \rightarrow 0,$$

where $VE = \ker d\pi$ is the vertical bundle. It follows by definition that the space \mathcal{A} of principal connections on E is an affine space modelled on $\Omega^1(\mathrm{ad} E)$. Let denote by \mathcal{G}

the *gauge group* of the bundle E , i.e. the group of G -equivariant diffeomorphisms of E that project to the identity on X . The action of $g \in \mathcal{G}$ on TE defines a left \mathcal{G} -action on \mathcal{A} given by $g \cdot A := g \circ A \circ g^{-1}$ and the 2-form on \mathcal{A} given by

$$(1.2.1) \quad \omega_{\mathcal{A}}(a, b) = \int_X (a \wedge b) \wedge \omega^{n-1},$$

for $a, b \in T_A \mathcal{A} = \Omega^1(\text{ad } E)$, with $A \in \mathcal{A}$, is a symplectic structure preserved by \mathcal{G} . It was first observed by Atiyah and Bott [6] when X is a Riemann surface and by Donaldson [21] in higher dimensions that the \mathcal{G} -action on \mathcal{A} is hamiltonian, with \mathcal{G} -equivariant moment map $\mu_{\mathcal{G}}: \mathcal{A} \rightarrow (\text{Lie } \mathcal{G})^*$ given by

$$(1.2.2) \quad \langle \mu_{\mathcal{G}}(A), \zeta \rangle = \int_X (\zeta \wedge F_A) \wedge \omega^{n-1},$$

where $\zeta \in \text{Lie } \mathcal{G}$ and $F_A \in \Omega^2(\text{ad } E)$ is the curvature of $A \in \mathcal{A}$. Note here that the \mathcal{G} -invariant bilinear pairing in $\text{Lie } \mathcal{G} = \Omega^0(\text{ad } E)$ given by

$$(1.2.3) \quad \langle \zeta_1, \zeta_2 \rangle_{\mathcal{G}} := \int_X (\zeta_1 \wedge \zeta_2) \text{vol}_{\omega},$$

where $\zeta_1, \zeta_2 \in \text{Lie } \mathcal{G}$, allows us to identify (up to multiplicative constant factors) the image of the moment map (1.2.2) with $\Lambda_{\omega} F_A \in \text{Lie } \mathcal{G}$.

We now suppose that X is a Kähler manifold with Kähler form ω and consider the associated principal G^c -bundle $E^c = E \times_G G^c$, where the group G^c denotes the complexification of G . There is a \mathcal{G} -invariant distinguished subspace $\mathcal{A}^{1,1} \subseteq \mathcal{A}$ of connections A with $F_A \in \Omega^{1,1}(\text{ad } E)$. This space is in correspondence with the space of holomorphic structures on the principal G^c -bundle E^c over the complex manifold X . Away from the singularities, it inherits a complex structure compatible with $\omega_{\mathcal{A}}$ and hence a Kähler structure. Thus, the smooth locus of $\mathcal{A}^{1,1}$ is a Kähler manifold endowed with a Hamiltonian \mathcal{G} -action and we can consider the Kähler reduction

$$(1.2.4) \quad \mu_{\mathcal{G}}^{-1}(z)/\mathcal{G},$$

where $\mu_{\mathcal{G}}$ is now the restriction of the moment map to $\mathcal{A}^{1,1}$ and z is an element of the centre \mathfrak{z} of \mathfrak{g} . Note here that any element $z \in \mathfrak{z}$ determines a \mathcal{G} -invariant element in $\text{Lie } \mathcal{G}$ and we use the identification $\mu_{\mathcal{G}}(A) \in \text{Lie } \mathcal{G}$, for $A \in \mathcal{A}$, provided by the \mathcal{G} -invariant pairing $\langle \cdot, \cdot \rangle_{\mathcal{G}}$ to make sense of (1.2.4). The elements in $\mu_{\mathcal{G}}^{-1}(z)$ for $z \in \mathfrak{z}$ are the *Hermitian-Yang-Mills* connections that we recall in the following definition.

Definition 1.2.1. A connection $A \in \mathcal{A}$ is called *Hermitian-Yang-Mills* if $F_A^{0,2} = 0$ and it satisfies the *Hermitian-Yang-Mills equation*

$$(1.2.5) \quad \Lambda F_A = z, \quad z \in \mathfrak{z}.$$

The element $z \in \mathfrak{z}$ in the right-hand side of (1.2.5) is determined by the cohomology class $\Omega := [\omega] \in H^2(X)$ and the topology of the principal bundle E . This follows because applying (z_j, \cdot) to (1.2.5), for an orthonormal basis $\{z_j\}$ of $\mathfrak{z} \subset \mathfrak{g}$, and then integrating over X , we obtain

$$(1.2.6) \quad z = \sum_j \frac{\langle z_j(E) \cup \Omega^{[n-1]}, [X] \rangle}{\text{Vol}(\Omega)} z_j.$$

Here, $\Omega^{[k]} := \Omega^k/k!$, $\text{Vol}(\Omega) := \langle \Omega^{[n]}, [X] \rangle$ and $z_j(E) := [z_j \wedge F_A] \in H^2(X)$ is the Chern–Weyl class associated to the G -invariant linear form (z_j, \cdot) on \mathfrak{g} , which only depends on the topology of the bundle E (see [42, Ch XII, §1]). When the associated Kähler reduction (1.2.4) of $\mathcal{A}^{1,1}$ is non-empty it turns out that this reduction has finite dimension and inherits, in the smooth locus, a structure of Kähler manifold (see [40, Corollary 3.32]).

1.2.3. The correspondence. Consider the space of isomorphism classes of holomorphic structures on the smooth principal G^c -bundle E^c over the complex manifold X , with fixed Kähler form ω . Set theoretically this space can be realized as the quotient of $\mathcal{A}^{1,1}$ by the left action of the gauge group \mathcal{G}^c of E^c , often referred as the *complex gauge group*. The action of $g \in \mathcal{G}^c$ on $\mathcal{A}^{1,1}$ can be thought of as the push-forward of the (integrable) almost complex structure on $\text{Tot}(E^c)$ defined by any element $A \in \mathcal{A}^{1,1}$. The notation for the complex gauge group is not accidental, since it can be considered a complexification of \mathcal{G} . We are thus in position to formulate an infinite dimensional version of Lemma 1.1.7 characterizing \mathcal{G}^c -orbits in $\mathcal{A}^{1,1}$ that contain a HYM connection. This correspondence is often referred as *the Hitchin-Kobayashi correspondence* for the HYM equation [21, 79, 62].

Following [57] we define stability in terms of maximal weights. A connection $A \in \mathcal{A}^{1,1}$ is *stable* if

$$w(A, \zeta) = \lim_{t \rightarrow \infty} \langle \mu_{\mathcal{G}}(e^{it\zeta} A), \zeta \rangle > 0$$

for any non-zero ζ in a suitable Sobolev completion of $\text{Lie } \mathcal{G} \cong \Omega^0(\text{ad } E)$ (see [57, Section 4.1.4]). The picture provided by §1.1 applies formally here giving an integral of the moment map $\mathcal{M}_A : \mathcal{G} \setminus \mathcal{G}^c \rightarrow \mathbb{R}$, the Donaldson functional in the case of vector bundles [21], whose critical points can be identified with zeros of $\mu_{\mathcal{G}}$. The existence of the functional is due to the fact that $\mathcal{G} \setminus \mathcal{G}^c$ is contractible. This functional is convex along geodesics $t \mapsto [e^{it\zeta} g]$ in the infinite dimensional symmetric space $\mathcal{G} \setminus \mathcal{G}^c$ and its behaviour at infinity is controlled by the corresponding maximal weights. Recall that a holomorphic G^c -bundle E^c is simple if the space of holomorphic global sections of $\text{ad } E^c$ coincides with the centre of \mathfrak{g}^c . Identifying $\mathcal{G} \setminus \mathcal{G}^c$ with the space of smooth G -reductions of the holomorphic bundle determined by A we can state the Hitchin-Kobayashi correspondence as follows (see [62, Remark 4])

Theorem 1.2.2 ([24, 62, 79]). *A holomorphic G^c -bundle is stable if and only if it is simple and there exists a unique G -reduction whose associated Chern connection is Hermite-Yang-Mills.*

The proof of the Hitchin-Kobayashi correspondence follows formally as in the finite dimensional case: uniqueness follows from the existence of geodesics in $\mathcal{G} \setminus \mathcal{G}^c$ joining any two points and the convexity of \mathcal{M}_A along them. The “if” part is the easy one and follows as in the proof of Lemma 1.1.7. For the “only if” part, a theorem of Uhlenbeck and Yau [79] on weak subbundles of vector bundles allows to prove that A is stable if and only if \mathcal{M}_A is linearly proper (see [57, Lemma 6.1]). Then, by taking a minimizing sequence and proving regularity the latter is equivalent to attaining a minimum in $\mathcal{G} \setminus \mathcal{G}^c$, thus achieving the proof of Theorem 1.2.2.

Remark 1.2.3. Note that the statement of Theorem 1.2.2 for the holomorphic structure determined by $A \in \mathcal{A}^{1,1}$ is analogue to that of Theorem 1.1.8. To see this note that a G -reduction (a Hermitian metric when $G = U(r)$) determines $g \in \mathcal{G}^c$. Then, the associated Chern connection is HYM if and only if so is $g \cdot A$.

1.2.4. Mumford-Takemoto slope stability. In this section we restrict to the algebraic setting and explain why the analytic stability condition defined in §1.2.3 coincides with the slope stability for coherent sheaves due to D. Mumford and F. Takemoto. First we recall the definition of slope stability for coherent sheaves over a polarised scheme.

Let X be an n -dimensional projective scheme with an ample invertible sheaf L . Then the *slope* $\mu_L(E) = d_1(E)/d_0(E)$ of a non-zero coherent sheaf E with n -dimensional support is defined (up to an additive factor independent of E) as the ratio of the first two leading coefficients of its Hilbert polynomial with respect to L (cf. e.g. [37, §1.2]), given by

$$P_L(E, k) := \sum_{p=0}^n (-1)^p \dim H^p(X, E \otimes L^k) = k^n d_0(E) + k^{n-1} d_1(E) + \cdots + d_n(E).$$

Following Simpson [65], a coherent sheaf on a scheme has *pure dimension* k if it has k -dimensional support and no subsheaves have lower dimensional support. Let E be a coherent sheaf on X which is pure of dimension n . Then E is called *Mumford semistable* (with respect to L) if

$$\mu_L(F) \leq \mu_L(E)$$

for each proper subsheaf $F \subset E$. Such an E is *Mumford stable* if the inequality is always strict and *Mumford polystable* if it is a direct sum of stable sheaves, all of them with the same slope.

Suppose now that (X, L) is a smooth compact polarised variety and let E be a holomorphic vector bundle (locally free sheaf) over X . As usual holomorphic $GL(r, \mathbb{C})$ -principal bundles are identified with holomorphic vector bundles over X via the standard representation on \mathbb{C}^r . Fix some Hermitian metric on E and suppose that a *weak endomorphism* ζ destabilizes the Chern connection A , i.e. has negative maximal weight, in the sense of the previous section. Then it has almost everywhere constant eigenvalues and defines a filtration of E by holomorphic subbundles in the complement of a complex codimension 2 subvariety of X . Using a theorem of Uhlenbeck and Yau [79] on weak subbundles of vector bundles this leads to a filtration of E by reflexive (coherent) subsheaves (see [79]) and the maximal weight $w(A, \zeta)$ has then an explicit expression in terms of the slopes of the elements of the filtration (see [57, Lemma 4.2]). From this it follows that both definitions are equivalent in the stable and semistable case and so the Theorem 1.2.2 can be stated in this context as follows (see [21, Corollary 4]): A holomorphic vector bundle E is polystable if and only if there exists a Hermitian metric on E whose associated Chern connection is Hermite-Yang-Mills.

1.3. The cscK problem

In this section we briefly review the state of the art of the cscK problem, centering the attention in those aspects relevant for the contents of this thesis. We discuss the moment map interpretation of the scalar curvature of a Kähler metric [28, 22] and how the problem fits into the general framework of Section §1.1 [24].

1.3.1. Overview. In the 1980's E. Calabi [15] settled the problem of finding Kähler metrics, within a fixed Kähler class on a compact complex manifold, which are critical points of the L^2 norm of the scalar curvature. These metrics g are called *extremals*, and those satisfying the cscK equation

$$(1.3.1) \quad S_g = \text{const.}$$

are absolute minimizers of the action. The notion of cscK metric can be considered the natural generalization of the Kähler Einstein (KE) condition when the fixed Kähler class is not a multiple of the first Chern class $c_1(X)$. The first result concerning the cscK equation (1.3.1) is the classical uniformisation theorem, which states that every compact Riemann surface admits a metric of constant curvature, unique modulo holomorphic automorphisms if we prescribe the total area. This result establishes a correspondence between cscK metrics and complex structures for compact surfaces. In higher dimensions this correspondence is not true, and some stability of the complex manifold is necessary in order to carry a cscK metric. The solution of the Hitchin-Kobayashi correspondence for the HYM equation gives a clue of this fact, since the existence of a KE metric on X implies polystability of the holomorphic tangent bundle.

A correspondence between stable complex structures and cscK metrics was first conjectured by S. T. Yau for the KE case. The problem of whether a compact complex manifold admits a KE metric was solved in the case $c_1(X) \leq 0$ in the 1970's with positive answer. The celebrated solution of Yau to the Calabi conjecture established the existence of a Ricci flat metric, now named Calabi-Yau metric, in each Kähler class on a compact Kähler manifold X with $c_1(X) = 0$. If $c_1(X) < 0$, the existence of KE metrics was proved by Aubin and Yau, independently. When $c_1(X) > 0$ it was known for a long time that there are obstructions to the existence of solutions coming from the holomorphic vector fields on the manifold (see [30, 15]). A remarkable obstruction for the existence of cscK metrics in a given Kähler class was defined by A. Futaki in [29] for the case of KE metrics. The *Futaki invariant* is a character of the Lie algebra of the automorphism group of the complex manifold associated to any Kähler class. In [72] G. Tian proved that any compact complex surface with $c_1(X) > 0$ admits a KE metric if and only if it has vanishing Futaki invariant. In [75] the same author proved that the existence of KE metrics in arbitrary dimensions is equivalent to a suitable properness of the *Mabuchi K-energy* [48], the integral of the moment map.

In the analytic setting concerning cscK and extremal metrics, i.e. for arbitrary Kähler class, several achievements have taken place during the last years. In [19] G. Tian and X. X. Chen have proved uniqueness of solutions for the extremal case in full generality and that the Mabuchi K-energy is bounded below if there exists a cscK metric in the Kähler class. Their work is based on the study of the geodesic equation in the space of Kähler potentials introduced independently by A. Mabuchi [48] and S.

K. Donaldson [23]. In [18] Chen defined a suitable notion of analytic stability, called *geodesic stability*, and proved that it is an obstruction for the existence of cscK metrics.

For the case of polarised Kähler manifolds (X, L) , i.e. when the Kähler class is $c_1(L)$ for an ample line bundle L , a new obstruction concerning GIT called K -stability was defined by Tian in [74] for the KE case. Tian conjectured that a Kähler manifold with positive first Chern class admits a KE metric if and only if it is K -stable and proved the “only if” part. In [24] S. K. Donaldson generalized the notion of K -stability to any polarised manifold (X, L) introducing the notion of (algebraic) test configuration. Roughly speaking, algebraic test configurations are degenerations of the polarized manifold to general schemes with a rational number attached (the *generalized Futaki invariant*) which will be positive if the manifold is stable. We know that K -stability is an obstruction for the existence of cscK metrics in polarized manifolds mainly due to the work of Donaldson [25], T. Mabuchi [49] and J. Stoppa [67]. Donaldson has proved the converse for the case of toric surfaces [25]. Recent work of Arezzo, Della and La Nave in [3], based on previous work of Ross and Thomas in [71], has simplified the notion of K -stability in certain cases showing that only suitable nice degenerations need to be considered. However, the work of Apostolov, Calderbank, Gauduchon, Tønnesen and Friedman in [2] and of G. Székelyhidi in [69] concerning extremal metrics suggests that the algebraic notion of K -stability should be strengthened for the general case, allowing more general analytic (instead of just algebraic) test configurations.

1.3.2. The scalar curvature as a moment map. We now briefly explain the moment map interpretation of the scalar curvature. This fact was first observed by Fujiki [28] for the scalar curvature of Kähler manifolds and independently by Donaldson [22] for the Hermitian scalar curvature of almost Kähler manifolds. We hasten to add that both quantities coincide in the Kähler case, up to positive multiplicative constant factor. In this section we follow closely Donaldson’s approach [22].

Throughout this section X is a fixed compact symplectic manifold of dimension $2n$, with symplectic form ω . First we recall the notion of Hermitian scalar curvature of an almost Kähler manifold. An almost complex structure $J: TX \rightarrow TX$ is compatible with ω if $\omega(J\cdot, J\cdot) = \omega$ and $g_J = \omega(\cdot, J\cdot)$ is a Riemannian metric in X . Any compatible almost complex structure J defines a hermitian metric on T^*X and there is a unique unitary connection on T^*X whose $(0, 1)$ component is the operator

$$\bar{\partial}_J: \Omega_J^{1,0} \rightarrow \Omega_J^{1,1}$$

induced by J . We denote by $\Omega_J^{p,q}$ the space of (p, q) -type forms with respect to the decomposition provided by J . The real 2-form ρ_J is defined as $-\mathbf{i}$ times the curvature form of the induced connection on the canonical line bundle $K_X = \Lambda_{\mathbb{C}}^n T^*X$. The *Hermitian scalar curvature* S_J is the real function on X defined by

$$(1.3.2) \quad S_J \omega^n = n \rho_J \wedge \omega^{n-1}.$$

Note that the average \hat{S} of the Hermitian scalar curvature does not depend on J but only on the cohomology class of ω ,

$$(1.3.3) \quad \hat{S} := \frac{\int_X S_J \text{vol}_\omega}{\int_X \text{vol}_\omega} = 2\pi n \cdot \frac{\langle c_1(X) \cup [\omega]^{n-1}, [X] \rangle}{\text{Vol}([\omega])}.$$

The space \mathcal{J} of almost complex structures on X which are compatible with ω has a natural structure of infinite dimensional Kähler manifold. The Kähler form $\omega_{\mathcal{J}}$ is

$$(1.3.4) \quad \omega_{\mathcal{J}}(\dot{J}_1, \dot{J}_2) := \int_X \text{tr}(J \dot{J}_1 \dot{J}_2) \text{vol}_\omega, \text{ for } \dot{J}_1, \dot{J}_2 \in T_J \mathcal{J},$$

whereas the complex structure $\mathbf{J}: T_J \mathcal{J} \rightarrow T_J \mathcal{J}$ is $\mathbf{J}\dot{J} := J\dot{J}$. Here we are identifying the tangent space $T_J \mathcal{J}$ at $J \in \mathcal{J}$ with the set of endomorphisms $\dot{J}: TX \rightarrow TX$ such that $\dot{J}J = -J\dot{J}$ and \dot{J} is symmetric with respect to the induced metric $g_J = \omega(\cdot, J\cdot)$.

The group $\mathcal{H} \subset \text{Diff } X$ of Hamiltonian symplectomorphisms consists of those diffeomorphisms $h: X \rightarrow X$ such that there exists a Hamiltonian isotopy

$$[0, 1] \times X \rightarrow X: (t, x) \mapsto h_t(x)$$

from $h_0 = \text{Id}$ to $h_1 = h$ (see e.g. [55, §3.2]). Note that the Lie algebra $\text{Lie } \mathcal{H}$ of \mathcal{H} , consisting of the hamiltonian vector fields on X , can be identified with the Lie algebra $C^\infty(X)/\mathbb{R}$ of smooth real functions on X modulo the constants, endowed with the Poisson bracket. There is a natural \mathcal{H} -action on \mathcal{J} by push-forward, that is $h \cdot J := h_* J = dh \circ J \circ dh^{-1}$, preserving the Kähler structure and, as was proved by Donaldson in [22, Proposition 9], it is hamiltonian, with equivariant moment map $\mu_{\mathcal{H}}: \mathcal{J} \rightarrow \text{Lie } \mathcal{H}^*$ given by

$$(1.3.5) \quad \mu_{\mathcal{H}}(J) = -(S_J - \hat{S}).$$

Here we use the L^2 pairing provided by vol_ω to identify the set $C_0^\infty(X)$ of functions with zero integral over X , with the dual of $\text{Lie } \mathcal{H} = C^\infty(X)/\mathbb{R}$.

The \mathcal{H} -invariant subspace $\mathcal{J}^i \subseteq \mathcal{J}$ of integrable almost complex structures is (away from the singularities) a complex submanifold so it inherits a Kähler structure. Thus, \mathcal{J}^i is a symplectic manifold endowed with a Hamiltonian \mathcal{H} -action and we can construct the Kähler reduction

$$(1.3.6) \quad \mu_{\mathcal{H}}^{-1}(0)/\mathcal{H}$$

where $\mu_{\mathcal{H}}$ is now the restriction of the moment map to \mathcal{J}^i . Over \mathcal{J}^i , the Hermitian scalar curvature S_J is, up to a multiplicative constant factor, the usual Riemannian scalar curvature of the Kähler metric determined by J and ω . We conclude that the Kähler reduction (1.3.6) is the moduli space of Kähler metrics with fixed Kähler form ω and constant scalar curvature.

1.3.3. Complexified orbits. We discuss now how the cscK problem fits into the general framework provided by Section §1.1. We follow the argument given by Donaldson in [23, Section 4].

Let (X, J) be a compact complex manifold where we fix a Kähler class $[\omega]$. Let \mathcal{K} be the space of Kähler metrics in (X, J) with this Kähler class. We fix a Kähler form ω and consider the corresponding group of Hamiltonian symplectomorphisms

and the infinite dimensional Kähler manifold \mathcal{J}^i of compatible complex structures. For simplicity we will suppose that the first cohomology group $H^1(X)$ vanishes, so the group of Hamiltonian symplectomorphisms \mathcal{H} coincides with the identity component of the group of symplectomorphisms. As was pointed out by Donaldson in [23], the main obstacle to apply the picture of §1.1 formally in this situation comes from the fact that \mathcal{H} does not admit a complexification (a reason for this will be given later in this section). However, there is an integrable distribution $D \subseteq \mathcal{J}^i$ whose leaves play the role of the orbits under the action of an ideal complexification. At a point $J \in \mathcal{J}^i$ this distribution is given by

$$(1.3.7) \quad D_J = \{Y_{\eta|J}; \eta \in \text{Lie } \mathcal{H}^c\},$$

where $\text{Lie } \mathcal{H}^c = \text{Lie } \mathcal{H} \otimes \mathbb{C}$ is the complexification of $\text{Lie } \mathcal{H}$ and

$$(1.3.8) \quad Y_\eta = Y_{\eta_0} + \mathbf{J}Y_{\eta_1} \text{ for all } \eta = \eta_0 + i\eta_1 \in \text{Lie } \mathcal{H}^c.$$

Recall that \mathbf{J} is the complex structure on \mathcal{J}^i . The distribution D is closed under the Lie bracket in \mathcal{J}^i , since the Nijenhuis tensor of \mathbf{J} vanishes, although it is not obvious that this fact implies the existence of leaves integrating D through every point $J \in \mathcal{J}^i$ (because \mathcal{J}^i is infinite dimensional). To see that D indeed integrates we consider the principal \mathcal{H} -bundle

$$(1.3.9) \quad \mathcal{Y} \rightarrow \mathcal{K},$$

given by pairs $(\omega', f) \in \mathcal{K} \times \text{Diff}_0(X)$, such that $f^*\omega' = \omega$. Here we use the assumption $H^1(X) = 0$ to have \mathcal{H} as the structure group of \mathcal{Y} , instead of the identity component of the group of symplectomorphisms. The map

$$(1.3.10) \quad \tau_J: \mathcal{Y} \rightarrow \mathcal{J}^i: (\omega', f) \mapsto f^*J,$$

plays the role of the complexified action, in the sense that the image $\tau_J(\mathcal{Y})$ of this map is a leaf of D through J (for this last fact it is important to work with \mathcal{J}^i instead of the whole \mathcal{J}). The image of τ_J can be then thought of as the orbit of an ideal complexification of \mathcal{H} acting on \mathcal{J}^i . It follows by definition of the complexified orbit $\tau_J(\mathcal{Y})$ of a point $J \in \mathcal{J}^i$ that it contains a zero of the \mathcal{H} -equivariant moment map $\mu_{\mathcal{H}}$ (1.3.5) if and only if there exists a cscK metric in \mathcal{K} .

The principal bundle \mathcal{Y} is endowed with a canonical connection that provides the link between the complex and the symplectic points of view (see [23, Section 2]). To define this connection on \mathcal{Y} we consider a smooth curve

$$\omega_t = \omega_{\phi_t} = \omega + dd^c\phi_t \subset \mathcal{K}$$

and define its horizontal lift. The horizontal lift f_t through $(\omega_0, f) \in \mathcal{Y}_J$ is given by the Moser curve of ω_t , i.e. the curve of diffeomorphisms with initial condition $f_0 = f$, integrating the vector $-Jy_t$ defined by

$$i_{y_t}\omega_t = d\dot{\phi}_t.$$

It is straightforward that f_t satisfies $f_t^*\omega_t = \omega_0$ for all t , as required. The curve of metrics ω_t determines another one $J_t = \tau_J(\omega_t, f_t) \subset \mathcal{J}^i$ which satisfies

$$(1.3.11) \quad \dot{J}_t = \frac{d}{dt}f_t^*J = f_t^*(L_{-Jy_t}J) = -J_t L_{f_t^*y_t}J_t = \mathbf{J}Y_{\eta_t},$$

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where \mathbf{J} is the complex structure on \mathcal{J}^i and Y_{η_t} is the infinitesimal action of the Hamiltonian vector field $\eta_t = y_{\dot{\phi}_t \circ f_t} \in \text{Lie } \mathcal{H}$. Special curves ϕ_t in the space of Kähler potentials are those for which $\dot{\phi}_t \circ f_t$ remains constant, since they determine curves J_t in \mathcal{J}^i via the horizontal lift that follow a pure imaginary direction $\mathbf{J}Y_{\eta_0}$ of the missed complexification of \mathcal{H} . This last condition is equivalent to

$$(1.3.12) \quad \ddot{\phi}_t - |d\dot{\phi}_t|_{\omega_t}^2 = 0,$$

that can be interpreted as the equation of geodesics with respect to suitable negatively curved metric in the space of Kähler potentials, as proved independently by T. Mabuchi [48] and S. K. Donaldson [23]. The analogue formulation of the geodesic equation in \mathcal{K} is obtained considering (1.3.12) modulo the constant functions. A deep analytic study of equation (1.3.12) has been carried out by X. X. Chen and G. Tian in [19], where the authors proved the existence of a unique $C^{1,1}$ geodesic (with bounded mixed second derivatives) joining any two points in \mathcal{K} and that this regularity is sharp. The lack of regularity of the solutions gives a clue of why the \mathcal{H} -action in \mathcal{J}^i cannot be complexified (at least in the C^∞ -category).

Example 1.3.1. We discuss now a well known example of solution to (1.3.12) provided by holomorphic isometries in (X, ω, J) . Let $\eta \in \text{Lie } \mathcal{H}_J$, the Lie algebra of the isotropy group of J in \mathcal{H} (equivalently, the Hamiltonian Killing vector fields on (X, J, ω)). Given any vector field y on X we denote by f_t^y the flow of y on X . We claim that

$$\phi_t = - \int_0^t \phi \circ (f_t^{J\eta})^{-1},$$

is a solution to (1.3.12). Note first that $\omega_t = \omega_{\phi_t} = (f_t^{J\eta})_* \omega$, since

$$\frac{d}{dt} (f_t^{J\eta})_* \omega = (f_t^{J\eta})_* d(i_{-J\eta} \omega) = -(f_t^{J\eta})_* (dd^c \phi) = dd^c \dot{\phi}_t.$$

This implies that $\eta = -y_{\dot{\phi}_t, \omega_t}$ from which it follows that

$$\ddot{\phi}_t = (f_t^{J\eta})_* d\phi(J\eta) = d\dot{\phi}_t (Jy_{\dot{\phi}_t, \omega_t}) = |d\dot{\phi}_t|_{\omega_t}^2,$$

as claimed.

We make now more explicit the parallelism with the finite dimensional picture explained in Section §1.1. Recall that Z , $G \subset G^c$ denote, respectively, a Kähler manifold, a compact group acting by holomorphic isometries and its complexification. The distribution D in (1.3.7) plays the role of the tangent space $T_z(G^c \cdot z)$ to the orbit of the complexified group G^c , while the space \mathcal{K} is the analogue of the negatively curved symmetric spaces $G \backslash G^c$. Note that the isomorphism

$$(1.3.13) \quad G \backslash G^c \cong \mathbf{i}\mathfrak{g}$$

given by the exponential map (where \mathfrak{g} denotes the Lie algebra of G) corresponds now to an inclusion

$$\mathcal{K} \subseteq C^\infty(X)/\mathbb{R} \cong \text{Lie } \mathcal{H},$$

where \mathcal{K} is viewed as the space of Kähler potentials of a fixed Kähler form modulo the constant functions. The \mathcal{H} -principal bundle $\mathcal{Y} \rightarrow \mathcal{K}$ and the map τ_J are respectively,

the analogues of the projection $G^c \rightarrow G \setminus G^c$ and the map $G^c \rightarrow G^c \cdot z \subset Z$, provided by the G^c -action on Z . The canonical connection on the principal G -bundle

$$G^c \rightarrow G \setminus G^c$$

is provided by the lift $[e^{i\zeta_t}] \rightarrow ge^{i\zeta_t}$ passing through $g \in G$ and so the smooth solutions to (1.3.12) are the analogues of the geodesics $t \rightarrow [e^{it\zeta}g]$ in (1.3.13).

1.3.4. The Futaki character. In this section we explain an obstruction for the existence of cscK metrics in a Kähler class $[\omega_0]$ due to A. Futaki [29]. This obstruction is the analogue of the character of a complex orbit introduced in §1.1.2 for the finite dimensional case. Together with the discussion in the previous sections, leads to conjecture of the existence of a Hitchin-Kobayashi correspondence for cscK metrics. We keep the notation of the previous sections.

To define the invariant, we proceed as in §1.1.2 and define first a closed 1-form in the space of Kähler metrics. Consider the space of Kähler metrics \mathcal{K} of (X, J) with fixed Kähler class $[\omega]$. Given any $\dot{\phi} \in C^\infty(X)$, we identify it with the constant vector field on \mathcal{K} given by the flow $\omega_{t\dot{\phi}} = \omega + tdd_J^c\dot{\phi}$, for any $\omega \in \mathcal{K}$. Then, the *Mabuchi 1-form* in \mathcal{K} is given by

$$(1.3.14) \quad \sigma_J(\dot{\phi}) = - \int_X \dot{\phi}(S_\omega - \hat{S}) \frac{\omega^n}{n!},$$

where $\dot{\phi}$ is regarded as a tangent to $\omega \in \mathcal{K}$. In order to see the parallelism with (1.1.4), recall from §1.3.3 that given a curve $\omega_t = \omega_{\phi_t}$ in \mathcal{K} we can associate, via its Moser curve of diffeomorphisms f_t , a curve of ω_0 -compatible complex structures $J_t \in \mathcal{J}^i$. Then, doing a change of variable, we can write

$$(1.3.15) \quad \sigma_J(\dot{\phi}_t) = \langle \mu_{\mathcal{H}}(J_t), \dot{\phi}_t \circ f_t \rangle$$

along the curve ω_t , where $\mu_{\mathcal{H}}$ denotes the moment map in the space of ω_0 -compatible complex structures \mathcal{J}^i . From this expression it is straightforward to prove that the 1-form (1.3.14) is closed. We give a proof of this fact for completeness.

Proposition 1.3.2 (T. Mabuchi). *The 1-form σ_J is closed.*

PROOF. Let $\dot{\phi}_1, \dot{\phi}_2 \in C^\infty(X)$, regarded as vector fields in \mathcal{K} . Then it is straightforward that their Lie bracket vanishes: $[\dot{\phi}_1, \dot{\phi}_2] = 0$. Hence, by (1.1.6), we just have to check that $\dot{\phi}_1(\sigma_J(\dot{\phi}_2))$ is symmetric in $j = 1, 2$. Let ω_t be the curve determined by $\dot{\phi}_1$ and consider its horizontal lift f_t starting at any point over ω_0 . Then, using the equations (1.3.15) and (1.3.11), we obtain

$$\frac{d}{dt}_{|t=0} (\sigma_J(\dot{\phi}_2))_{\phi+t\dot{\phi}_1} = \omega_{\mathcal{J}^i}(Y_{\dot{\phi}_2}, \mathbf{J}Y_{\dot{\phi}_1}) - \langle \mu_{\mathcal{H}}(J), \omega(y_{\dot{\phi}_2}, Jy_{\dot{\phi}_1}) \rangle,$$

that is symmetric, as required. \square

Note that the 1-form (1.3.14) is invariant under the right action on \mathcal{K} of the group of automorphisms $\text{Aut}(X, J)$ given by pullback. The infinitesimal action of $\text{Lie Aut}(X, J)$ defines then an invariant $\mathcal{F} = \mathcal{F}_{J, [\omega]}$ of the complex structure and the Kähler class $[\omega]$, due to A. Futaki [29], exactly as in §1.1.2. To give an explicit expression of the *Futaki*

invariant, choose any Kähler metric $\omega \in [\omega]$ and recall that any $\eta \in \text{Lie Aut}(X, J)$ can be written as (see e.g. [44])

$$\eta = y_{\dot{\phi}_1} + Jy_{\dot{\phi}_2} + \beta,$$

where $y_{\dot{\phi}_j}$ is the ω -Hamiltonian field associated to $\dot{\phi}_j$, $j = 1, 2$, and β is the dual of a harmonic 1-form in (X, J, ω) . Then the infinitesimal action of η at the Kähler metric ω is $\eta^l = 2\mathbf{i}\partial_J\bar{\partial}_J\dot{\phi}_2$ and so the Futaki invariant given by $\mathbf{i}\sigma_J(\eta^l) + \sigma_J(J\eta)^l$ can be written as

$$(1.3.16) \quad \mathcal{F}(\eta) = - \int_X (\dot{\phi}_1 + \mathbf{i}\dot{\phi}_2)(S_\omega - \hat{S}) \frac{\omega^n}{n!}$$

for any $\eta \in \text{Lie Aut}(X, J)$. The following result is the analogue of Proposition 1.1.5.

Proposition 1.3.3 (A. Futaki). *The map $\mathcal{F}: \text{Lie Aut}(X, J) \rightarrow \mathbb{C}$ is independent of the Kähler metric $\omega \in [\omega]$. It defines a character of $\text{Lie Aut}(X, J)$ that vanishes if $[\omega]$ contains a cscK metric.*

An elegant proof, similar to the one of Proposition 1.1.5, was provided by J. P Bourguignon [14].

1.3.5. Towards an analytic Hitchin-Kobayashi correspondence. As we have seen in the previous section, not all the complexified orbits of the \mathcal{H} -action in \mathcal{J}^i carry a cscK metric, due to the obstruction given by the Futaki invariant. This leads to conjecture the existence of an analytic Hitchin-Kobayashi correspondence for cscK metrics in a given Kähler class similar to the one explained in §1.2.3 for the HYM equation. The key ingredient in this correspondence will be the integral of the moment map, that in this context was defined by T. Mabuchi [48] and it is usually called the *Mabuchi K-Energy*. In this section we recall the definition of the K-energy and discuss some recent progress in [19] and [18] concerning crucial steps towards an analytic Hitchin-Kobayashi correspondence for the cscK problem.

Let (X, J) be a compact complex manifold. Consider the space \mathcal{K} of Kähler metrics with fixed Kähler class $[\omega]$. Recall from the previous section that (1.3.14) defines a closed 1-form in \mathcal{K} . Since the space \mathcal{K} is contractible, the 1-form σ_J integrates giving the K-energy map

$$\mathcal{M}_\omega : \mathcal{K} \rightarrow \mathbb{R},$$

which satisfies $d\mathcal{M}_\omega = \sigma_J$ and $\mathcal{M}_\omega(\omega) = 0$. This implies that the critical points of the K-energy are the cscK metrics with fixed Kähler class $[\omega]$. Given any curve of metrics $\omega_t = \omega_{\phi_t}$ in \mathcal{K} , we can write explicitly

$$(1.3.17) \quad \mathcal{M}_\omega(\omega_t) = \mathcal{M}_\omega(\omega_0) - \int_0^t \int_X \dot{\phi}_s(S_{\omega_s} - \hat{S}) \frac{\omega_s^n}{n!} \wedge ds,$$

that combined with (1.3.15) shows that \mathcal{M}_ω is convex along geodesics in the space of Kähler metrics. The lack of regularity of the geodesic equation makes very difficult to apply the recipe provided by the proof of Kempf-Ness Theorem, however it has been sufficient for the proof with full generality of the following two fundamental results, due to Chen and Tian [19]. The first one is that there exists a unique solution of the cscK equation in $[\omega_0]$ modulo the action $\text{Aut}(X, J)$. The second one is that the Mabuchi K-energy is bounded from below if there exists a cscK metric in $[\omega_0]$. Other

remarkable progress has been made by X. X. Chen in [18], where a notion of geodesic stability has been defined and proved to be an obstruction for the existence of cscK metrics in $[\omega_0]$. As in Section §1.1, Chen defines maximal weights using the 1-form σ_J and suitable weak solutions to (1.3.12) defined for all time.

1.3.6. Algebraic K-stability. In this section we recall the definition of K -stability for polarized manifolds. This stability notion was introduced by G. Tian in [74] and generalized by S.K. Donaldson in [24].

Let (X, L) be a pair consisting of a projective scheme X and an ample invertible sheaf L over X . The notion of K -stability for such pairs is formulated in terms of flat degenerations of the pair over \mathbb{C} . The central fiber of each degeneration is endowed with a \mathbb{C}^* -action and the stability of the pair, is measured in terms of a suitable weight $-F_1$ associated to this action, the generalized Futaki invariant. When the degeneration considered is smooth it was proved by Donaldson in [24] that $-F_1$ coincides, up to multiplicative constant factor, with the evaluation of the Futaki character (1.3.16) in the generator of the $S^1 \subset \mathbb{C}^*$ action.

We start defining the degenerations of (X, L) that we will consider. In what follows a \mathbb{C}^* -action on such a pair means a \mathbb{C}^* -action on X with a \mathbb{C}^* -linearisation on L .

Definition 1.3.4. A *test configuration* for (X, L) is a pair $(\mathcal{X}, \mathcal{L})$ consisting of a scheme \mathcal{X} and an invertible sheaf \mathcal{L} over \mathcal{X} , together with a \mathbb{C}^* -action on $(\mathcal{X}, \mathcal{L})$ and a flat morphism $\pi: \mathcal{X} \rightarrow \mathbb{C}$, such that

- (1) the invertible sheaf \mathcal{L} is ample,
- (2) $\pi: \mathcal{X} \rightarrow \mathbb{C}$ is \mathbb{C}^* -equivariant, where \mathbb{C}^* acts on \mathbb{C} by multiplication in the standard way,
- (3) the fibre (X_t, L_t) is isomorphic to (X, L) for all $t \in \mathbb{C} \setminus \{0\}$, where $X_t = \pi^{-1}(t)$ and $L_t = \mathcal{L}|_{X_t}$.

Any \mathbb{C}^* -action on (X, L) determines a special type of test configuration, called a *product configuration*, where $\mathcal{X} \cong X \times \mathbb{C}$ with the induced \mathbb{C}^* -action, \mathcal{L} is obtained by pulling the \mathbb{C}^* -linearised sheaf L , and π is the canonical projection. This product configuration is called a *trivial configuration* if the given \mathbb{C}^* -action on (X, L) is trivial, i.e. if it consists of the trivial \mathbb{C}^* -action on X and the \mathbb{C}^* -action defined by scalar multiplication on L .

Let (X, L) be as in Definition 4.1.1. Let us consider a test configuration $(\mathcal{X}, \mathcal{L})$ for (X, L) . As $0 \in \mathbb{C}$ is fixed, the central fibre (X_0, L_0) has an induced \mathbb{C}^* -action. For each integer k we consider the integer $w(k)$ given by the weight of the \mathbb{C}^* -action on the determinant of the cohomology of L_0^k , i.e. on the line (cf. e.g. [37, §§1.2 and 2.1])

$$\det H^*(X_0, L_0^k) := \bigotimes_{p=0}^n \left(\det H^p(X_0, L_0^k) \right)^{(-1)^p}.$$

Then, by the equivariant Riemann–Roch theorem, $w(k)$ is a polynomial of degree at most $n+1$ in k (see [24, §2.1]). Thus the following quotient has an expansion

$$(1.3.18) \quad F(k) := \frac{w(k)}{kP(k)} = F_0 + k^{-1}F_1 + k^{-2}F_2 + O(k^{-3})$$

with rational coefficients F_i , where $P(k)$ is the Hilbert polynomial of \mathcal{O}_{X_0} which has degree n in k . Now, the Futaki invariant of the test configuration $(\mathcal{X}, \mathcal{L}, \cdot)$ is

$$(1.3.19) \quad -F_1.$$

We already have all the ingredients to give the notion of K-stability.

Definition 1.3.5. Let (X, L) be a pair consisting of a projective scheme and an ample invertible sheaf. Then the pair (X, L) is said to be *K-semistable* if

$$(1.3.20) \quad -F_1 \geq 0$$

for all integers $k > 0$ and all test configurations $(\mathcal{X}, \mathcal{L})$ for (X, L^k) . This pair is *K-stable* if the inequality (1.3.20) is strict for all non-trivial $(\mathcal{X}, \mathcal{L}, \mathcal{E})$. It is called α -*K-polystable* if it is K-semistable and the inequality (1.3.20) is always strict unless $(\mathcal{X}, \mathcal{L})$ is a product configuration.

As usual, any pair as in Definition 1.3.5 satisfies the following chain of implications: (X, L) is K-stable \Rightarrow (X, L) is K-polystable \Rightarrow (X, L) is K-semistable. A *K-unstable* triple is a triple which is not K-semistable. We know that any smooth compact polarized manifold in K-polystable mainly due to the work of Donaldson [25], T. Mabuchi [49] and J. Stoppa [67]. Donaldson has proved the converse for the case of toric surfaces [25]. However, as remarked by J. Ross and R. Thomas [63], a recent example in [2] suggests that the algebraic notion of K-stability we have just defined for may not be sufficient to guarantee the existence of a cscK metric and so a stronger notion may be required.

CHAPTER 2

The coupled equations

In this chapter we explain the moment map interpretation of the *coupled equations* (0.0.3), that arises when we consider an extension of the infinite dimensional Lie groups involved in the moment map problems for the HYM and the cscK equation. Symplectic reductions by Lie group extensions have been studied in the literature in various degrees of generality (see [51] and references therein). Previous work includes split group extensions and more general ones, although it seems that the moment map calculations of §2.2, based on Proposition 2.1.1, have not been previously made (cf. [51, §3.2]). We also describe solutions of (0.0.3) in terms of the Kaluza-Klein Theory for G -invariant metrics on E and relate them with the Einstein–Yang–Mills equation.

Throughout this chapter, X is a fixed compact symplectic manifold of dimension $2n$, with symplectic form ω and volume form $\text{vol}_\omega = \frac{\omega^n}{n!}$. We fix a real compact Lie group G with Lie algebra \mathfrak{g} and a smooth principal G -bundle E over X . We also fix a positive definite inner product

$$(\cdot, \cdot) : \mathfrak{g} \times \mathfrak{g} \rightarrow \mathbb{R}$$

on \mathfrak{g} invariant under the adjoint action, that induces a metric on the adjoint bundle $\text{ad } E$. The space of smooth k -forms on X is denoted by Ω^k and for any associated vector bundle F , $\Omega^k(F)$ denote the space of smooth F -valued k -forms on X . Considering the space $\Omega^k(\text{ad } E)$, the metric on $\text{ad } E$ extends to give a pairing

$$(2.0.21) \quad \Omega^p(\text{ad } E) \times \Omega^q(\text{ad } E) \rightarrow \Omega^{p+q}.$$

If X has a fixed almost complex structure compatible with ω the operator

$$(2.0.22) \quad \Lambda : \Omega^{p,q} \rightarrow \Omega^{p-1,q-1}$$

acting on the space of smooth (p, q) -forms is the adjoint of the Lefschetz operator $\Omega^{p-1,q-1} \rightarrow \Omega^{p,q} : \alpha \mapsto \alpha \wedge \omega$, with respect to the metric determined by the symplectic form ω and the given almost complex structure. Its linear extension $\Lambda : \Omega^{p,q}(\text{ad } E) \rightarrow \Omega^{p-1,q-1}(\text{ad } E)$ will be denoted in the same way (we use the same notation as, e.g., in [21]).

2.1. Hamiltonian actions of extended Lie groups

The aim of this section is to describe, under certain assumptions, hamiltonian actions of a Lie group on a symplectic manifold, in terms of a normal Lie subgroup and the quotient Lie group. Later in this thesis we will be interested in the case in which the symplectic manifold and the Lie groups are infinite dimensional.

Suppose that there is an extension of Lie groups

$$(2.1.1) \quad 1 \rightarrow \mathcal{G} \xrightarrow{\iota} \tilde{\mathcal{G}} \xrightarrow{p} \mathcal{H} \rightarrow 1.$$

Such an extension determines another extension of Lie algebras

$$(2.1.2) \quad 0 \rightarrow \text{Lie } \mathcal{G} \xrightarrow{\iota} \text{Lie } \tilde{\mathcal{G}} \xrightarrow{p} \text{Lie } \mathcal{H} \rightarrow 0,$$

where the use of the same symbols ι and p should lead to no confusion. Let \mathcal{A} be a manifold with an action of the ‘extended’ Lie group $\tilde{\mathcal{G}}$. For each $\zeta \in \text{Lie } \tilde{\mathcal{G}}$, let Y_ζ be the vector field on \mathcal{A} generated by the infinitesimal action of ζ .

Note that the short exact sequence (2.1.2) does not generally split as a sequence of Lie algebras, but it always does as a short exact sequence of vector spaces. Let $W \subset \text{Hom}(\text{Lie } \tilde{\mathcal{G}}, \text{Lie } \mathcal{G})$ be the affine space of vector space splittings. Since $\mathcal{G} \subset \tilde{\mathcal{G}}$ is a normal subgroup, there is a well-defined $\tilde{\mathcal{G}}$ -action on W , given by

$$g \cdot \theta := \text{Ad}(g) \circ \theta \circ \text{Ad}(g^{-1}), \text{ for } g \in \tilde{\mathcal{G}}, \theta \in W.$$

Let $\mathcal{W} \subset C^\infty(\mathcal{A}, W)$ be the space of $\tilde{\mathcal{G}}$ -equivariant smooth maps $\theta: \mathcal{A} \rightarrow W$. The case considered in this thesis (see §2.2) is an example where the space \mathcal{W} is non-empty. Observe that if \mathcal{A} is a point, then \mathcal{W} is the space of $\tilde{\mathcal{G}}$ -equivariant splittings of (2.1.2) and any such equivariant splitting determines an isomorphism $\text{Lie } \tilde{\mathcal{G}} \cong \text{Lie } \mathcal{G} \rtimes \text{Lie } \mathcal{H}$, so in this case the condition that $\mathcal{W} \neq \emptyset$ is very strong.

Let $\omega_{\mathcal{A}}$ be a symplectic form on \mathcal{A} such that the $\tilde{\mathcal{G}}$ -action on \mathcal{A} is symplectic. We want to characterise hamiltonian $\tilde{\mathcal{G}}$ -actions on \mathcal{A} in terms of \mathcal{G} and \mathcal{H} under the assumption that there exists some $\theta \in \mathcal{W}$. Suppose that the $\tilde{\mathcal{G}}$ -action is Hamiltonian with $\tilde{\mathcal{G}}$ -equivariant moment map $\mu_{\tilde{\mathcal{G}}}: \mathcal{A} \rightarrow (\text{Lie } \tilde{\mathcal{G}})^*$. Using $\theta \in \mathcal{W}$ we will break this map into pieces corresponding to the Lie algebras $\text{Lie } \mathcal{G}$ and $\text{Lie } \mathcal{H}$. Note first that the isomorphism

$$W \rightarrow \text{Hom}(\text{Lie } \mathcal{H}, \text{Lie } \tilde{\mathcal{G}}): \theta \mapsto \theta^\perp,$$

where $\theta^\perp: \text{Lie } \mathcal{H} \rightarrow \text{Lie } \tilde{\mathcal{G}}$ is uniquely defined by the equation

$$(2.1.3) \quad \text{Id} - \iota \circ \theta = \theta^\perp \circ p,$$

is $\tilde{\mathcal{G}}$ -equivariant with respect to the action in $\text{Hom}(\text{Lie } \mathcal{H}, \text{Lie } \tilde{\mathcal{G}})$ given by $g \cdot \theta^\perp = \text{Ad}(g) \circ \theta^\perp \circ \text{Ad}(p(g^{-1}))$ for all $g \in \tilde{\mathcal{G}}$. Here we are using the $\tilde{\mathcal{G}}$ -action in $\text{Lie } \mathcal{H}$ induced by the surjective homomorphism p in the exact sequence (2.1.1). Thus, given now $\theta \in \mathcal{W}$, the map

$$\theta^\perp: \mathcal{A} \rightarrow \text{Hom}(\text{Lie } \mathcal{H}, \text{Lie } \tilde{\mathcal{G}})$$

is $\tilde{\mathcal{G}}$ -equivariant and we can break the moment map into

$$(2.1.4) \quad \langle \mu_{\tilde{\mathcal{G}}}, \zeta \rangle = \langle \mu_{\tilde{\mathcal{G}}}, \iota \theta \zeta \rangle + \langle \mu_{\tilde{\mathcal{G}}}, \theta^\perp p(\zeta) \rangle, \text{ for all } \zeta \in \text{Lie } \tilde{\mathcal{G}},$$

where the summands in the right hand side define a pair of $\tilde{\mathcal{G}}$ -equivariant maps $\mu_{\mathcal{G}}: \mathcal{A} \rightarrow (\text{Lie } \mathcal{G})^*$ given by

$$\langle \mu_{\mathcal{G}}, \zeta \rangle := \langle \mu_{\tilde{\mathcal{G}}}, \iota \zeta \rangle$$

for all $\zeta \in \text{Lie } \mathcal{G}$, and $\sigma_\theta: \mathcal{A} \rightarrow (\text{Lie } \mathcal{H})^*$

$$(2.1.5) \quad \langle \sigma_\theta, \eta \rangle := \langle \mu_{\tilde{\mathcal{G}}}, \theta^\perp \eta \rangle, \text{ for all } \eta \in \text{Lie } \mathcal{H}.$$

Note that since \mathcal{G} is a normal subgroup of $\tilde{\mathcal{G}}$ we can require the map $\mu_{\mathcal{G}}$ to be $\tilde{\mathcal{G}}$ -equivariant. It is now straightforward from the moment map condition for $\mu_{\tilde{\mathcal{G}}}$ to check that $\mu_{\mathcal{G}}$ is a \mathcal{G} -equivariant moment map for the \mathcal{G} -action on \mathcal{A} , i.e. $d\langle \mu_{\mathcal{G}}, \zeta \rangle = i_{Y_{\zeta}}\omega_{\mathcal{A}}$ for all $\zeta \in \text{Lie } \mathcal{G}$. We will see in the next proposition that σ_{θ} satisfies a similar infinitesimal condition that provides our characterization of Hamiltonian $\tilde{\mathcal{G}}$ -action, for what we first introduce some notation. Given a smooth map $\zeta: \mathcal{A} \rightarrow \text{Lie } \tilde{\mathcal{G}}$, let us denote by Y_{ζ} the vector field on \mathcal{A} given by

$$(Y_{\zeta})_A := \frac{d}{dt} \Big|_{t=0} \exp(t\zeta_A) \cdot A, \text{ for all } A \in \mathcal{A}.$$

In particular, in the notation (2.1.3), any smooth $\theta: \mathcal{A} \rightarrow W$ defines a map

$$Y_{\theta^{\perp}}: \text{Lie } \mathcal{H} \rightarrow \Omega^0(T\mathcal{A}): \eta \mapsto Y_{\theta^{\perp}\eta}.$$

Note also that $d\theta$ is a $\tilde{\mathcal{G}}$ -invariant $\text{Hom}(\text{Lie } \mathcal{H}, \text{Lie } \mathcal{G})$ -valued 1-form on \mathcal{A} , as $\theta \in \mathcal{W}$ and the affine space \mathcal{W} is modelled on the space of $\tilde{\mathcal{G}}$ -equivariant smooth maps $\tau: \mathcal{A} \rightarrow \text{Hom}(\text{Lie } \mathcal{H}, \text{Lie } \mathcal{G})$.

Proposition 2.1.1. *Suppose that there exists some $\theta \in \mathcal{W}$. Then the $\tilde{\mathcal{G}}$ -action on \mathcal{A} is hamiltonian if and only if so is the action of $\mathcal{G} \subset \tilde{\mathcal{G}}$ on \mathcal{A} , with a $\tilde{\mathcal{G}}$ -equivariant moment map $\mu_{\mathcal{G}}: \mathcal{A} \rightarrow (\text{Lie } \mathcal{G})^*$, and there exists a smooth $\tilde{\mathcal{G}}$ -equivariant map $\sigma_{\theta}: \mathcal{A} \rightarrow (\text{Lie } \mathcal{H})^*$ satisfying*

$$(2.1.6) \quad i_{Y_{\theta^{\perp}\eta}}\omega_{\mathcal{A}} = \langle \mu_{\mathcal{G}}, (d\theta)\eta \rangle + d\langle \sigma_{\theta}, \eta \rangle, \text{ for all } \eta \in \text{Lie } \mathcal{H}.$$

In this case, a $\tilde{\mathcal{G}}$ -equivariant moment map $\mu_{\tilde{\mathcal{G}}}: \mathcal{A} \rightarrow (\text{Lie } \tilde{\mathcal{G}})^*$ is given by

$$(2.1.7) \quad \langle \mu_{\tilde{\mathcal{G}}}, \zeta \rangle = \langle \mu_{\mathcal{G}}, \theta\zeta \rangle + \langle \sigma_{\theta}, p(\zeta) \rangle, \text{ for all } \zeta \in \text{Lie } \tilde{\mathcal{G}}.$$

PROOF. To prove the “only if” part it remains to check that (2.1.6) holds. This follows by definition, by taking differentials in (2.1.4) and using that

$$\begin{aligned} d\langle \mu_{\mathcal{G}}, \theta\zeta \rangle &= \langle d\mu_{\mathcal{G}}, \theta\zeta \rangle + \langle \mu_{\mathcal{G}}, (d\theta)\zeta \rangle \text{ and} \\ i_{Y_{\zeta}}\omega &= i_{Y_{\theta\zeta}}\omega + i_{Y_{\theta^{\perp}\zeta}}\omega, \text{ with } \eta := p(\zeta), \end{aligned}$$

where the first equation is obtained applying the chain rule, and the second one holds because $\zeta = \theta\zeta + \theta^{\perp}\eta$ and Y_{ζ} is linear in ζ . The “if” part is straightforward from the statement and is left to the reader. \square

To see the dependence of σ_{θ} on the choice of $\theta \in \mathcal{W}$, recall that \mathcal{W} is an affine space modelled on the vector space of $\tilde{\mathcal{G}}$ -equivariant smooth maps $\tau: \mathcal{A} \rightarrow \text{Hom}(\text{Lie } \mathcal{H}, \text{Lie } \mathcal{G})$, so the elements of \mathcal{W} are all of the form $\theta + p^*\tau$. Now, if θ and σ_{θ} satisfy the conditions of Proposition 2.1.1, then the same conditions is obviously satisfied by $\sigma_{\theta+p^*\tau}$ with respect to $\theta + p^*\tau$, where

$$\langle \sigma_{\theta+p^*\tau}, \eta \rangle = \langle \sigma_{\theta}, \eta \rangle - \langle \mu_{\mathcal{G}}, \tau\eta \rangle, \text{ for all } \eta \in \text{Lie } \mathcal{H}.$$

Condition (2.1.6) for σ_{θ} in proposition 2.1.1 generalizes the usual infinitesimal condition $i_{Y_{\eta}}\omega_{\mathcal{A}} = d\langle \mu_{\mathcal{H}}, \eta \rangle$ ($\eta \in \mathcal{H}$) for moment maps $\mu_{\mathcal{H}}$ for the induced \mathcal{H} -action on \mathcal{A} when the Lie group extension (2.1.1) splits. To see this, suppose that there exists such a splitting. It determines an \mathcal{H} -action on \mathcal{A} and a $\tilde{\mathcal{G}}$ -invariant $\theta_0 \in \mathcal{W}$, so for all $\eta \in \mathcal{H}$,

$Y_{\theta_0 \pm \eta}$ is the vector field induced by the infinitesimal action of η . Then the constant map $\theta: \mathcal{A} \rightarrow W: A \mapsto \theta_0$ is in \mathcal{W} and $d\theta = 0$, so (2.1.6) is precisely the condition that σ_θ is a moment map for the \mathcal{H} -action on \mathcal{A} . Thus, for a fixed Lie group splitting of (2.1.1), the $\tilde{\mathcal{G}}$ -action on \mathcal{A} is Hamiltonian if and only if so are the induced actions of \mathcal{G} and \mathcal{H} , with $\tilde{\mathcal{G}}$ -equivariant moment maps $\mu_{\mathcal{G}}$ and σ_θ , respectively.

2.2. The extended gauge group action on the space of connections

Let X be a compact symplectic manifold of dimension $2n$, with symplectic form ω . We fix a compact Lie group G and a (smooth) principal G -bundle E on X , with projection map $\pi: E \rightarrow X$. Let \mathcal{H} be the group of Hamiltonian symplectomorphisms of X , as defined in §1.3.2. An *automorphism* of E is a G -equivariant diffeomorphism $g: E \rightarrow E$. Any such automorphism covers a unique diffeomorphism $h: X \rightarrow X$, i.e. a unique h such that $\pi \circ g = h \circ \pi$. The (Hamiltonian) *extended gauge group* $\tilde{\mathcal{G}}$ of E is the group of automorphisms of E which cover elements of \mathcal{H} . Then the gauge group of E , already mentioned in §1.2.2, is the normal subgroup $\mathcal{G} \subset \tilde{\mathcal{G}}$ of automorphisms covering the identity. Furthermore, there is a canonical short exact sequence of Lie groups

$$(2.2.1) \quad 1 \rightarrow \mathcal{G} \xrightarrow{\iota} \tilde{\mathcal{G}} \xrightarrow{p} \mathcal{H} \rightarrow 1.$$

Here ι is the inclusion map and p is the map that assigns to each automorphism the Hamiltonian symplectomorphisms that it covers.

The only non-trivial fact that needs to be proved to show that (2.2.1) is exact is that p is surjective. Choose a connection A on E . Given $h \in \mathcal{H}$, by definition there exists a Hamiltonian isotopy $[0, 1] \times X \rightarrow X: (t, x) \mapsto h_t(x)$ from $h_0 = \text{Id}$ to $h_1 = h$, which is the flow of a smooth family of vector fields $\eta_t \in \text{Lie } \mathcal{H}$, i.e. with $dh_t/dt = \eta_t \circ h_t^{-1}$ (see e.g. [55, §3.2]). Let $\zeta_t \in \text{Lie } \tilde{\mathcal{G}}$ be the horizontal lift to E of η_t given by the connection A . The vector fields ζ_t are G -invariant so its time-dependent flow g_t exists for all $t \in [0, 1]$ and the $g_t: E \rightarrow E$ are G -equivariant. Since ζ_t is a lift of η_t to E , its flow g_t covers h_t , so in particular $g_t \in \tilde{\mathcal{G}}$ for all t and in particular $g_1 \in \tilde{\mathcal{G}}$ covers $h = h_1$. Thus p is surjective and so (2.2.1) is exact, as claimed.

Remark 2.2.1. Note that the existence of the short exact sequence (2.2.1) holds even when the structure group or even the base are non-compact. The crucial fact is that \mathcal{H} lies in the connected component of the identity of $\text{Diff } X$. For generalizations of the exact sequence (2.2.1) see [1].

There is an action of the extended gauge group on the space \mathcal{A} of connections on E . To define this action, we view the elements of \mathcal{A} as G -equivariant splittings $A: TE \rightarrow VE$ of the short exact sequence

$$(2.2.2) \quad 0 \rightarrow VE \longrightarrow TE \longrightarrow \pi^* TX \rightarrow 0,$$

where $VE = \ker d\pi$ is the vertical bundle. Using the action of $g \in \tilde{\mathcal{G}}$ on TE , the $\tilde{\mathcal{G}}$ -action on \mathcal{A} is $g \cdot A := g \circ A \circ g^{-1}$. Any such splitting A induces a vector space splitting of the Atiyah short exact sequence

$$(2.2.3) \quad 0 \rightarrow \text{Lie } \mathcal{G} \xrightarrow{\iota} \text{Lie}(\text{Aut } E) \xrightarrow{p} \text{Lie}(\text{Diff } X) \rightarrow 0$$

(cf. [6, equation (3.4)]), where $\text{Lie}(\text{Diff } X)$ is the Lie algebra of vector fields on X and $\text{Lie}(\text{Aut } E)$ is the Lie algebra of G -invariant vector fields on E . This splitting is given by maps

$$\theta_A: \text{Lie}(\text{Aut } E) \rightarrow \text{Lie } \mathcal{G}, \quad \theta_A^\perp: \text{Lie}(\text{Diff } X) \rightarrow \text{Lie}(\text{Aut } E)$$

such that $\iota \circ \theta_A + \theta_A^\perp \circ p = \text{Id}$, where θ_A^\perp is the horizontal lift of vector fields on X to vector fields on E given by A . This splitting restricts to another one of the exact sequence

$$(2.2.4) \quad 0 \rightarrow \text{Lie } \mathcal{G} \xrightarrow{\iota} \text{Lie } \tilde{\mathcal{G}} \xrightarrow{p} \text{Lie } \mathcal{H} \rightarrow 0$$

induced by (2.2.1), because $\text{Lie } \tilde{\mathcal{G}} = p^{-1}(\text{Lie } \mathcal{H})$. In the notation of §2.1, it is easy to see that the map

$$(2.2.5) \quad \theta: \mathcal{A} \rightarrow W: A \mapsto \theta_A$$

is an element of \mathcal{W} . It is also clear that the $\tilde{\mathcal{G}}$ -action on \mathcal{A} is symplectic, for the symplectic form (1.2.1). The methods of §2.1 apply here to provide a moment map. Let \hat{c} be the real constant

$$(2.2.6) \quad \hat{c} := \frac{\int_X (F_A \wedge F_A) \wedge \omega^{n-2}}{\int_X \text{vol}_\omega} = \frac{\langle c(E) \cup [\omega]^{n-2}, [X] \rangle}{\text{Vol}([\omega])}$$

where $c(E) \in H^4(X)$ is the Chern-Weyl class associated to the biinvariant bilinear form (\cdot, \cdot) in \mathfrak{g} , that depends only on the topology of the bundle (see Theorem 1.1, p. 293 in [42]). Then,

Proposition 2.2.2. *The $\tilde{\mathcal{G}}$ -action on \mathcal{A} is hamiltonian, with $\tilde{\mathcal{G}}$ -equivariant moment map $\mu_{\tilde{\mathcal{G}}}: \mathcal{A} \rightarrow (\text{Lie } \tilde{\mathcal{G}})^*$ given by*

$$(2.2.7) \quad \langle \mu_{\tilde{\mathcal{G}}}, \zeta \rangle = \langle \mu_{\mathcal{G}}, \theta \zeta \rangle + \langle \sigma, p(\zeta) \rangle, \text{ for all } \zeta \in \text{Lie } \tilde{\mathcal{G}},$$

where $\mu_{\mathcal{G}}: \mathcal{A} \rightarrow (\text{Lie } \mathcal{G})^*$ and $\sigma: \mathcal{A} \rightarrow (\text{Lie } \mathcal{H})^*$ are

$$(2.2.8) \quad \begin{aligned} \langle \mu_{\mathcal{G}}, \theta \zeta \rangle(A) &= \int_X (\theta_A \zeta \wedge F_A) \wedge \omega^{n-1} \\ \langle \sigma, \eta \rangle(A) &= -\frac{(n-1)}{2} \int_X \eta((F_A \wedge F_A) \wedge \omega^{n-2} - \hat{c} \text{vol}_\omega) \end{aligned}$$

for any $\eta \in \text{Lie } \mathcal{H} \cong C^\infty(X)/\mathbb{R}$.

PROOF. By Proposition 2.1.1 this will follow from the facts that $\mu_{\mathcal{G}}$ and σ are $\tilde{\mathcal{G}}$ -equivariant and that the map σ defined by (2.2.8) satisfies (2.1.6). The $\tilde{\mathcal{G}}$ -equivariance is immediate from the formulae (2.2.8) and the Change of Variable Theorem. To prove the last fact, we first compute the terms involved in (2.1.6).

Let $\zeta \in \text{Lie}(\text{Aut } E)$, $A \in \mathcal{A}$. By the Leibniz rule, for all $v \in \Omega^0(TE)$,

$$\left. \frac{d}{dt} \right|_{t=0} (e^{t\zeta} \circ A \circ e^{-t\zeta}(v)) = \theta_A[\zeta, v] - [\zeta, \theta_A(v)] = \theta_A[\zeta, v - \theta_A(v)],$$

where in the second equality we have used the fact that ζ covers a vector field on X , so that the vector field $[\zeta, \theta_A(v)]$ is vertical. It is easy to see that this expression is tensorial in v , so at each point of E it only depends on its projection $p(v)$. Thus for

any $\zeta \in \text{Lie } \tilde{\mathcal{G}}$ the infinitesimal action $Y_\zeta(A)$ on \mathcal{A} regarded as an element of $\Omega^1(\text{ad } E)$ is given by

$$\begin{aligned} Y_\zeta(A)(y) &= \theta_A[\zeta, \theta_A^\perp(y)] \\ (2.2.9) \quad &= [\theta_A\zeta, \theta_A^\perp(y)] + \theta_A[\theta_A^\perp(p\zeta), \theta_A^\perp(y)] \\ &= -d_A(\theta_A\zeta) - i_{p(\zeta)}F_A, \end{aligned}$$

for any $y \in \Omega^0(TX)$ (cf. with the equation before (4.2) and the equation after (3.4) in [6]. Note that we are using a different sign for the curvature).

Note that (2.2.9) also applies for maps $\zeta: \mathcal{A} \rightarrow \text{Lie}(\text{Aut } E)$. In particular,

$$(2.2.10) \quad Y_{\theta_A^\perp\eta}(A) = -i_\eta F_A, \text{ for } \eta \in \text{Lie } \mathcal{H}.$$

The $\text{Hom}(\text{Lie } \mathcal{H}, \text{Lie } \tilde{\mathcal{G}})$ -valued 1-form $d\theta$ on \mathcal{A} is given by

$$(2.2.11) \quad i_a(d\theta) = a: \text{Lie } \mathcal{H} \rightarrow \text{Lie } \tilde{\mathcal{G}}: \eta \mapsto i_\eta a,$$

for $A \in \mathcal{A}$ and $a \in T_A\mathcal{A} = \Omega^1(\text{ad } E)$. Hence, with the same notation,

$$(2.2.12) \quad ((d\theta)\eta)(a) = a(\eta), \text{ for } \eta \in \text{Lie } \mathcal{H}.$$

For the last term of (2.1.6), let $A \in \mathcal{A}$, $a \in T_A\mathcal{A}$, $\zeta \in \text{Lie } \tilde{\mathcal{G}}$. Let $\eta = \eta_\phi \in \text{Lie } \mathcal{H}$, where $\phi \in C^\infty(X)/\mathbb{R}$. Using the path $A_t = A + ta$, we obtain

$$\begin{aligned} d\langle \sigma, \eta \rangle(a) &= -\frac{n-1}{2} \frac{d}{dt} \Big|_{t=0} \int_X \phi(F_{A_t} \wedge F_{A_t}) \wedge \omega^{n-2} - \hat{c} \text{vol}_\omega \\ (2.2.13) \quad &= -(n-1) \int_X \phi(d_A a \wedge F_A) \wedge \omega^{n-2} \\ &= -(n-1) \int_X \phi d(a \wedge F_A) \wedge \omega^{n-2} \\ &= (n-1) \int_X i_\eta \omega \wedge (a \wedge F_A) \wedge \omega^{n-2}. \end{aligned}$$

Here we have used $dF_{A_t}/dt = d_A a$ for $t = 0$ in the second equality, the Bianchi identity $d_A F_A = 0$ in the third equality, and we have integrated by parts and used the formula $d\phi = i_\eta \omega$ in the last equality. To compute the integral in the last equality, note that $(a \wedge F_A) \wedge \omega^{n-1} = 0$, as $\dim X = 2n$, so $i_\eta((a \wedge F_A) \wedge \omega^{n-1})$ equals

$$(a(\eta) \wedge F_A) \wedge \omega^{n-1} - (a \wedge i_\eta F_A) \wedge \omega^{n-1} - (n-1)(a \wedge F_A) \wedge i_\eta \omega \wedge \omega^{n-2} = 0.$$

Combined with (2.2.13), we thus obtain equation (2.1.6):

$$\begin{aligned} d\langle \sigma, \eta \rangle(a) &= \int_X (a \wedge i_\eta F_A) \wedge \omega^{n-1} - (a(\eta) \wedge F_A) \wedge \omega^{n-1} \\ &= \int_X (Y_{\theta_A^\perp\eta}(A) \wedge a) \wedge \omega^{n-1} - (((d\theta)\eta)(a) \wedge F_A) \wedge \omega^{n-1} \\ &= i_{Y_{\theta_A^\perp\eta}} \omega_{\mathcal{A}}(a) - \langle \mu_{\tilde{\mathcal{G}}}, (d\theta)\eta \rangle(a). \end{aligned} \quad \square$$

To finish this section we identify the image of the moment map $\mu_{\tilde{\mathcal{G}}} : \mathcal{A} \rightarrow (\text{Lie } \tilde{\mathcal{G}})^*$ in the above proposition with elements in $\text{Lie } \tilde{\mathcal{G}}$ as we did in §1.2.2 and §1.3.2 for the groups \mathcal{G} and \mathcal{H} . Though we do not have now a canonical $\tilde{\mathcal{G}}$ -invariant pairing in $\text{Lie } \tilde{\mathcal{G}}$,

we can use the $\tilde{\mathcal{G}}$ -equivariant map defined in (2.2.5), $\theta: \mathcal{A} \rightarrow \text{Hom}(\text{Lie } \mathcal{H}, \text{Lie } \mathcal{G})$, to construct an \mathcal{A} -dependent pairing. Let us identify $\text{Lie } \mathcal{H} \cong C^\infty(X)/\mathbb{R}$ with $C_0^\infty(X)$, the space of functions with zero integral over X with respect to vol_ω . Then, the L^2 pairing on functions on X provides an \mathcal{H} -invariant pairing

$$(2.2.14) \quad \langle \eta_1, \eta_2 \rangle_{\mathcal{H}} = \int_X \phi_1 \phi_2 \text{vol}_\omega,$$

where $\eta_j \in \text{Lie } \mathcal{H}$ and $\phi_j \in C_0^\infty(X)$ is the Hamiltonian function determined by η_j , $j = 1, 2$. Then, we obtain a map

$$\langle \cdot, \cdot \rangle_{\tilde{\mathcal{G}}} : \mathcal{A} \rightarrow S^2 \text{Lie } \tilde{\mathcal{G}}^*$$

that to each $A \in \mathcal{A}$ assigns a symmetric pairing on $\text{Lie } \tilde{\mathcal{G}}$ given by

$$(2.2.15) \quad \langle \zeta_1, \zeta_2 \rangle_{\tilde{\mathcal{G}}|A} = \langle p(\zeta_1), p(\zeta_2) \rangle_{\mathcal{H}} + \langle \theta_A \zeta_1, \theta_A \zeta_2 \rangle_{\mathcal{G}}, \text{ for } \zeta_1, \zeta_2 \in \text{Lie } \tilde{\mathcal{G}}.$$

Since both, θ and $\langle \cdot, \cdot \rangle_{\mathcal{G}}$ are $\tilde{\mathcal{G}}$ -equivariant (recall that \mathcal{G} is normal in $\tilde{\mathcal{G}}$), the map $\langle \cdot, \cdot \rangle_{\tilde{\mathcal{G}}}$ is $\tilde{\mathcal{G}}$ -equivariant. Given $A \in \mathcal{A}$ let

$$\phi_A = \frac{1}{4} \Lambda_\omega^2(F_A \wedge F_A) \in C^\infty(X),$$

and consider the element of the Lie algebra of $\tilde{\mathcal{G}}$ given by

$$(2.2.16) \quad \zeta_A = \Lambda_\omega F_A - \theta_A^\perp(\eta_A) \in \text{Lie } \tilde{\mathcal{G}},$$

where $\Lambda_\omega F_A \in \Omega^0(\text{ad } E)$ is regarded as a vertical G -invariant vector field in E and $\eta_A \in \text{Lie } \mathcal{H}$ is the Hamiltonian vector field associated to ϕ_A . Then, given $\zeta \in \text{Lie } \tilde{\mathcal{G}}$ covering $\eta_\phi \in \text{Lie } \mathcal{H}$ we have

$$(2.2.17) \quad \begin{aligned} \langle \mu_{\tilde{\mathcal{G}}}(A), \zeta \rangle &= \int_X (\theta_A \zeta \wedge F_A) \wedge \omega^{n-1} - \frac{(n-1)}{2} \int_X \phi((F_A \wedge F_A) \wedge \omega^{n-2} - \hat{c} \text{vol}_\omega) \\ &= \frac{(n-1)!}{4} \left(\int_X (\theta_A \zeta \wedge 4\Lambda_\omega F_A) - \phi(\Lambda_\omega^2(F_A \wedge F_A) - \frac{2\hat{c}}{(n-2)!}) \cdot \text{vol}_\omega \right) \\ &= \frac{(n-1)!}{4} \left(\langle \theta_A \zeta \wedge 4\Lambda_\omega F_A \rangle_{\mathcal{G}} + \langle \phi, -(\Lambda_\omega^2(F_A \wedge F_A) - \frac{2\hat{c}}{(n-2)!}) \rangle_{\mathcal{H}} \right) \\ &= (n-1)! \langle \zeta, \zeta_A \rangle_{\tilde{\mathcal{G}}|A}, \end{aligned}$$

and so the moment map $\mu_{\tilde{\mathcal{G}}}$ satisfies

$$(2.2.18) \quad \mu_{\tilde{\mathcal{G}}}(A) = (n-1)! \langle \zeta_A, \cdot \rangle_{\tilde{\mathcal{G}}|A} \in \text{Lie } \tilde{\mathcal{G}}^*.$$

For the previous computation in (2.2.17) we have used the identities

$$(2.2.19) \quad \begin{aligned} (n-1)! \Lambda_\omega F_A \text{vol}_\omega &= F_A \wedge \omega^{n-1} \\ (n-2)! \Lambda_\omega^2(F_A \wedge F_A) \text{vol}_\omega &= 2(F_A \wedge F_A) \wedge \omega^{n-2}. \end{aligned}$$

The $\tilde{\mathcal{G}}$ -equivariance of the r.h.s. of (2.2.18) is justified by the equality $\zeta(gA) = g \cdot \zeta_A$, for any $A \in \mathcal{A}$, $g \in \tilde{\mathcal{G}}$.

2.3. The coupled equations as a moment map condition

As in the previous section, let X be a compact symplectic manifold of dimension $2n$, with symplectic form ω , G be a compact Lie group and E be a smooth principal G -bundle on X . Let \mathcal{J} be the space of almost complex structures compatible with ω and \mathcal{A} be the space of connections on E . Using the symplectic structures on \mathcal{J} and \mathcal{A} induced by ω (see (1.3.4) and (1.2.1)), we can define a symplectic form on the product $\mathcal{J} \times \mathcal{A}$, for each pair of non-zero real constants $\alpha = (\alpha_0, \alpha_1)$, given by

$$(2.3.1) \quad \omega_\alpha = \alpha_0 \cdot \omega_{\mathcal{J}} + \frac{4\alpha_1}{(n-1)!} \cdot \omega_{\mathcal{A}}.$$

The extended gauge group $\tilde{\mathcal{G}}$ has a canonical action on $\mathcal{J} \times \mathcal{A}$ and this action is symplectic for any ω_α . In the notation of §2.2, this action is given by

$$g \cdot (J, A) = (p(g) \cdot J, g \cdot A), \text{ for } g \in \tilde{\mathcal{G}} \text{ and } (J, A) \in \mathcal{J} \times \mathcal{A},$$

where p is the map in the short exact sequence (2.2.1). Let $z \in \mathfrak{g}$ be an element in the centre of \mathfrak{g} . We define the real constant

$$(2.3.2) \quad c_z = \alpha_0 \hat{S} + \alpha_1 \left(\frac{2\hat{c}}{(n-2)!} - 4|z|^2 \right),$$

where \hat{S} is as in (1.3.3) and \hat{c} is given by (2.2.6). Combining the moment map for the \mathcal{H} -action on \mathcal{J} (see §1.3.2), Proposition 2.2.2 and equation (2.2.17) we obtain the following.

Proposition 2.3.1. *The $\tilde{\mathcal{G}}$ -action on $\mathcal{J} \times \mathcal{A}$ is hamiltonian with respect to (2.3.1), with $\tilde{\mathcal{G}}$ -equivariant moment map $\mu_\alpha: \mathcal{J} \times \mathcal{A} \rightarrow (\text{Lie } \tilde{\mathcal{G}})^*$ given by*

$$(2.3.3) \quad \begin{aligned} \langle \mu_\alpha(J, A), \zeta \rangle &= 4\alpha_1 \int_X (\theta_A \zeta \wedge (\Lambda_\omega F_A - z)) \cdot \text{vol}_\omega \\ &\quad - \int_X \phi(\alpha_0 S_J + \alpha_1 \Lambda_\omega^2(F_A \wedge F_A) - 4\alpha_1 (\Lambda_\omega F_A \wedge z) - c_z) \text{vol}_\omega \end{aligned}$$

for any $\zeta \in \text{Lie } \tilde{\mathcal{G}}$ covering $\eta_\phi \in \text{Lie } \mathcal{H}$.

PROOF. Simply note that

$$\int_X ((\theta_A \zeta - \phi \Lambda_\omega F_A) \wedge z) \text{vol}_\omega$$

is constant over \mathcal{A} . □

The space $\mathcal{J} \times \mathcal{A}$ has a complex structure preserved by the $\tilde{\mathcal{G}}$ -action given by

$$(2.3.4) \quad \mathbf{I}_{|(J,A)}(\dot{J}, a) = (J\dot{J}, -a(J\cdot)); \quad (\dot{J}, a) \in T_J \mathcal{J} \times T_A \mathcal{A}.$$

The projection $\mathcal{J} \times \mathcal{A} \rightarrow \mathcal{J}$ becomes now a holomorphic submersion with respect to the complex structure in \mathcal{J} and it turns out that, for α_0, α_1 positive, this complex structure is compatible with the family of symplectic structures (2.3.1). This endows $\mathcal{J} \times \mathcal{A}$ with a structure of Kähler fibration, i.e. a holomorphic submersion with a Kähler metric ω_α that is the sum of the closed form $\frac{4\alpha}{(n-1)!} \cdot \omega_{\mathcal{A}}$ that restricts to a Kähler metric in the fibers and the pull-back of a Kähler metric in the base $\alpha_0 \cdot \omega_{\mathcal{J}}$.

The formal integrability of the almost complex structure \mathbf{I} is a priori not obvious and we give a proof in the following proposition. By formal integrability we mean, as in [23], that the associated Nijenhuis tensor vanishes.

Proposition 2.3.2. *The almost complex structure \mathbf{I} is integrable.*

PROOF. For this we regard $\mathcal{J} \times \mathcal{A}$ as an almost complex fibration over \mathcal{J} with respect to the complex structure introduced in §1.3.2. The complex structure on the base and the one on each fibre are integrable, and hence the integrability condition of \mathbf{I} reduces to evaluate the Nijenhuis tensor $N_{\mathbf{I}}$ on a pair of vectors $\dot{J} \in T_J\mathcal{J}$, $a \in T_A\mathcal{A}$ for each (J, A) in the product $\mathcal{J} \times \mathcal{A}$. Since \mathcal{A} is an affine space we can regard a as a constant vector field on \mathcal{A} and so on the product $\mathcal{J} \times \mathcal{A}$. We can also extend \dot{J} to a vector field on \mathcal{J} taking

$$\dot{J}_{|J'} = \frac{1}{2}(J\dot{J}J' - J'\dot{J}J).$$

Using this extension it make sense to write

$$\begin{aligned} N_{\mathbf{I}}(\dot{J}, a) &= [\mathbf{I}\dot{J}, \mathbf{I}a] - \mathbf{I}[\mathbf{I}\dot{J}, a] - \mathbf{I}[\dot{J}, \mathbf{I}a] - [\dot{J}, a] \\ &= [\mathbf{I}\dot{J}, \mathbf{I}a] - \mathbf{I}[\dot{J}, \mathbf{I}a], \end{aligned}$$

where the brackets denote the Lie bracket between vector fields in $\mathcal{J} \times \mathcal{A}$. We have used that $[\mathbf{I}\dot{J}, a] = [\dot{J}, a] = 0$, since the flow of a is the identity on \mathcal{J} . Given any vector field \dot{J} on \mathcal{J} , regarded as a vector field on $\mathcal{J} \times \mathcal{A}$, we denote its flow by $J_t^{\dot{J}}$. Note that $J_t^{\dot{J}}$ induces the identity on \mathcal{A} and so

$$\begin{aligned} N_{\mathbf{I}}(\dot{J}, a) &= \frac{d}{dt}|_{t=0} \mathbf{I}a_{|J_t^{\dot{J}}} - \mathbf{I}_{|J} \frac{d}{dt}|_{t=0} \mathbf{I}a_{|J_t^{\dot{J}}} \\ &= -\frac{d}{dt}|_{t=0} a(J_t^{\dot{J}} \cdot) + \mathbf{I}_{|J} \frac{d}{dt}|_{t=0} a(J_t^{\dot{J}} \cdot) \\ &= -a(J\dot{J} \cdot) - a(\dot{J}J \cdot) \\ &= -a(J\dot{J} + \dot{J}J \cdot) = 0, \end{aligned}$$

where a is regarded as an element of $\Omega^1(\text{ad } E)$. Note that the integrability of \mathbf{I} follows without any assumption on the compatibility of $\dot{J} \in T_J\mathcal{J}$ with the symplectic structure. \square

Remark 2.3.3. Note that there is another $\tilde{\mathcal{G}}$ -invariant complex structures in $\mathcal{J} \times \mathcal{A}$ given by $\mathbf{I}_+(\dot{J}, a) = (J\dot{J}, a(J \cdot))$, such that the projection onto \mathcal{J} is pseudoholomorphic. However, it follows from the proof of the proposition above that this almost complex structure is not integrable and corresponds to a $\tilde{\mathcal{G}}$ -invariant almost Kähler structure for coupling constants $\alpha_0 > 0$, $\alpha_1 < 0$ in (2.3.1).

Suppose now that there exist Kähler structures on X with Kähler form ω . In the notation of §1.3.2, this means that the subspace $\mathcal{J}^i \subset \mathcal{J}$ of integrable almost complex structures compatible with ω is non-empty. For each $J \in \mathcal{J}^i$, let $\mathcal{A}_J^{1,1} \subset \mathcal{A}$ be the subspace of connections A such that $F_A \in \Omega_J^{1,1}(\text{ad } E)$ (cf. §1.2.2). Then we have a $\tilde{\mathcal{G}}$ -invariant (maybe singular) complex submanifold

$$(2.3.5) \quad \mathcal{P} \subset \mathcal{J} \times \mathcal{A},$$

consisting of pairs (J, A) with $J \in \mathcal{J}^i$ and $A \in \mathcal{A}_J^{1,1}$. Since \mathcal{P} is complex, it inherits a Kähler form by restriction and by $\tilde{\mathcal{G}}$ -invariance, μ_α restricts to a moment map on \mathcal{P} . The following system of coupled equations

$$(2.3.6) \quad \left. \begin{aligned} \Lambda_\omega F_A &= z \\ \alpha_0 S_J + \alpha_1 \Lambda_\omega^2(F_A \wedge F_A) &= c \end{aligned} \right\}$$

is the condition satisfied by the points in \mathcal{P} whose orbits are in the corresponding Kähler reduction

$$(2.3.7) \quad \mu_\alpha^{-1}(0)/\tilde{\mathcal{G}}.$$

Here, S_J is the scalar curvature determined by the metric $g_J = \omega(\cdot, J\cdot)$ and c is the topological constant c_0 given by (2.3.2).

The space of pairs \mathcal{P} has a geometric interpretation in terms of holomorphic structures that we will use in the rest of the thesis. Let G^c be the complexification of G with Lie algebra \mathfrak{g}^c and consider the associated principal G^c -bundle $E^c = E \times_G G^c$. Then \mathcal{P} parameterizes pairs consisting of a Kähler structure on X , with Kähler form ω , together with a structure of holomorphic principal G^c -bundle on E^c over the corresponding complex manifold X . Alternatively, \mathcal{P} is the space G^c -invariant integrable structures on the principal G^c -bundle E^c inducing holomorphic structures on X which are compatible with ω . Fixing the holomorphic structure on E^c and X , the coupled equations (2.3.6) have another formulation: we can consider the system as equations for pairs (ω, H) , where ω is a Kähler metric on X in a fixed Kähler class and H is a smooth G -reduction of the holomorphic G^c -bundle. The role of the connection A in (2.3.6) is played now by the Chern connection of H .

In order to interpret \mathcal{P} as holomorphic structures on E^c , we interpret first the bigger space $\mathcal{J} \times \mathcal{A}$ in terms of G^c -invariant almost complex structures on $\text{Tot}(E^c)$, the total space of E^c . The G^c -invariant complex structure on \mathfrak{g}^c induces an almost complex structure \mathbf{i} on the vertical bundle $VE^c \subset TE$, via the isomorphism $VE^c \cong \text{ad } E^c$. We have a map from $\mathcal{J} \times \mathcal{A}$, that assigns to each pair (J, A) the G^c -invariant complex structure

$$(2.3.8) \quad (J, A) \mapsto I_{J,A} := \mathbf{i}A + A^\perp J\pi,$$

on $\text{Tot}(E^c)$, where $\pi: E^c \rightarrow X$ is the projection and

$$A: TE^c \rightarrow VE^c \quad A^\perp: \pi^*TX \rightarrow TE^c$$

are the maps induced by the connection A (cf. (2.2.2)). On the other hand, any G^c -invariant almost complex structure I on $\text{Tot}(E^c)$ that preserves VE^c induces a unique pair of almost complex structures I_V on VE^c and I_X on TX . We claim that the image of (2.3.8) correspond to the space \mathcal{I} of G^c -invariant almost complex structures I on $\text{Tot}(E^c)$ that preserve VE^c and such that $I_V = \mathbf{i}$ and $I_X \in \mathcal{J}$. By definition, given $I \in \mathcal{I}$ it induces a unique element $J_I \in \mathcal{J}$. Moreover, for such an I there exists a unique G -connection $A \in \mathcal{A}$ such that $I = I_{J,A}$ (see [42], Theorem 1.1 p. 178 and the remark on p. 185), so the map

$$(2.3.9) \quad \mathcal{J} \times \mathcal{A} \rightarrow \mathcal{I}$$

given by (2.3.8) is a bijection. Note that (2.3.9) commutes with the projections of both spaces onto \mathcal{J} and it is $\tilde{\mathcal{G}}$ -equivariant, where the action of $\tilde{\mathcal{G}}$ on \mathcal{I} is given by pushforward. Moreover, (2.3.9) is a biholomorphisms with respect to the complex structure \mathbf{I} on $\mathcal{J} \times \mathcal{A}$ and the complex structure on \mathcal{I} given by

$$\mathbf{I}_I B = IB$$

for any $I \in \mathcal{I}$, where we regard $B \in T_I \mathcal{I}$ as a (G^c -invariant) endomorphism of TE^c . The integrability condition for an almost complex structure $I \in \mathcal{I}$ is equivalent in terms of the associated pair (J, A) to $J \in \mathcal{J}^i$ and $A \in \mathcal{A}_J^{1,1}$, hence equivalently to the condition $(J, A) \in \mathcal{P}$. It follows easily now from the naturality of the Nijenhuis tensor of any almost complex structure (see Proposition 4.11 in [55]) that \mathcal{P} is a $\tilde{\mathcal{G}}$ -invariant complex submanifold of $\mathcal{J} \times \mathcal{A}$.

Remark 2.3.4. Alternatively, taking $z = 0$ in (2.3.3), the coupled equations (2.3.6) can be defined by the condition (see (2.3.7) and (2.3.6))

$$\mu_\alpha(p) \in \mathfrak{z} \subset \text{Lie } \tilde{\mathcal{G}},$$

$p \in \mathcal{P}$, provided by the identification given by the $\tilde{\mathcal{G}}$ -equivariant family of parings $\langle \cdot, \cdot \rangle_{\tilde{\mathcal{G}}}$ given in (2.2.15). Given $p \in \mathcal{P}$ let $\tilde{\mathcal{G}}_p$ be its isotropy group in $\tilde{\mathcal{G}}$. This is a finite dimensional compact Lie group whose Lie algebra $\text{Lie } \tilde{\mathcal{G}}_p$ contains \mathfrak{z} . When $\mathfrak{z} \not\subseteq \text{Lie } \tilde{\mathcal{G}}_p$ it makes sense to generalize the condition $\mu_\alpha(p) \in \mathfrak{z}$ (above) to

$$(2.3.10) \quad \mu_\alpha(p) \in \text{Lie } \tilde{\mathcal{G}}_p$$

obtaining new, weaker, equations. This new condition is the analogue for our situation of the well-known notion of extremal Kähler metric introduced by Calabi in [15]. In fact the extremal pairs defined by (2.3.10) are critical points of the quadratic functional

$$(2.3.11) \quad p = (J, A) \rightarrow \|\mu_\alpha(p)\|_{\tilde{\mathcal{G}}|A}^2,$$

over the space $\mathcal{J} \times \mathcal{A}^{HYM}$, where \mathcal{A}^{HYM} is the set of connections satisfying $\Lambda_\omega F_A = z$. This functional can be considered as the norm-squared of the moment map with respect to the family of metrics $\langle \cdot, \cdot \rangle_{\tilde{\mathcal{G}}}$ indexed by \mathcal{A} . To see this, given any $p \in \mathcal{P}$ we define (cf. (2.2.16))

$$\zeta_p = 4\alpha_1 \Lambda_\omega F_A - \theta_A^\perp(\eta_p) \in \text{Lie } \tilde{\mathcal{G}},$$

with coupling constants $\alpha_0, \alpha_1, c \in \mathbb{R}$ as in (2.3.6), where η_p is the Hamiltonian vector field associated to the smooth function

$$\phi_p = \alpha_0 S_J + \alpha_1 \Lambda_\omega^2(F_A \wedge F_A) - c \in C_0^\infty(X).$$

Then, a solution to (2.3.10) is a critical point of (2.3.11) over $\mathcal{J} \times \mathcal{A}^{HYM}$ since for any curve $p_t \subset \mathcal{J} \times \mathcal{A}^{HYM}$ we have

$$\frac{d}{dt} \langle \mu_\alpha(p_t), \zeta_t \rangle = 2\omega_\alpha(Y_{\zeta_t}, \dot{p}_t),$$

where Y_{ζ_t} is the infinitesimal action of $\zeta_t = \zeta_{p_t}$.

2.4. Variational interpretation and canonical invariant metrics

In this section we give a variational interpretation of the coupled equations (2.3.6), as absolute minima of a suitable Calabi–Yang–Mills type functional. The unknown variables in the system (2.3.6) are considered to be a Kähler structure on the base, with fixed symplectic form, and a connection on the bundle. We will also interpret the solutions as a choice of canonical invariant metrics on the total space of the bundle, linking with the classical Kaluza–Klein Theory for metrics and connections (see e.g. [20]). Throughout this section we suppose that the coupling constants α_0 and α_1 in (2.3.6) have positive ratio

$$\frac{\alpha_1}{\alpha_0} > 0.$$

This will be a crucial assumption for our arguments. For simplicity, we consider that the compact Lie group G , with Lie algebra \mathfrak{g} , is semisimple. As in the previous section, we fix a metric (\cdot, \cdot) on \mathfrak{g} invariant by the adjoint action.

Let (X, ω) be a compact symplectic manifold and consider the space \mathcal{J} of ω -compatible almost complex structures on X . We suppose that (X, ω) is Kählerian, i.e. that the subspace $\mathcal{J}^i \subset \mathcal{J}$ of integrable complex structures is non empty. Let E be a smooth principal G -bundle over X . Given any $J \in \mathcal{J}$ we consider the associated Riemannian metric

$$(2.4.1) \quad J \rightarrow g_J = \omega(\cdot, J\cdot).$$

Since ω is fixed we can recover J from g_J , and so we will denote the latter simply by g . Given a pair (g, A) , where g is given by (2.4.1) and $A \in \mathcal{A}$ is a connection on E , we consider the functional

$$(2.4.2) \quad CYM(g, A) = \int_X (\alpha_0 S_g - 2\alpha_1 |F_A|^2)^2 \cdot \text{vol}_g + 2\alpha_1 \cdot \|F_A\|^2,$$

where vol_g is the volume form of the metric g . The term $\|F_A\|^2$ is the Yang–Mills functional, considered as a functional in both the metric and the connection, and the function $|F_A|^2$ is the point-wise norm with respect to metric g and the fixed inner product on \mathfrak{g} . Note that (2.4.2) is well defined for an arbitrary Riemannian metric in the base. However, only when considered as a functional for the pairs (g, A) , as before, we will obtain that our solutions are absolute minima of the functional. Let $c \in \mathbb{R}$, \hat{c} and \hat{S} be the real constant defined, respectively, in (2.3.2), (2.2.6) and (1.3.3). Recall that c , involved in the definition of the coupled system (2.3.6), equals

$$c = \alpha_0 \cdot \hat{S} + \alpha_1 \cdot \frac{2\hat{c}}{(n-2)!},$$

where \hat{S} and \hat{c} only depend on the topology of X and E . Then we have the following.

Proposition 2.4.1. *Suppose that α_0 and α_1 are positive and small enough such that $2c < 1$. Then any pair (g_J, A) satisfying the coupled equations (2.3.6) and the compatibility condition $F_A^{0,2} = 0$ is an absolute minimum of CYM .*

PROOF. For any pair (g, A) as in (2.4.2), we have (cf. e.g. proof of Lemma 7.9 in [57])

$$(2.4.3) \quad \frac{1}{2} \Lambda_\omega^2(F_A \wedge F_A) = -|F_A|^2 + |\Lambda_\omega F_A|^2 + 4|F_A^{0,2}|^2$$

where the point-wise norms are taken with respect to the metric g and the inner product on \mathfrak{g} . Then,

$$\begin{aligned} CYM(g, A) &= \int_X (\alpha_0 S_g - 2\alpha_1 |F_A|^2 - c)^2 \cdot \text{vol}_g + 2\alpha_1 \cdot \|F_A\|^2 \\ &\quad + 2c \cdot \int_X (\alpha_0 S_g - 2\alpha_1 |F_A|^2) \cdot \text{vol}_g + c^2 \cdot \text{Vol}([\omega]) \\ &= \|\alpha_0 S_g - 2\alpha_1 |F_A|^2 - c\|^2 \\ &\quad + 2\alpha_1(1-2c) \cdot (\|\Lambda_\omega F_A\|^2 + 4\|F_A^{0,2}\|^2) \\ &\quad + (2c \cdot \alpha_0 \hat{S} + c^2 + \alpha_1(1-2c) \cdot \frac{2\hat{c}}{(n-2)!}) \cdot \text{Vol}([\omega]), \end{aligned}$$

where $\text{Vol}([\omega]) = \frac{[\omega]^n}{n!}$, that only depends on the fixed symplectic form. The statement follows now from (2.4.3). \square

Remark 2.4.2. Note that the condition $2c < 1$ imposes no restriction since we can always multiply the scalar equation in the system (2.3.6) by a small positive constant. Note also that $F_A^{0,2} = 0$ is the condition for the pair (J, A) to be in the space of integrable pairs \mathcal{P} (2.3.5), provided that $J \in \mathcal{J}^i$. An analogue result, as in Proposition 2.4.1, is obtained if we fix a complex structure on X and let ω vary in a fixed Kähler class.

In order to understand better this functional we think in the data (g, A) as defining a G -invariant metric \hat{g} on $\text{Tot}(E)$, the total space of E . Let $\pi: E \rightarrow X$ be the standard projection. The biinvariant positive definite inner product (\cdot, \cdot) on \mathfrak{g} induces a Riemannian metric g_V on the vertical bundle VE , via its identification with $\text{ad } E$. Given a connection A on E and a Riemannian metric g on X we can associate a G -invariant metric \hat{g} on $\text{Tot}(E)$ given by

$$\hat{g} = \pi^* g + g_V(\theta_A \cdot, \theta_A \cdot),$$

where $\theta_A: TE \rightarrow VE$ denotes the vertical projection determined by A . Endowed with this metric the projection $\pi: (E, \hat{g}) \rightarrow (X, g)$ is a Riemannian submersion with totally geodesic fibres (see [11, Theorem 9.59]). The scalar curvature \hat{S} of \hat{g} is given by (see [11, Proposition 9.70])

$$(2.4.4) \quad S_{\hat{g}} = S_g - |F_A|^2 + s,$$

where S_g is the scalar curvature of g , $|F_A|^2$ is as in (2.4.3) and s is a real constant that corresponds to the scalar curvature of the metric in the fibres. Note that, since \hat{g} is G -invariant, $S_{\hat{g}}$ induces a well defined function on X . We can scale the metric in the fibres multiplying by a constant factor $t > 0$ and then the associated metric \hat{g}_t has scalar curvature (see [11, Proposition 9.70])

$$\hat{S}_t = S_g - t|F_A|^2 + \frac{1}{t}s.$$

We set $t = \frac{2\alpha_1}{\alpha_0} > 0$, where the coupling constants are as in (2.4.2). Then, the first summand in the definition of the functional (2.4.2) can be thought of, up to scaling, as a Calabi functional for G -invariant metrics in $\text{Tot}(E)$. This follows from the invariance of the scalar curvature of a invariant metric on E , which implies that the integration on the bundle reduces to an integration on the base (see [20]).

The previous expression for the scalar curvature (2.4.4) of the metric \hat{g} links with our next interpretation of the equations (2.3.6): they define a ‘good’ choice of canonical G -invariant metrics on $\text{Tot}(E)$, in the sense that a moduli construction for such metrics would lead to a Kähler moduli space. Let \hat{g} be the G -invariant metric on $\text{Tot}(E)$ determined by a pair (g, A) , where g is given by (2.4.1) and $A \in \mathcal{A}$ is a connection on E . We fix the scale on the fibers

$$(2.4.5) \quad t = \frac{2\alpha_1}{\alpha_0} > 0$$

as before. We assume further that the connection A is irreducible, $J \in \mathcal{J}^i$ (in the definition of g), and that

$$F_A^{0,2} = 0.$$

Therefore, (g, J, ω) is a Kähler structure on X and A determines a holomorphic structure on the G^c -bundle $E^c = E \times_G G^c$. In order to understand the conditions imposed by the system (2.3.6) to the metric \hat{g} , we will first interpret the HYM condition (1.2.5) in terms of the Einstein equation for \hat{g} . Using the irreducibility of the connection, it follows from the Kähler identities that the HYM equation is equivalent to the a priori weaker Yang-Mills equation

$$d_A^* F_A = 0.$$

For the case $G = U(r)$ see [21, Proposition 3]. In general it follows by taking a faithful representation of the complexification G^c in $GL(r, \mathbb{C})$. The isomorphism

$$TE \cong VE \oplus \pi^* TX,$$

provided by the connection A , splits the Ricci tensor $R_{\hat{g}}$ and so the Einstein equation can be written as a system with respect to this splitting. Let R_{hv} be the horizontal-vertical part of $R_{\hat{g}}$. It follows now from [11, Proposition 9.61] that the Yang-Mills equation for \hat{g} is equivalent to the vertical-horizontal Einstein equation $R_{hv} = 0$. On the other hand, note that the scalar equation in the coupled system (2.3.6) can be written as

$$\alpha_0 S_g - 2\alpha_1 \cdot |F_A|^2 = c,$$

using (2.4.3). Hence, it follows from (2.4.4) that: the tuple (g, ω, J, A) satisfies the coupled equations with coupling constants α_0 and α_1 if and only if \hat{g} , with scale (2.4.5) on the fibers, has constant scalar curvature and satisfies the vertical-horizontal Einstein equation $R_{hv} = 0$. Since any Einstein metric has constant scalar curvature it follows from (2.4.6) that

$$R_{\hat{g}} = \lambda \hat{g} \quad \Rightarrow \quad \left. \begin{aligned} \Lambda F_A &= 0 \\ \alpha_0 S_g + \alpha_1 \Lambda^2(F_A \wedge F_A) &= c \end{aligned} \right\} \quad \Rightarrow \quad S_{\hat{g}} = c,$$

where λ and c are real constants. A crucial ingredient for this chain of implications is the compatibility condition $F_A^{0,2} = 0$, between the Kähler structure (g, ω, J) and the connection A . In the language of the moment map interpretation given in §2.3,

this condition implies that $(J, A) \in \mathcal{P}$ (0.0.1). An unexpected point here is that the positivity of the ratio $\frac{\alpha_1}{\alpha_0} > 0$, that we have assumed to build our Riemannian metrics \hat{g} from the data (g, A) , is precisely the condition for the symplectic form ω_α (0.0.2) in \mathcal{P} to be Kähler. Thus, a moduli construction for such metrics using the symplectic reduction process will lead to a Kähler moduli space. In this sense, the coupled equations (0.0.3) provide a good choice of canonical G -invariant metrics in $\text{Tot}(E)$, well suited for the framework of Kähler geometry.

It is natural to ask which coupled system of equations is equivalent to the Einstein condition for the metric \hat{g} . From [11, Proposition 9.64] it follows that the Einstein equation for such a metric (with irreducible connection) and arbitrary compact group G reads

$$(2.4.6) \quad \left. \begin{aligned} \Lambda_\omega F_A &= z \\ \alpha_0(\rho - c'\omega) &= \alpha_1(2(\Lambda_\omega F_A, F_A) - \Lambda_\omega(F_A \wedge F_A) - c''\omega) \end{aligned} \right\},$$

where z is an element in the centre of \mathfrak{g} and $c', c'' \in \mathbb{R}$ are suitable constants. The 2-form ρ denotes the Ricci form of the Kähler structure on the base.

Remark 2.4.3. The system (2.4.6) can be considered a Kähler analogue of the Einstein-Yang-Mills equation considered in the pseudoRiemannian context (see e.g. [20]). Note that in complex coordinates we have that

$$2(\Lambda_\omega F_A, F_A) - \Lambda_\omega(F_A \wedge F_A) = -\mathbf{i} \operatorname{tr} g^{\alpha, \bar{\beta}} (F_{\alpha, \bar{\delta}} F_{\gamma, \bar{\beta}} + F_{\gamma, \bar{\beta}} F_{\alpha, \bar{\delta}}) \cdot dz_\gamma \wedge d\bar{z}_\delta,$$

where

$$\omega = \mathbf{i} g_{\alpha, \bar{\beta}} \cdot dz_\alpha \wedge d\bar{z}_\beta \quad \text{and} \quad F_A = F_{\alpha, \bar{\beta}} \cdot dz_\alpha \wedge d\bar{z}_\beta.$$

Remark 2.4.4. The Einstein condition for the metric \hat{g} implies strong restrictions on the tuple (g, ω, J, A) . To see this, suppose that \hat{g} is an Einstein metric. Then, since the Ricci form ρ and ω in (2.4.6) are closed so $\Lambda_\omega(F_A \wedge F_A)$ is also closed. It follows then from the Kähler identities that the 4-form $(F_A \wedge F_A)$ is harmonic. From this one has that $\Lambda^2(F_A \wedge F_A)$ is constant. Then, taking traces in the second equation of (2.4.6) it follows that the Kähler metric on the base g has constant scalar curvature.

CHAPTER 3

Analytic obstructions and existence of solutions

Let (X, J) be a compact complex manifold, G^c be a complex Lie group and (E^c, I) a holomorphic principal G^c -bundle over (X, J) . We fix a maximal compact subgroup $G \subset G^c$. Then, we can consider the coupled equations (0.0.3) as equations in the space of pairs $(\omega, H) \in \tilde{\mathcal{K}}$, consisting of a Kähler metric ω in a fixed Kähler class and a reduction H of E^c to a principal G -bundle. Thus, for positive real coupling constants α_0, α_1 the *coupled equations* can be rewritten as

$$(3.0.7) \quad \left. \begin{array}{l} \Lambda_\omega F_H = z \\ \alpha_0 S_\omega + \alpha_1 \Lambda_\omega^2(F_H \wedge F_H) = c \end{array} \right\}.$$

Here, S_ω is the scalar curvature determined by the Kähler metric, F_H is the curvature of the Chern connection determined by the reduction H and the complex structure I , and $z \in \mathfrak{z}$ and $c \in \mathbb{R}$ are as in (2.3.6). In this chapter we give an existence criterion for the existence of solutions to the coupled system (3.0.7) for small values of the ratio of the coupling constants $\frac{\alpha_1}{\alpha_0} > 0$ (Theorem 3.2.4). This criterion is formulated in terms of a real subgroup $\tilde{\mathcal{G}}_I$ of the automorphism group $\text{Aut}(E^c, I)$ of the holomorphic bundle. We will also give an obstruction (Proposition (3.3.3)) controlled by a character

$$\mathcal{F}_\alpha: \text{Lie Aut}(E^c, I) \rightarrow \mathbb{C}$$

that generalizes the usual Futaki invariant for the cscK problem [29] and also other invariants defined by A. Futaki in [30] (see §3.3). In §3.4 we construct the integral of the moment map

$$\mathcal{M}_\kappa: \tilde{\mathcal{K}} \rightarrow \mathbb{R}$$

and determine a sufficient condition for its convexity and existence of lower bounds in terms of the existence of suitable smooth curves on $\tilde{\mathcal{K}}$ joining any two points. In §3.5 we restrict to the case in which the base is a compact complex surface. We use the integral of the moment map to prove some results concerning the uniqueness and the existence problem for the coupled equations (3.0.7).

To obtain these results, we first discuss in §3.1 the general framework provided by Kempf-Ness Theorem where our problem fits and establish its relation with §2.3: the pairs (ω, H) satisfying (3.0.7) correspond to zeros of the moment map μ_α (2.3.3) in the “complexified orbit” of the G^c -invariant complex structure I of E^c with respect to the action of the extended gauge group.

Throughout this chapter, we suppose that the G -invariant metric (\cdot, \cdot) in \mathfrak{g} used to define the system (3.0.7) extends to a G^c -invariant (maybe degenerate) bilinear pairing

$$(3.0.8) \quad (\cdot, \cdot): \mathfrak{g}^c \otimes \mathfrak{g}^c \rightarrow \mathbb{C}.$$

Note that the pullback of minus the trace with respect to any faithful representation of G^c in $GL(r, \mathbb{C})$ mapping G into $U(r)$ always satisfies this property.

3.1. General framework

In this section we show how the problem of finding solutions to the coupled equations (3.0.7) fits into the general framework provided by the Kemp–Ness Theorem, already explained in §1.1. Our argument combines the one for HYM connections explained in §1.2 with the one given by Donaldson for the cscK problem in [23] (see §1.3.3). We will study the system (3.0.7) with fixed positive coupling constants $\alpha_0, \alpha_1 \in \mathbb{R}$.

Let (X, J) be a compact complex manifold, where we fix a Kähler class $[\omega]$, and denote by \mathcal{K} the space of Kähler metrics on X with this Kähler class. Let (E^c, I) be a holomorphic principal G^c -bundle over (X, J) and denote by \mathcal{R} the space of smooth G -reductions on the underlying smooth bundle E^c . An element $H \in \mathcal{R}$ will be regarded, depending on the context, as a smooth cross section of the bundle $E^c/G \rightarrow X$ or as a G -subbundle of E^c . Let us consider the product space

$$(3.1.1) \quad \tilde{\mathcal{K}} = \mathcal{K} \times \mathcal{R}.$$

The elements in $\tilde{\mathcal{K}}$ will be denoted either by (ω, H) or κ , when there is no confusion. First we will establish the relation with §2.3, showing that the pairs (ω, H) correspond to points in the “complexified orbit” of the holomorphic G^c -invariant structure I on E^c for the action of the extended gauge group. Those pairs satisfying the coupled equations (3.0.7) will be given by the zero locus of the moment map μ_α (2.3.3) inside the complexified orbit. It is convenient to consider the following definition (cf. [51, Section 1.1]).

Definition 3.1.1. A *marked Kähler Hamiltonian $\tilde{\mathcal{G}}$ -space* (mK-Hamiltonian space, for short) is a tuple $(\mathcal{P}, \mathbf{I}, \omega_\alpha, \tilde{\mathcal{G}}, \mu_\alpha, p)$, where $(\mathcal{P}, \mathbf{I}, \omega_\alpha)$ is a Kähler manifold, $\tilde{\mathcal{G}}$ is a Lie group that acts on \mathcal{P} by holomorphic isometries,

$$\mu_\alpha : \mathcal{P} \rightarrow \text{Lie } \tilde{\mathcal{G}}^*$$

is a $\tilde{\mathcal{G}}$ -equivariant moment map and p is a marked point in \mathcal{P} .

Relying on the results in §2.3, any $\kappa = (\omega, H) \in \tilde{\mathcal{K}}$ determines a mK-Hamiltonian space

$$(3.1.2) \quad (\mathcal{P}, \mathbf{I}, \omega_\alpha, \tilde{\mathcal{G}}, \mu_\alpha, p_I),$$

where $\mathcal{P} \subseteq \mathcal{J}^i \times \mathcal{A}$ (see (2.3.5)) is given by pairs $p = (J', A')$ such that $J' \in \mathcal{J}^i$, the space of compatible complex structures on (X, ω) , and A' is a compatible connection with J' on the G -reduction H , i.e such that $F_A^{0,2} = 0$. Recall that \mathcal{P} parameterizes G^c -invariant complex structures on E^c that cover elements in \mathcal{J}^i (see (2.3.8)). Therefore, the marked point $p_I = (J, A) \in \mathcal{P}$ is determined by the complex structure I , where A denotes the Chern connection of H on (E^c, I) . The extended gauge group $\tilde{\mathcal{G}}$ fits into the short exact sequence

$$1 \rightarrow \mathcal{H} \rightarrow \tilde{\mathcal{G}} \rightarrow \mathcal{G} \rightarrow 1,$$

where \mathcal{H} and \mathcal{G} denote respectively the Hamiltonian group of symplectomorphisms of (X, ω) and the gauge group of the reduction H . The symplectic structure, the moment map and the complex structure in (3.1.2) are given respectively by (2.3.1), (2.3.3) and (2.3.4), where the two former are determined by the fixed coupling constants α_0 and α_1 .

Let us fix $\kappa \in \tilde{\mathcal{K}}$ and consider the associated mK-Hamiltonian space (3.1.2). As \mathcal{H} , the group $\tilde{\mathcal{G}}$ does not admit a complexification extending the action on \mathcal{P} . However, there is an integrable distribution $D \subseteq T\mathcal{P}$ whose leaves play the role of the orbits under the action of an ideal complexification. Recall that given $\zeta \in \text{Lie } \tilde{\mathcal{G}}$, we denote by Y_ζ the infinitesimal action of ζ on \mathcal{P} . Then, at a point $p \in \mathcal{P}$ the distribution is given by

$$(3.1.3) \quad D_p = \{Y_{\zeta|p}; \zeta \in \text{Lie } \tilde{\mathcal{G}}^c\},$$

where $\text{Lie } \tilde{\mathcal{G}}^c$ is the complexification of $\text{Lie } \tilde{\mathcal{G}}$ and

$$(3.1.4) \quad Y_\zeta = Y_{\zeta_0} + \mathbf{I}Y_{\zeta_1} \text{ for all } \zeta = \zeta_0 + \mathbf{i}\zeta_1 \in \text{Lie } \tilde{\mathcal{G}}^c.$$

The distribution D is closed under the Lie bracket, since the Nijenhuis tensor of the complex structure \mathbf{I} vanishes (see Proposition 2.3.2), although it is not obvious from this fact that it integrates, since \mathcal{P} is infinite dimensional. To see that D indeed integrates, we define a principal $\tilde{\mathcal{G}}$ -bundle

$$(3.1.5) \quad \tilde{\mathcal{Y}} \rightarrow \tilde{\mathcal{K}}$$

that plays the role of the missed complexification, with $\tilde{\mathcal{K}}$ regarded as the symmetric space in §1.1 (see Remark 3.1.2 below). For simplicity, we suppose that the first cohomology group $H^1(X)$ vanishes, so the group \mathcal{H} coincides with $\text{Symp}_0(X, \omega)$, the identity component of the group of symplectomorphisms. The assumption $H^1(X) = 0$ can be easily removed in the discussion replacing $\tilde{\mathcal{G}}$ by the bigger group of smooth G -invariant automorphisms of the reduction H that project to $\text{Symp}_0(X, \omega)$. To define $\tilde{\mathcal{Y}}$, let us consider $\text{Aut}_0(E^c)$, the connected component of Id in the group $\text{Aut}(E^c)$ of G^c -equivariant diffeomorphisms of E^c . Note that any element $f \in \text{Aut}_0(E^c)$ covers a unique diffeomorphism \check{f} in $\text{Diff}_0 X$. Then, the $\tilde{\mathcal{G}}$ -bundle $\tilde{\mathcal{Y}}$ determined by $\kappa = (\omega, H)$ consists of triples

$$(3.1.6) \quad \tilde{\mathcal{Y}} := \{(\omega', H', f) \in \tilde{\mathcal{K}} \times \text{Aut}_0(E^c); (\check{f}^*\omega', f^*H') = (\omega, H)\}.$$

The reduction H' is viewed here as a cross section of the bundle $E^c/G \rightarrow X$ and so for any point $x \in X$ we have $f^*H'(x) = f^{-1}(H'(\check{f}(x)))$. Given any $p \in \mathcal{P}$, the map

$$(3.1.7) \quad \tau_p: \tilde{\mathcal{Y}} \rightarrow \mathcal{P}: (\omega', H', f) \mapsto f^*p$$

plays the role of the complexified action, in the sense that the image of this map is a leaf of D through p . Let us consider the *complexified orbit* $\tau_{p_I}(\tilde{\mathcal{Y}})$ of the marked point $p_I \in \mathcal{P}$ determined by the complex structure I on E^c . Given a pair in $\tilde{\mathcal{K}}$, we can associate a point in the complexified orbit of p_I via the map (3.1.7), modulo an indeterminacy coming from the extended gauge group $\tilde{\mathcal{G}}$. Then, it can be readily checked that $\tau_{p_I}(\tilde{\mathcal{Y}})$ contains a zero of the $\tilde{\mathcal{G}}$ -equivariant moment map μ_α (see (2.3.3)) if and only if there exists a pair in $\tilde{\mathcal{K}}$ satisfying the coupled equations (3.0.7).

Remark 3.1.2. When the isotropy group of p_I vanishes, the previous construction shows that $\tilde{\mathcal{Y}}$ is an infinitesimal complexification of $\tilde{\mathcal{G}}$ in the sense of [23]. This would endow $\tilde{\mathcal{K}} \cong \tilde{\mathcal{Y}}/\tilde{\mathcal{G}}$ with a structure of Riemannian symmetric space, applying [23, Proposition 2], provided that $\text{Lie } \tilde{\mathcal{G}}$ has a biinvariant metric. However such a metric is not known to exist by the author (cf. (2.2.15)). A priori, i.e. removing the fixed point in $\tilde{\mathcal{K}}$, there is a Riemannian submersion structure on the fibration

$$\tilde{\mathcal{K}} = \mathcal{K} \times \mathcal{R} \rightarrow \mathcal{K},$$

where the space of Kähler metrics \mathcal{K} is regarded as a symmetric space with the Mabuchi metric (see [48]). Note that for each $\omega \in \mathcal{K}$, the fiber $\mathcal{R} \cong \Omega^0(X, E^c/G)$ has a structure of Riemannian symmetric space given by the one on G^c/G and integration over X with respect to ω . Thus, the structure of Riemannian submersion on $\tilde{\mathcal{K}}$ has trivial connection and varying Riemannian metric on each fiber. It is very likely that this structure is in fact Riemannian symmetric, in the sense that it admits a Levi–Civita connection whose curvature is negative, covariantly constant and admits an explicit expression in terms of the Lie bracket on the Lie algebra of the extended gauge group (see [23, Proposition 2]). To make sense of the latter condition, note that for any element $\kappa = (\omega, H) \in \tilde{\mathcal{K}}$ there is a canonical isomorphism

$$(3.1.8) \quad T_\kappa \tilde{\mathcal{K}} \cong \text{Lie } \tilde{\mathcal{G}}$$

via the identifications $T_\omega \mathcal{K} \cong C^\infty(X)/\mathbb{R} \cong \text{Lie } \mathcal{H}$ and $T_H \mathcal{R} \cong \mathbf{i} \text{Lie } \mathcal{G}$ and the vector space splitting $\text{Lie } \tilde{\mathcal{G}} \cong \text{Lie } \mathcal{H} \times \text{Lie } \mathcal{G}$ provided by the Chern connection of H on (E^c, I) . Here the groups \mathcal{G} , $\tilde{\mathcal{G}}$ and \mathcal{H} are those corresponding to the mK-Hamiltonian space (3.1.2) determined by κ .

We will now construct a map that will be used in §3.2 to deform a solution $\kappa \in \tilde{\mathcal{K}}$ of the coupled equations (3.0.7) on the attached complexified orbit $\tau_{p_I}(\tilde{\mathcal{Y}})$. This will provide an existence criterion for (3.0.7) and also canonical curves in $\tilde{\mathcal{K}}$ along which the integral of the moment map (see §3.4), whose zeros locus correspond to the solutions of (3.0.7), will be convex. It is very likely that these curves are, in fact, the geodesics on $\tilde{\mathcal{K}}$ with respect to the Riemannian structure discussed in Remark 3.1.2. Let us fix a point $\kappa \in \tilde{\mathcal{K}}$. We will construct a map

$$\rho: T_\kappa \tilde{\mathcal{K}} \rightarrow \tau_{p_I}(\tilde{\mathcal{Y}}) \subset \mathcal{P},$$

where $p_I \in \mathcal{P}$, τ_{p_I} and $\tilde{\mathcal{Y}}$ are, as before, associated to the mK-Hamiltonian space (3.1.2) determined by κ . Note that in the finite dimensional picture provided by §1.1 the complex exponential defines a canonical connection on the principal bundle

$$G^c \rightarrow G^c/G \cong \mathbf{i} \mathfrak{g},$$

that can be used to map G^c/G into the Kähler manifold where the complex group acts. Hence, via the identification $T_\kappa \tilde{\mathcal{K}} \cong \mathbf{i} \text{Lie } \tilde{\mathcal{G}}$ (see (3.1.8)), the map ρ can be considered as an analogue of the composition of the (missed) complex exponential with the action on p_I of the (missed) complexification of $\tilde{\mathcal{G}}$. The strategy we will follow is: first, we will define a connection on the $\tilde{\mathcal{G}}$ -bundle $\tilde{\mathcal{Y}} \rightarrow \tilde{\mathcal{K}}$. Then, choosing a suitable family of curves in $\tilde{\mathcal{K}}$ joining any point to κ , we will lift the curves starting at the canonical

element on the fiber of $\tilde{\mathcal{Y}}$ over κ given by the identity. Finally, using the map τ_{p_I} (3.1.7) we will map the lifted $\tilde{\mathcal{K}}$ into \mathcal{P} .

To define the connection we define the horizontal lift of a smooth curve κ_t in $\tilde{\mathcal{K}}$ passing through $f \in \tilde{\mathcal{Y}}$ in the fiber over the initial point κ_0 . In the following we will denote by a t -subscript the elements of the mK-Hamiltonian space determined by κ_t (3.1.2). Let $\kappa_t = (\omega_t, H_t)$ be a curve in $\tilde{\mathcal{K}}$. As in §1.3.3, let $\check{f}_t \subset \text{Diff}_0(X)$ be the Moser curve of diffeomorphisms of ω_t determined by the time dependent vector field \check{y}_t with initial condition $\check{f}_0 = f$, the diffeomorphism covered by $f \in \text{Aut}_0(E^c)$. Recall that \check{f}_t satisfies the condition $\check{f}_t^* \omega_t = \omega_0$ for all t . We now lift \check{f}_t to a curve in the group $\text{Aut}_0(E^c)$ with the desired property $f_t^* H_t = H$ (see (3.1.6)). Let $f_t \subset \text{Aut}_0(E^c)$ be the curve integrating

$$(3.1.9) \quad y_t = -\frac{1}{2} \dot{H}_t + \theta_t^\perp(\check{y}_t)$$

with initial condition $f_0 = f$, where $\dot{H}_t \in \mathbf{i} \text{Lie } \mathcal{G}_t$ is the H_t -symmetric vector field determined by $\frac{d}{dt} H_t$ and $\theta_t^\perp(\check{y}_t)$ is the horizontal lift of \check{y}_t with respect to the (Chern) connection of H_t on (E^c, I) .

Proposition 3.1.3. *The curve $f_t \subset \text{Aut}_0(E^c)$ defined by (3.1.9) is a lift of (ω_t, H_t) to $\tilde{\mathcal{Y}}$ passing through f . The induced curve $p_t = \tau_p(\omega_t, H_t, f_t) \subset \mathcal{P}$ determined by (3.1.7) satisfies*

$$(3.1.10) \quad \dot{p}_t = \mathbf{I} Y_{\zeta_t},$$

where Y_{ζ_t} is the infinitesimal action at $p_t \in \mathcal{P}$ of

$$(3.1.11) \quad \zeta_t = f_t^*(I y_t) \in \text{Lie } \tilde{\mathcal{G}}.$$

PROOF. For the first part of the statement it is enough to check that $f_t^* H_t$ is constant on t . Note first that the derivative of $f_t^* H$ for any $H \in \mathcal{R}$ is given by

$$\frac{d}{dt} f_t^* H = f_t^*(2\pi_1 \theta_H(y_t))$$

where θ_H is the (Chern) connection 1-form of H on the holomorphic bundle (E^c, I) and

$$\pi_1: \text{Lie } \mathcal{G}^c \rightarrow \mathbf{i} \text{Lie } \mathcal{G}$$

is the standard projection from the Lie algebra of the complex gauge group of E^c onto the Lie algebra of the gauge group of the G -reduction H . When the G -reduction is parameterized by t , we denote this map by π_1^t . Then we have

$$\frac{d}{dt} f_t^* H_t = f_t^*(2\pi_1^t \theta_t(y_t) + \dot{H}_t) = f_t^*(-\dot{H}_t + \dot{H}_t) = 0.$$

To prove (3.1.10), recall from §2.3 that the space \mathcal{P} is complex isomorphic to a suitable space of integrable G^c -invariant complex structures on E^c and that this isomorphism is $\tilde{\mathcal{G}}$ -equivariant (see (2.3.9)). Then, any element of the curve $p_t \in \mathcal{P}$ determines a unique G^c -invariant complex structure $I_t = f_t^* I$ on E^c and so

$$\frac{d}{dt} I_t = L_{f_t^* y_t} I_t = -I_t L_{f_t^*(I y_t)} I_t = \mathbf{I} Y_{f_t^*(I y_t)}.$$

Thus, $\zeta_t = f_t^*(Iy_t)$ as required. Note that $Iy_t \in \tilde{\mathcal{G}}_t$ implies that $\zeta_t \in \text{Lie } \tilde{\mathcal{G}}$. \square

The previous proposition shows that the connection we have just defined satisfies reasonable properties that bring us close to the map ρ of our interest. Recall that we have fixed a pair $\kappa = (\omega, H)$. Then, given any function $\dot{\phi} \in C_0^\infty(X)$, with zero integral with respect to vol_ω , and any $\dot{H} \in \mathbf{i}\text{Lie } \mathcal{G}$, we identify $(\dot{\phi}, \dot{H})$ with the vector field in the tangent space $T_\kappa \tilde{\mathcal{K}}$ given by the curve

$$(3.1.12) \quad \kappa_t = (\omega_t, H_t) := (\omega_{t\dot{\phi}}, (e^{t\dot{H}})^* H) \subset \tilde{\mathcal{K}},$$

where

$$\omega_{t\dot{\phi}} = \omega + tdd^c \dot{\phi}$$

and e denotes the exponential map in the complex gauge group \mathcal{G}^c of the smooth G^c -bundle E^c . Given $\dot{\kappa} = (\dot{\phi}, \dot{H}) \in T_\kappa \tilde{\mathcal{K}}$ we define the map ρ as

$$(3.1.13) \quad \rho: T_\kappa \tilde{\mathcal{K}} \rightarrow \tau_{p_I}(\tilde{\mathcal{Y}}) : \dot{\kappa} \mapsto \rho(\dot{\kappa}) := f_1^* p_I,$$

where f_1 is the value at $t = 1$ of the horizontal lift $f_t \subset \text{Aut}_0(E^c)$ (see (3.1.9)) of the curve κ_t (3.1.12) starting at $f_0 = \text{Id}$. Equivalently we can write (see (3.1.7))

$$\rho(\dot{\kappa}) = \tau_{p_I}(\kappa_1, f_1).$$

Note that this map is only well defined on the open subset given by those pairs $(\dot{\phi}, \dot{H})$ such that $\dot{\phi}$ is a Kähler potential on (X, J, ω) . Given $\dot{\kappa} = (\dot{\phi}, \dot{H}) \in T_\kappa \tilde{\mathcal{K}}$ as before, we will denote

$$(3.1.14) \quad \zeta_{\dot{\kappa}} = -\mathbf{i}\dot{H} + \theta_A^\perp(\eta_{\dot{\phi}}) \in \text{Lie } \tilde{\mathcal{G}},$$

where $\eta_{\dot{\phi}} \in \text{Lie } \mathcal{H}$ is the Hamiltonian vector field on (X, ω) corresponding to $\dot{\phi}$ and $\theta_A^\perp(\eta_{\dot{\phi}})$ denotes its horizontal lift with respect to the connection A determined by the marked point $p_I \in \mathcal{P}$. Then, by Proposition 3.1.3, we have the following.

Corollary 3.1.4. *The differential of ρ at the origin $0 \in T_\kappa \tilde{\mathcal{K}}$ is*

$$(3.1.15) \quad d\rho|_0(\dot{\kappa}) = \mathbf{I}Y_{\zeta_{\dot{\kappa}}},$$

where $Y_{\zeta_{\dot{\kappa}}}$ is the infinitesimal action of $\zeta_{\dot{\kappa}}$ given by (3.1.14) at $p_I \in \mathcal{P}$.

PROOF. We claim that the horizontal lift f_t of the curve κ_t (3.1.12) is defined by the time dependent vector field

$$(3.1.16) \quad y_t = -I(-\mathbf{i}\dot{H} + \theta_t^\perp(J\check{y}_t)),$$

where \check{y}_t is the vector field on X covered by y_t . To prove this, note that the H_t -symmetric tangent vector \dot{H}_t is given by

$$(3.1.17) \quad \dot{H}_t = s_t^*(s_t^{-1}\dot{s}_t + \dot{s}_t s_t^{-1}) = 2s_t^*(\pi_1 y_{s_t}) \in \mathbf{i}\text{Lie } \mathcal{G}_t,$$

where $s_t = e^{t\dot{H}}$, and so $y_{s_t} = \dot{s}_t s_t^{-1} = \dot{H}$ implies $\dot{H}_t = 2\dot{H}$ for all t . Then, it is straightforward from (3.1.9) that

$$\begin{aligned} y_t &= -\dot{H} + \theta_t^\perp(\check{y}_t) \\ &= -I(-\mathbf{i}\dot{H} + \theta_t^\perp(J\check{y}_t)), \end{aligned}$$

as claimed. Recall from §1.3.3 that $J\check{y}_t = y_{\dot{\phi}, \omega_{t\dot{\phi}}}$, where the latter is the Hamiltonian vector field of $\dot{\phi}$ with respect to $\omega_{t\dot{\phi}}$. Then, given a curve $s \mapsto (s\dot{\phi}, s\dot{H}) \in T_\kappa \tilde{\mathcal{K}}$, we have $\rho(s\dot{\phi}, s\dot{\xi}) = (f_1(s))^* p_I$ by definition, where $f_1(s)$ is the value at $t = 1$ of the curve $f_t(s) \subset \text{Aut}(E^c)$ determined by

$$\begin{aligned} y_t(s) &= -I(-\mathbf{i}s\dot{H} + \theta_{ts}^\perp(y_{s\dot{\phi}, \omega_{ts\dot{\phi}}})) \\ &= -sI(-\mathbf{i}\dot{H} + \theta_{ts}^\perp(y_{\dot{\phi}, \omega_{ts\dot{\phi}}})) \\ &= y_{ts} \end{aligned}$$

and $f_0(s) = \text{Id}$. Then, $\rho(s\dot{\phi}, s\dot{\xi}) = (f_{ts})_{|t=1}^* p_I = f_s^* p_I$ and (3.1.15) follows from (3.1.10) in Proposition 3.1.3. \square

The map (3.1.13) behaves at the origin, due to the identity (3.1.15), as the composition of the missed complex exponential with the action of the complexified group on the marked point p_I . In the next section we will compose ρ with the moment map μ_α given by (2.3.3) obtaining a new map that will allow us to deform a given solution to the coupled equations (3.0.7), thereby obtaining a criterion of existence. Note that the choice of curve κ_t made in (3.1.12) for the definition of ρ is not canonical. In fact, due to the identity (3.1.10), it would have been more natural to consider curves κ_t such that the associated vector field $f_t^*(Iy_t) \in \text{Lie } \tilde{\mathcal{G}}$ (3.1.11) satisfies

$$(3.1.18) \quad \frac{d}{dt} f_t^*(Iy_t) = 0.$$

If $\kappa_t = (\omega_t, H_t)$ solves (3.1.18), the Kähler metric ω_t is a geodesic in \mathcal{K} (see Proposition (3.4.1)). This provides an evidence of the fact that (3.1.18) is the geodesic equation with respect to the Riemannian structure on $\tilde{\mathcal{K}}$ described in Remark 3.1.2. Recall that that the Riemannian submersion structure on $\tilde{\mathcal{K}}$ has trivial connection and varying Riemannian metric on each fiber that, together with the previous fact, fits with [80, Theorem 3.3].

The existence of smooth solutions to (3.1.18) has important consequences for the existence and uniqueness problem for the coupled equations, as we will see in §3.4. In this sense, the solutions to (3.1.18) are the analogue in our context of the geodesics in the space of Kähler metrics in the cscK problem. We finish this section with an example of smooth solution to (3.1.18) defined for all $t \in [0, \infty[$, coming from the Lie algebra of the isotropy group $\tilde{\mathcal{G}}_I$ of the marked point $p_I \in \mathcal{P}$ on the extended gauge group $\tilde{\mathcal{G}}$ (cf. §1.3.1). It is very likely that a general existence result for short time, as the one for the geodesics in \mathcal{K} (see [48, Remark 3.3]), should hold for (3.1.18). Note however that, the lack of a uniform $0 < t_0 \in \mathbb{R}$ over $\tilde{\mathcal{K}}$ for which the solutions to (3.1.18) are defined does not allow one to construct a suitable map ρ as in (3.1.13). This can be considered an indirect proof of the fact that the extended gauge group $\tilde{\mathcal{G}}$ does not admit a complexification extending the $\tilde{\mathcal{G}}$ -action on the complex manifold \mathcal{P} .

Example 3.1.5. Let $\zeta \in \text{Lie } \tilde{\mathcal{G}}_I \subset \text{Lie } \tilde{\mathcal{G}}$, where we suppose fixed $(\omega, H) \in \tilde{\mathcal{K}}$. Then we can write

$$\zeta = \xi + \theta_A^\perp(\eta),$$

where $\xi \in \text{Lie } \mathcal{G}$ and $\eta \in \text{Lie } \mathcal{H}$ is the Hamiltonian vector field of $\phi \in C^\infty(X)$. Given any vector field y on X (resp. on E^c), we denote by f_t^y the flow of y on X (resp. on E^c). Let us define the curve

$$(\omega_t, H_t) = ((f_t^{J\eta})_*\omega, (f_t^{I\zeta})_*H),$$

where $I\zeta = \mathbf{i}\xi + \theta_A^\perp(J\eta)$. We claim that (ω_t, H_t) is a solution to $\frac{d}{dt}\zeta_t = 0$ in $\tilde{\mathcal{K}}$, where ζ_t is as in (3.1.11). Recall from Example 1.3.1 that $J\eta = \check{y}_t$ in (3.1.9). Combined with

$$\dot{H}_t = (f_t^{I\zeta})_*(-2\pi_1\theta_H(I\zeta)) = -2(f_t^{I\zeta})_*(\mathbf{i}\xi),$$

it implies that

$$\begin{aligned} y_t &= (f_t^{I\zeta})_*(\mathbf{i}\xi) + \theta_t^\perp(J\eta) \\ &= (f_t^{I\zeta})_*(\mathbf{i}\xi + \theta_{f_t^{-I\zeta}\cdot A}^\perp(J\eta)) \\ &= (f_t^{I\zeta})_*I\zeta = I\zeta, \end{aligned}$$

and so $\zeta_t = f_t^*(Iy_t) = -\zeta$, since $I\zeta$ is holomorphic.

Remark 3.1.6. For further computations it is convenient to express the time dependent vector fields y_t and ζ_t , defined respectively in (3.1.9) and (3.1.11), in a different way. Given a curve $(\omega_t, H_t) \in \tilde{\mathcal{K}}$ starting at the fixed point $\kappa = (\omega, H) \in \tilde{\mathcal{K}}$, we can suppose that $H_t = s_t^*H$ for a smooth curve $s_t = e^{i\xi_t}$, where $\xi_t \subset \mathbf{i}\text{Lie } \mathcal{G}$. Recall from (3.1.17) that the H_t -symmetric tangent vector \dot{H}_t is given by

$$\dot{H}_t = 2s_t^*(\pi_1y_{s_t}) \in \mathbf{i}\text{Lie } \mathcal{G}_t,$$

where $y_{s_t} = \dot{s}_t s_t^{-1}$. Setting $\omega_t = \omega + dd^c\phi_t$, it follows from (3.1.9) that

$$\begin{aligned} (3.1.19) \quad y_t &= -\pi_1^t y_{s_t} + \theta_t^\perp(\check{y}_t) \\ &= s_t^*(-\pi_1y_{s_t} + \theta_{s_t\cdot A}^\perp(\check{y}_t)) \\ &= -Is_t^*(-\mathbf{i}\pi_1y_{s_t} + \theta_{s_t A}^\perp(J\check{y}_t)), \end{aligned}$$

where $\pi_1^t: \text{Lie } \mathcal{G}^c \rightarrow \mathbf{i}\text{Lie } \mathcal{G}_t$ is the standard projection of H_t and $p_I = (J, A)$ is the marked point associated to κ (3.1.2). Let us define

$$(3.1.20) \quad y_{g_t} = \pi_0(y_{s_t}) + \theta_{s_t\cdot A}^\perp(\check{y}_t).$$

Note that y_{g_t} preserves H_0 . From (3.1.19) it follows that $y_t = s_t^*(-y_{s_t} + y_{g_t})$ and so the lift equals $f_t = s_t^{-1}g_t$. Finally, this gives

$$(3.1.21) \quad \zeta_t = g_t^*(-\mathbf{i}\pi_1y_{s_t} + \theta_{s_t\cdot A}^\perp(y_{\dot{\phi}_t, \omega_t})) = g_t^*(-\mathbf{i}\pi_1y_{s_t}) + \theta_{f_t^{-1}\cdot A}^\perp(y_{\dot{\phi}_t \circ \check{f}_t}).$$

3.2. Deforming solutions in a complexified orbit

In this section we perturb a solution to the coupled equations (3.0.7) with coupling constants α_0, α_1 to new solutions for close values of the coupling constants. In order to do this we need to impose conditions on the spaces of holomorphic vector fields of the manifold and of the total space of the bundle.

Let us first explain the underlying geometric idea of our deformation argument. Recall that (Z, J, ω) , $G \subset G^c$ are a Kähler manifold, a compact group acting by

holomorphic isometries on Z and its complexification. Pick $z \in Z$, a zero of the moment map $\mu : Z \rightarrow \mathfrak{g}^*$, and consider the map

$$(3.2.1) \quad L : \mathfrak{g} \rightarrow \mathfrak{g}^* : \zeta \mapsto L(\zeta) := \mu(e^{i\zeta} \cdot z).$$

The kernel of $dL|_0$ is the Lie algebra of the isotropy group G_z of z in G , due to

$$\langle dL|_0(\zeta), \zeta' \rangle = \omega(Y_\zeta, JY_{\zeta'})|_z$$

for any $\zeta, \zeta' \in \mathfrak{g}$, where Y_ζ denotes the infinitesimal action of ζ on Z . Therefore, when G_z is finite, $dL|_0$ is an isomorphism. In the situation of our interest the moment map μ_α is parameterized by a pair of coupling constants $\alpha_0, \alpha_1 \in \mathbb{R}$. When G_z is finite the initial zero $z \in Z$ of the moment map μ_α can be perturbed, via the Implicit Function Theorem, to zeros of $\mu_{\alpha'}$ for closed values α' of the coupling constants.

We now apply this picture to our infinite dimensional setup. First we will construct an analogue of (3.2.1) using the map ρ (3.1.13) defined in §3.1. We keep the notation of the previous section for a holomorphic principal G^c -bundle (E^c, I) over a compact complex manifold (X, J) . We fix a pair

$$\kappa = (\omega, H) \in \mathcal{K} \times \mathcal{R} = \tilde{\mathcal{K}}$$

consisting of a Kähler metric ω on the fixed Kähler class $[\omega]$ and a G -reduction H of the underlying smooth bundle E^c . For any pair $\alpha = (\alpha_0, \alpha_1) \in \mathbb{R}^2$ of positive coupling constants we have an associated mK-Hamiltonian space (see Definition 3.1.1)

$$\kappa \rightarrow (p_I, \mathcal{P}, \mathbf{I}, \omega_\alpha, \tilde{\mathcal{G}}, \mu_\alpha)$$

as in (3.1.2), where only the symplectic structure ω_α (2.3.1) and the $\tilde{\mathcal{G}}$ -equivariant moment map μ_α (2.3.3) depend on the coupling constants. Recall that (see (2.3.5))

$$\mathcal{P} \subset \mathcal{J}^i \times \mathcal{A}$$

is given by pairs $p = (J', A')$ such that $J' \in \mathcal{J}^i$, the space of compatible complex structures on (X, ω) , and A' is a $(1, 1)$ -type connection with respect to J' on the G -reduction H . We identify \mathcal{K} with the space of Kähler potentials $\phi \in C_0^\infty(X)$, with zero integral with respect to vol_ω . The construction of §3.1 gives a well defined map (see (3.1.13))

$$\rho : \mathcal{K} \times \mathbf{i} \text{Lie } \mathcal{G} \rightarrow \tau_{p_I}(\tilde{\mathcal{Y}}) \subset \mathcal{P},$$

where $\tau_{p_I}(\tilde{\mathcal{Y}})$ is the complexified $\tilde{\mathcal{G}}$ -orbit of the marked point p_I in the Kähler manifold given by $(\mathcal{P}, \mathbf{I}, \omega_\alpha)$. We will identify $\text{Lie } \mathcal{G} \cong \mathbf{i} \text{Lie } \mathcal{G}$ via the change of variable $\dot{H} = \mathbf{i} \xi$ in (3.1.13) to simplify the notation. Then, composing with the moment map μ_α , (3.1.13) provides a well defined map

$$L : \mathbb{R}^2 \times \mathcal{K} \times \text{Lie } \mathcal{G} \rightarrow C_0^\infty(X) \times \text{Lie } \mathcal{G}$$

given by

$$(3.2.2) \quad L(\alpha_0, \alpha_1, \phi, \xi) := (\alpha_0 S_{J_{\phi, \xi}} + \alpha_1 \Lambda^2(F_{A_{\phi, \xi}} \wedge F_{A_{\phi, \xi}}) - c, \Lambda F_{A_{\phi, \xi}}),$$

where $\rho(\phi, \xi) = (J_{\phi, \xi}, A_{\phi, \xi}) \in \mathcal{P}$ and $c \in \mathbb{R}$ is as in (3.0.7). Note that we have normalized the $\text{Lie } \mathcal{G}$ factor in (3.2.2). This allows us to perturb solutions also in the limit $\alpha_1 = 0$.

Remark 3.2.1. Here μ_α is regarded as a map

$$\mu_\alpha: \mathcal{P} \rightarrow \text{Lie } \mathcal{H} \times \text{Lie } \mathcal{G},$$

provided by the vector space splitting $\text{Lie } \tilde{\mathcal{G}} \cong \text{Lie } \mathcal{H} \times \text{Lie } \mathcal{G}$ and the pairing (2.2.15) given by any G -connection A on the reduction H .

Recall that \mathcal{J}^i denotes the space ω -compatible complex structures in X . Let

$$P: C_0^\infty(X) \rightarrow T_J \mathcal{J}^i$$

be the infinitesimal action of $C_0^\infty(X)$ on \mathcal{J}^i , where the former is regarded as the Lie algebra of the group of Hamiltonian sysmplectomorphisms on (X, ω) . Then it follows from (3.1.15), from the definition of the complex structure \mathbf{I} (2.3.4) and from the moment map property of the scalar curvature (see [22]) that the differential at the origin of L is given by

$$(3.2.3) \quad \begin{aligned} dL_{0|0}(\phi, \xi) &= \alpha_0 P^* P\phi + 2\alpha_1 \Lambda^2(F_A \wedge d_A J(d_A \xi + i_{\eta_\phi} F_A)), \\ dL_{1|0}(\phi, \xi) &= \Lambda d_A J(d_A \xi + i_{\eta_\phi} F_A) = d_A^*(d_A \xi + i_{\eta_\phi} F_A), \end{aligned}$$

where $L = (L_0, L_1)$, $p_I = (J, A)$ is the marked point in \mathcal{P} and P^* denotes the adjoint operator with respect to the metric on $\Omega^0(\text{End } TX)$ determined by J and ω . The kernel of the system (3.2.3) has a geometric interpretation in terms of holomorphic G^c -invariant vector fields in E^c when α_0 and α_1 have the same sign (cf. Proposition 1 in [44]). Given $(\phi, \xi) \in C_0^\infty(X) \times \text{Lie } \mathcal{G}$ we will denote

$$(3.2.4) \quad \zeta_{\phi, \xi} = \xi + \theta_A^\perp(\eta_\phi) \in \text{Lie } \tilde{\mathcal{G}}$$

as in (3.1.14), where η_ϕ is the Hamiltonian vector field on (X, ω) corresponding to ϕ and $\theta_A^\perp(\eta_\phi)$ denotes its horizontal lift with respect to the connection A determined by the marked point $p_I = (J, A) \in \mathcal{P}$.

Lemma 3.2.2. *If $\alpha_0 \cdot \alpha_1 > 0$ and $\Lambda_\omega F_A = z$, then a pair $(\phi, \xi) \in C_0^\infty(X) \times \text{Lie } \mathcal{G}$ is in the kernel of (3.2.3) if and only if $\zeta_{\phi, \xi}$ (3.2.4) is a G^c -invariant real holomorphic vector field on (E^c, I) .*

PROOF. Note that the real G^c -invariant vector field (3.2.4) is holomorphic if and only if

$$P\phi = 0, \quad d_A \xi + i_{\eta_\phi} F_A = 0,$$

or, equivalently, if its infinitesimal action $(Y_{\zeta_{\phi, \xi}})_{|p_I}$ at the point $p_I \in \mathcal{P}$ vanishes (see (2.2.9)). Hence, if (3.2.4) is holomorphic, $(\phi, \xi) \in \text{Ker } dL_{0|0}$, by (3.2.3). To prove the converse we fix $(\alpha_0, \alpha_1) \in \mathbb{R}^2$ and consider the pairings

$$\begin{aligned} \langle \cdot, \cdot \rangle_{\mathcal{H}}: \text{Lie } \mathcal{H} \times \text{Lie } \mathcal{H} &\rightarrow \mathbb{R}, \\ \langle \cdot, \cdot \rangle_{\mathcal{G}}: \text{Lie } \mathcal{G} \times \text{Lie } \mathcal{G} &\rightarrow \mathbb{R}, \end{aligned}$$

defined in (2.2.14) and (1.2.3). Then, given $(\phi_j, \xi_j) \in C_0^\infty(X) \times \text{Lie } \mathcal{G}$ with $j = 0, 1$, the differential $dL = (dL_0, dL_1)$ satisfies

$$\begin{aligned} \langle dL_{0|0}(\phi_0, \xi_0), \phi_1 \rangle_{\mathcal{H}} + \alpha_1 \cdot \langle dL_{1|0}(\phi_0, \xi_0), \xi_1 \rangle_{\mathcal{G}} &= \langle d\mu_\alpha(\mathbf{I}Y_{\zeta_0}), \zeta_1 \rangle_{|p} \\ &= \omega_\alpha(Y_{\zeta_1}, \mathbf{I}Y_{\zeta_0})_{|p}, \end{aligned}$$

where $\zeta_j = \zeta_{\phi_j, \xi_j} \in \text{Lie } \tilde{\mathcal{G}}$ is as in (3.2.4). Recall that the symplectic structure ω_α of (2.3.1) and the moment map μ_α of (2.3.3) (where we set $z = 0$) depend on the coupling constants α_0 and α_1 . If $\alpha_0 \cdot \alpha_1 > 0$ then ω_α is compatible with either \mathbf{I} or $-\mathbf{I}$ and so the previous equation implies that if $(\phi_0, \xi_0) \in \text{Ker } dL|_0$ then $Y_{\zeta_0} = 0$ (setting $\zeta_0 = \zeta_1$). We conclude that the vector field ζ_0 is holomorphic as claimed. \square

Note that when $\alpha_0 \cdot \alpha_1 > 0$, the kernel of $dL|_0$ can be identified with the Lie algebra of the isotropy group $\tilde{\mathcal{G}}_I$ of $p_I \in \mathcal{P}$ inside the extended gauge group $\tilde{\mathcal{G}}$. Let $H^0(TX)$ be the space of holomorphic sections of TX with respect to complex structure J . Taking suitable Sobolev completions of $C_0^\infty(X) \times \text{Lie } \mathcal{G}$ and applying the Implicit Function Theorem we obtain the following result.

Proposition 3.2.3. *Suppose that $\kappa = (\omega, H) \in \tilde{\mathcal{K}}$ is a solution to the coupled equations (3.0.7) with coupling constants $(\alpha_0, \alpha_1) \in \mathbb{R}^2$. Let $(p_I, \mathcal{P}, \mathbf{I}, \omega_\alpha, \tilde{\mathcal{G}}, \mu_\alpha)$ be the associated mK -Hamiltonian space and $\mathfrak{z} \subset \text{Lie } \tilde{\mathcal{G}}_I$ the centre of the Lie algebra of the compact group $G \subset G^c$. Then,*

- i) *If $\text{Lie } \tilde{\mathcal{G}}_I = \mathfrak{z}$ and $\alpha_0 \cdot \alpha_1 > 0$, then there exists an open neighbourhood U of (α_0, α_1) in \mathbb{R}^2 and a C^1 map f from U to the complexified orbit of the marked point p_I such that $f(t, s)$ is locally the unique solution of the coupled equations (6.0.2) with coupling constants (t, s) over this complexified orbit modulo the action of $\tilde{\mathcal{G}}$.*
- ii) *If $\text{Lie } \tilde{\mathcal{G}}_I = \mathfrak{z}$, $H^0(TX) = 0$, $\alpha_0 \neq 0$ and $\alpha_1 = 0$, then there exists an open neighbourhood U of $(\alpha_0, 0)$ with the same properties as in i).*

PROOF. Given a positive integer $k > 0$, let $C^k(X)$ and $\Omega^{0,k}(X, \text{ad } E)$ be, respectively, the spaces of C^k functions and sections of X and $\text{ad } E$ and denote by H^k and W^k their L^2 completions. Let

$$(3.2.5) \quad L_k: U_k \subseteq H^k \times W^{k-1} \rightarrow H^{k-4} \times W^{k-3}$$

be the map induced by L (3.2.2) in a neighbourhood U_k of the origin. Since the map L of (3.2.2) is not explicit, we check that it is well defined. We take $k \gg 0$ such that H^k is an algebra and all the equations below hold in a strong sense. Given $(\phi, \xi) \in U_k$ the automorphism $f_1 \in \text{Aut}(E^c)$ involved in the definition of the map (see (3.1.13))

$$\rho(\phi, \xi) = f_1^* p$$

has $k - 2$ derivatives in L^2 , due to the formula (3.1.16) and the Picard's Theorem on ordinary differential equations for Banach spaces (see [50, Theorem 3.1]). Recall from Remark 3.1.6 that we have a decomposition $f_1 = s^{-1} \cdot g$ where s and g have, respectively, $k - 1$ and $k - 2$ derivatives in L^2 . Therefore, the equality

$$F_{A_{\phi, \xi}} = F_{f_1^* A} = g^* F_{s \cdot A}$$

implies that the terms $\Lambda F_{A_{\phi, \xi}}$ and $\Lambda^2(F_{A_{\phi, \xi}} \wedge F_{A_{\phi, \xi}})$ in (3.2.2) are respectively in W^{k-3} and H^{k-3} (note that the symplectic form is fixed and so Λ is a C^∞ operator). On the other hand, if \check{f}_1 is the diffeomorphism on X that f_1 covers, then

$$S_{J_{\phi, \xi}} = S_{\check{f}_1^* J} = S_{\omega_\phi} \circ \check{f}_1,$$

in (3.2.2), where S_{ω_ϕ} denotes the scalar curvature of the H^{k-2} metric provided by the complex structure J and the H^{k-2} symplectic form $\omega_\phi = \omega + dd^c\phi$. Therefore, the function $S_{J_{\phi,\xi}}$ is of class H^{k-4} . Hence, the previous computation shows that

$$L_k(\phi, \xi) \subset H^{k-4} \times W^{k-3}$$

for k large enough, as required, provided that $(\phi, \xi) \in H^k \times W^{k-1}$.

We now proof the statement using a deformation argument. For simplicity we suppose $\mathfrak{z} = \{0\}$, i.e. \mathfrak{g} is semi-simple, and we explain later how to adapt the argument. The linearization of L_k at the origin in the two cases *i*) and *ii*) is given by (3.2.3) and it is invertible by hypothesis, provided that the induced operators $P^*P: H^k \rightarrow H^{k-4}$ and

$$\Delta_A = d_A^* d_A: W^{k-1} \rightarrow W^{k-3}$$

are self-adjoint elliptic operators. The result follows applying the Implicit Function Theorem in Banach spaces (see e.g. [8, p. 72, Theorem 3.10]). The regularity of the solutions follows, by local uniqueness, taking arbitrary large values of k . If \mathfrak{g} is not semi-simple the argument can be adapted by changing W^k and W^{k-2} in (3.2.5) by the orthogonal complement \mathfrak{z}^\perp of \mathfrak{z} on each space and considering the map induced by (3.2.2) and the orthogonal projection. The same argument holds since this restricted map has the same linearization as (3.2.2) at the origin. \square

As a straightforward consequence of the previous result we obtain,

Theorem 3.2.4. *Let (X, L) be a polarised complex projective manifold, G^c be a complex reductive Lie group and E^c be a holomorphic G^c -bundle over X . If there exists a cscK metric $\omega \in c_1(L)$, X has finite automorphism group and E^c is stable with respect to L then, given a pair of positive real constants $\alpha_0, \alpha_1 > 0$ with small ratio $0 < \frac{\alpha_1}{\alpha_0} \ll 1$, there exists a solution $(\omega_\alpha, H_\alpha)$ to (3.0.7) with these coupling constants and $\omega_\alpha \in c_1(L)$.*

PROOF. Let $\omega \in c_1(L)$ be a cscK metric. By the Hitchin-Kobayashi correspondence for the HYM equation [21, 79, 62], E^c admits a reduction H to a maximal compact subgroup $G \subset G^c$ with HYM Chern connection A . The pair (ω, H) is then a solution to the coupled equation (3.0.7) with $\alpha_0 \neq 0, \alpha_1 = 0$. Let $\tilde{\mathcal{G}}$ be the extended gauge group (2.1.1) associated to the symplectic form ω and the reduction H . We claim that the Lie algebra of the subgroup $\tilde{\mathcal{G}}_I \subset \tilde{\mathcal{G}}$ preserving the holomorphic structure in E^c reduces to the centre \mathfrak{z} of the Lie algebra of G . To see this, consider any G -invariant (real) holomorphic vector field $\zeta \in \text{Lie } \tilde{\mathcal{G}}_I$ and let $\eta \in \Omega^0(TX)$ be its projection onto X . By definition of $\tilde{\mathcal{G}}_I$ (see also (3.2.2)), η is a Hamiltonian Killing vector field that defines a 1-parameter group of isometries of (X, ω) and hence, by hypothesis, identically vanishes. The vector field ζ is then an A -parallel vertical vector field and can be identified with a holomorphic G^c -invariant vector field on E^c . Since E^c is stable, the Lie algebra of holomorphic G^c -invariant vector fields, canonically identified with $H^0(X, \text{ad } E^c)$, is isomorphic to the centre \mathfrak{z}^c of G^c . Finally, since ζ preserves the reduction H , we conclude that $\zeta \in \mathfrak{z} \subset \mathfrak{z}^c$, as claimed. In the situation considered, we are under the hypothesis of Proposition 3.2.3 and the statement holds. \square

Remark 3.2.5. The hypothesis of Theorem 3.2.4 are satisfied in a large class of examples. In §5.3 we will discuss the special case in which the polarized manifold (X, L) satisfies $c_1(L) = \lambda \cdot c_1(X)$, for $\lambda \in \mathbb{Z}$.

Remark 3.2.6. The statement of Theorem 3.2.4 is also true for arbitrary Kähler manifolds, not necessarily algebraic, if we impose the additional condition of simplicity on (E^c, I) (see [79, p. 261]).

3.3. Integral invariants

Let (E^c, I) be a holomorphic G^c -bundle over a compact complex manifold (X, J) . We fix a Kähler class $[\omega]$ on X and denote by $\tilde{\mathcal{K}} = \mathcal{K} \times \mathcal{R}$ the space of pairs (ω, H) consisting of a Kähler metric $\omega \in [\omega]$ and a smooth G -reduction $H \in \mathcal{R}$. We consider the coupled equations (3.0.7) for elements in $\tilde{\mathcal{K}}$, with fixed positive coupling constants

$$0 < \alpha_0, \alpha_1 \in \mathbb{R}.$$

In this section we give an obstruction to the existence of solutions to (3.0.7) controlled by a character of the Lie algebra of $\text{Aut}(E^c, I)$, the group of automorphisms of the holomorphic bundle over X . This generalizes previous constructions of A. Futaki in [29, 30]. To define the character we adapt the argument of J. P. Bourguignon in [14]. To simplify some formulae we suppose $z = 0$ in the system (3.0.7), as there is no significant difference with the general case.

To define the character we proceed as in §1.1.2, so we define first a closed 1-form σ_I in $\tilde{\mathcal{K}}$. To do this we need a technical lemma that relates σ_I with the moment map μ_α whose zero locus correspond to the solutions to the coupled equations (3.0.7). We will use this lemma to prove the analogue of Proposition 1.3.2, about suitable properties of the 1-form necessary for the definition of the invariant. For a first reading one can skip Lemma 3.3.1 and the proof of Proposition 3.3.2.

Let us first define the 1-form. Fix $\kappa = (\omega, H) \in \tilde{\mathcal{K}}$. As in §3.1, we identify the tangent space

$$T_\kappa \tilde{\mathcal{K}} \cong C_0^\infty(X) \times \mathbf{i} \text{Lie } \mathcal{G},$$

where $C_0^\infty(X)$ are smooth functions with zero integral with respect to vol_ω and $\text{Lie } \mathcal{G}$ is the gauge group of the G -reduction H . Recall from (3.1.12) that given $(\dot{\phi}, \dot{H}) \in T_\kappa \tilde{\mathcal{K}}$, with the previous identification, we consider it as the vector given by the curve

$$(3.3.1) \quad \kappa_t = (\omega_t, H_t) := (\omega_{t\dot{\phi}}, (e^{t\dot{H}})^* H) \subset \tilde{\mathcal{K}},$$

where

$$\omega_{t\dot{\phi}} = \omega + tdd^c \dot{\phi}$$

and e denotes the exponential map in the complex gauge group \mathcal{G}^c of the smooth G^c -bundle E^c . Given $\dot{\kappa} = (\dot{\phi}, \dot{H}) \in T_\kappa \tilde{\mathcal{K}}$, we define the 1-form σ_I at $\kappa \in \tilde{\mathcal{K}}$ as

$$(3.3.2) \quad \sigma_I(\dot{\phi}, \dot{H}) = - \int_X (\dot{\phi}(\alpha_0 S_\omega + \alpha_1 \Lambda_\omega^2(F_H \wedge F_H) - c) + 2\alpha_1 (\mathbf{i} \dot{H}, \Lambda_\omega F_H)) \frac{\omega^n}{n!},$$

where $c \in \mathbb{R}$ is as in (3.0.7). The relation of the 1-form (3.3.2) with the moment map of Section §2.3 is given by the following result, that we will use systematically in the following.

Lemma 3.3.1. Fix $\kappa \in \tilde{\mathcal{K}}$ and consider $(p_I, \mathcal{P}, \mathbf{I}, \omega_\alpha, \tilde{\mathcal{G}}, \mu_\alpha)$, the associated mK -Hamiltonian space, where the $\tilde{\mathcal{G}}$ -equivariant moment map μ_α is determined by the coupling constants α_0, α_1 . Then, given any curve κ_t in $\tilde{\mathcal{K}}$ with initial condition κ , we have

$$\sigma_I(\dot{\kappa}_t) = \langle \mu_\alpha(f_t^{-1} \cdot p), f_t^*(Iy_t) \rangle,$$

where $f_t \subset \text{Aut}(E^c)$ is the horizontal lift of κ_t defined in (3.1.9).

PROOF. Let $\kappa = (\omega, H)$ and consider a curve $\kappa_t = (\omega_t, H_t) = (\omega_{\phi_t}, s_t^*H)$ in $\tilde{\mathcal{K}}$ with initial condition κ , $s_t = e^{\xi_t}$ and $\xi_t \in \mathbf{i}\text{Lie}\mathcal{G}$. Consider the horizontal lift f_t of this curve through $f_0 = \text{Id}$ defined by (3.1.9). Recall that

$$(\check{f}_t^*\omega_t, f_t^*H_t) = (\omega, H)$$

and also that $f_t = s_t^{-1}g_t$, where g_t preserves the fixed G -reduction H on E^c , for all t . Define the following curves:

$$\omega_t = \omega_{\phi_t}, \quad A_t = f_t^{-1} \cdot A, \quad J_t = \check{f}_t^*J, \quad \text{and} \quad \eta_t,$$

where the latter is the Hamiltonian vector field of the function $\dot{\phi}_t \circ \check{f}_t$ on X with respect to the fixed symplectic form ω . Recall from Remark 3.1.6 that

$$\dot{H}_t = 2s_t^*(\pi_1y_{s_t}) \in \mathbf{i}\text{Lie}\mathcal{G}_t,$$

where $y_{s_t} = \dot{s}_t s_t^{-1}$ and π_1 is, as in §3.1, the projection onto $\mathbf{i}\text{Lie}\mathcal{G}$. Combined with the formula $F_{H_t} = s_t^*(F_{s_t A})$ (see [21, Section 1.1]), we have

$$\begin{aligned} \sigma_I(\dot{\phi}_t, \dot{H}_t) &= - \int_X \dot{\phi}_t (\alpha_0 S_{\omega_t} + \alpha_1 \Lambda_{\omega_t}^2(F_{s_t A}^2) - c) + 4\alpha_1 (\mathbf{i}\pi_1y_{s_t}, \Lambda_{\omega_t} F_{s_t A}) \frac{\omega_t^n}{n!} \\ &= - \int_X \dot{\phi}_t \circ \check{f}_t (\alpha_0 S_{J_t} + \alpha_1 \Lambda_{\omega}^2 F_{A_t}^2 - c) + 4\alpha_1 (\mathbf{i}\pi_1g_t^*y_{s_t}, \Lambda_{\omega} F_{A_t}) \frac{\omega^n}{n!} \\ &= \langle \mu_\alpha(J_t, A_t), -\mathbf{i}\pi_1g_t^*(y_{s_t}) + \theta_{A_t}^\perp(\eta_{\dot{\phi}_t \circ \check{f}_t}) \rangle \\ &= \langle \mu_\alpha(f_t^{-1} \cdot p), f_t^*(Iy_t) \rangle, \end{aligned}$$

where $\hat{c} \in \mathbb{R}$ is defined by (2.2.6). To do the changes of variable we have used the fact that the pairing $\Omega^p(\text{ad } E^c) \times \Omega^q(\text{ad } E^c) \rightarrow \Omega^{p+q} \otimes \mathbb{C}$ induced by (3.0.8) is \mathcal{G}^c -invariant. The last equality follows from the identity (3.1.21). \square

Note that $\tilde{\mathcal{K}}$ has a right action by pull-back of the group of holomorphic automorphisms $\text{Aut}(E^c, I)$. Then, the 1-form defined in (3.3.2) satisfies the following.

Proposition 3.3.2. σ_I is $\text{Aut}(E^c, I)$ -invariant and closed.

PROOF. To prove the invariance note that for any pair $(\omega, H) \in \tilde{\mathcal{K}}$, any Kähler potential ϕ and any $f \in \text{Aut}(E^c, I)$ we have $F_{f^*H} = f^*F_H$ and

$$\check{f}^*\omega_\phi = (\check{f}^*\omega)_{\check{f}^*\phi}.$$

Then, the $\text{Aut}(E^c, I)$ -invariance follows, after a change of variable in (3.3.2), from the \mathcal{G}^c -invariance of the pairing $\Omega^p(\text{ad } E^c) \times \Omega^q(\text{ad } E^c) \rightarrow \Omega^{p+q} \otimes \mathbb{C}$ induced by (3.0.8). To prove the second part of the statement, fix $\kappa \in \tilde{\mathcal{K}}$ and consider

$$(p_I, \mathcal{P}, \mathbf{I}, \omega_\alpha, \tilde{\mathcal{G}}, \mu_\alpha),$$

the associated mK-Hamiltonian space, where the $\tilde{\mathcal{G}}$ -equivariant moment map μ_α is determined by the coupling constants α_0, α_1 . We will prove that σ_I is closed at the arbitrary fixed point κ . Given

$$\dot{\kappa}_j = (\dot{\phi}_j, \dot{H}_j) \in C_0^\infty(X) \times \mathbf{i} \operatorname{Lie} \mathcal{G}, \quad j = 1, 2,$$

regarded as a vector field on $\tilde{\mathcal{K}}$, the Lie bracket $[\dot{\kappa}_1, \dot{\kappa}_2]$ vanishes at κ and so, combined with (1.1.6), it is enough to prove that $\dot{\kappa}_1(\sigma_I(\dot{\kappa}_2))$ is symmetric in $j = 1, 2$ at κ . Let

$$\kappa_t = (\omega_{t\phi_1}, s_t^* H),$$

where $s_t = e^{t\dot{H}_1}$. Then, using the same notation as in the previous lemma, we obtain

$$\begin{aligned} \dot{\kappa}_1(\sigma_I(\dot{\kappa}_2)) &= \frac{d}{dt}_{|t=0} \sigma_I(\dot{\kappa}_2)_{|\kappa_t} \\ &= \frac{d}{dt}_{|t=0} \langle \mu_\alpha(g_t^{-1}s_t \cdot p_I), g_t^*(-\mathbf{i} \dot{H}_2) + \theta_A^{\perp -1} s_t A(\eta_{\dot{\phi}_2 \circ \tilde{f}_t}) \rangle = \\ &= \omega_\alpha(Y_2, \mathbf{I} Y_1) + \langle \mu_\alpha(p_I), \dot{\kappa}_{2,1} \rangle, \end{aligned}$$

where ω_α is a symplectic form on \mathcal{P} compatible with \mathbf{I} , Y_j is the infinitesimal action of the vector field $-\mathbf{i} \dot{H}_j + \theta_A^{\perp -1}(\eta_{\dot{\phi}_j}) \in \operatorname{Lie} \tilde{\mathcal{G}}$ at $p_I \in \mathcal{P}$ and

$$\begin{aligned} \dot{\kappa}_{2,1} &= \frac{d}{dt}_{|t=0} g_t^*(-\mathbf{i} \dot{H}_2) + \theta_A^{\perp -1} s_t A(\eta_{\dot{\phi}_2 \circ \tilde{f}_t}) \\ &= \theta_A[\theta_A^\perp(-J\eta_{\dot{\phi}_1}), -\mathbf{i} \dot{H}_2 + \theta_A^\perp(\eta_{\dot{\phi}_2})] - [\mathbf{i} \dot{H}_1, \theta_A^\perp(J\eta_{\dot{\phi}_2})] - \theta_A^\perp(\eta_{\omega(\eta_{\dot{\phi}_2}, J\eta_{\dot{\phi}_1})}) \\ &= F_A(J\eta_{\dot{\phi}_1}, \eta_{\dot{\phi}_2}) + d_A(\mathbf{i} \dot{H}_1)(J\eta_{\dot{\phi}_2}) + d_A(\mathbf{i} \dot{H}_2)(J\eta_{\dot{\phi}_1}) - \theta_A^\perp(\eta_{\omega(\eta_{\dot{\phi}_2}, J\eta_{\dot{\phi}_1})}). \end{aligned}$$

Hence, $\dot{\kappa}_1(\sigma_p(\dot{\kappa}_2)) = \dot{\kappa}_2(\sigma_p(\dot{\kappa}_1))$ and so $d\sigma_I = 0$ holds. \square

Due to Proposition 3.3.2, the holomorphic structure I and the Kähler class $[\omega]$ determine, as in §1.1.2, a complex character

$$\mathcal{F}_{\alpha, I, [\omega]}: \operatorname{Lie} \operatorname{Aut}(E^c, I) \rightarrow \mathbb{C}: \zeta \rightarrow \mathbf{i} \sigma_I(\zeta^l) + \sigma_I(I\zeta)^l,$$

where ζ^l denotes the infinitesimal action of $\zeta \in \operatorname{Lie} \operatorname{Aut}(E^c, I)$ on $\tilde{\mathcal{K}}$. When there is no possible confusion we write simply \mathcal{F}_α for the previous character. To give an explicit expression for \mathcal{F}_α , recall from §1.3.4 that, given a Kähler form $\omega \in \mathcal{K}$, any holomorphic vector field $y \in \operatorname{Lie} \operatorname{Aut}(X, J)$ can be written as

$$\eta = \eta_{\phi_1} + J\eta_{\phi_2} + \beta,$$

where η_{ϕ_j} is the ω -Hamiltonian associated to ϕ_j , $j = 1, 2$, and β is the dual of an harmonic 1-form with respect to the Kähler structure (X, J, ω) . Then the infinitesimal action of $\zeta \in \operatorname{Lie} \operatorname{Aut}(E^c, I)$ at $(\omega, H) \in \tilde{\mathcal{K}}$ is

$$\zeta^l = (\phi_2, 2\pi_1 \theta_H(\zeta))$$

if ζ covers y and the value of the character \mathcal{F}_α at ζ can be written as

$$(3.3.3) \quad \mathcal{F}_\alpha(\zeta) = - \int_X (\phi(\alpha_0 S_\omega + \alpha_1 \Lambda_\omega^2(F_H \wedge F_H) - c) - 4\alpha_1(\theta_H \zeta, \Lambda_\omega F_H)) \frac{\omega^n}{n!},$$

where $\phi = \phi_1 + \mathbf{i} \phi_2$. The following result is the analogue of Proposition 1.1.5.

Proposition 3.3.3. *The map $\mathcal{F}_\alpha: \text{Lie Aut}(E^c, I) \rightarrow \mathbb{C}$ is independent of the $(\omega, H) \in \tilde{\mathcal{K}}$. It defines a character of $\text{Lie Aut}(E^c, I)$ that vanishes if $\tilde{\mathcal{K}}$ contains a solution to the coupled equations (3.0.7).*

PROOF. Formally the proof is the same as the proof of Proposition 1.1.5. \square

When $\alpha_1 = 0$ the expression (3.3.3) is, up to multiplicative constant factor, the *Futaki invariant* of the complex manifold (X, J) and the Kähler class $[\omega]$. When $\alpha_0 = 0$ the character \mathcal{F}_α generalizes an invariant due also to A. Futaki (see Theorem 1.1 in [30]), who considered the smaller Lie algebra of G^c -invariant holomorphic vector fields on (E^c, I) that project onto a vector field that vanishes somewhere on X (i.e. onto a holomorphic complex Hamiltonian vector field).

We finish this section with an example of how the invariant \mathcal{F}_α controls the existence of solutions to the *coupled equations* for small ratio of the coupling constants.

Example 3.3.4. Let (X, L) be a polarized manifold that admits an extremal, non cscK, metric ω in $c_1(L)$ (e.g. \mathbb{CP}^2 blown up in one point with the anticanonical polarization [15]). Recall that this means that the scalar curvature S_ω is the Hamiltonian function of a real holomorphic Killing vector field η . The classical Futaki invariant of the Kähler class $c_1(L)$ (take $\alpha_0 = 1$ and $\alpha_1 = 0$ in (3.3.3)) evaluated on η is

$$\mathcal{F}(\eta) = - \int_X (S_\omega - \hat{S})^2 \text{vol}_\omega < 0.$$

Let E^c be the holomorphic \mathbb{C}^* -principal bundle of frames of L . The holomorphic vector field η can be trivially lifted to a holomorphic vector field on E^c and it follows by definition that the associated character \mathcal{F}_α given by (3.3.3) evaluated at η is negative for sufficiently small values of $\alpha_1/\alpha_0 > 0$. Hence, the triple (X, L, E^c) does not admit a solution to the coupled system (3.0.7) for these values of the coupling constants. Given an arbitrary holomorphic principal bundle E^c over X , the obstruction to lift the holomorphic vector field η on X to a G^c -invariant one in E^c lies in $H^1(X, \text{ad } E^c)$ (cf. equation (2.2.9)). When this obstruction vanishes the same argument can be applied.

3.4. The integral of the moment map

In this section we show that there exists a suitable integral of the moment map for the coupled equations that extends the Donaldson functional for the HYM equation and the Mabuchi K-Energy for the cscK equation. We determine a sufficient condition for its convexity and the existence of lower bounds in terms of the existence of smooth solutions to certain partial differential equation on $\tilde{\mathcal{K}}$ that generalizes the geodesic equation on \mathcal{K} . When this condition holds, we generalize the obstruction defined in §3.3. Throughout this section we will consider the coupled equations (3.0.7) with fixed positive coupling constants $0 < \alpha_0, \alpha_1 \in \mathbb{R}$.

As in the previous section, let (E^c, I) be a holomorphic G^c -bundle over a compact complex manifold (X, J) . We fix a Kähler class $[\omega]$ on X and denote by $\tilde{\mathcal{K}} = \mathcal{K} \times \mathcal{R}$ the space of pairs (ω, H) consisting of a Kähler metric $\omega \in [\omega]$ and a smooth G -reduction

$H \in \mathcal{R}$. Since the space $\tilde{\mathcal{K}}$ is contractible, the 1-form σ_I given by (3.3.2) integrates, i.e there exists an integral of the moment map

$$(3.4.1) \quad \mathcal{M}_\kappa : \tilde{\mathcal{K}} \rightarrow \mathbb{R},$$

that satisfies $d\mathcal{M}_\kappa = \sigma_I$ and $\mathcal{M}_\kappa(\kappa) = 0$ for $\kappa = (\omega, H) \in \tilde{\mathcal{K}}$. This implies that the critical points of \mathcal{M}_κ are the solutions to the coupled equations in $\tilde{\mathcal{K}}$. Given any curve κ_t in $\tilde{\mathcal{K}}$, we can write

$$(3.4.2) \quad \mathcal{M}_\kappa(\kappa_t) = \mathcal{M}_\kappa(\kappa_0) + \int_0^t \sigma_I(\dot{\kappa}_s) \wedge ds.$$

Combined with Lemma 3.3.1 this implies the following.

Proposition 3.4.1. *The functional \mathcal{M}_κ given by (3.4.2) is convex along smooth solutions $\kappa_t = (\omega_{\phi_t}, H_t)$ of the system*

$$(3.4.3) \quad \left. \begin{aligned} \frac{1}{2} \ddot{H}_t - d_t \dot{H}_t(Jy_{\dot{\phi}_t, \omega_t}) - \mathbf{i} F_t(y_{\dot{\phi}_t, \omega_t}, Jy_{\dot{\phi}_t, \omega_t}) &= 0 \\ \ddot{\phi}_t - |d\dot{\phi}_t|_{\omega_t}^2 &= 0 \end{aligned} \right\},$$

where d_t and F_t denote, respectively, the covariant derivative and the curvature of the Chern connection determined by H_t .

PROOF. We fix $\kappa = (\omega, H) \in \tilde{\mathcal{K}}$ and consider the mK-Hamiltonian space (3.1.2) and the integral of the moment map \mathcal{M}_κ determined by κ . Given a smooth curve $\kappa_t = (\omega_t, H_t) \subset \tilde{\mathcal{K}}$ with $\omega_t = \omega_{\phi_t}$ let $f_t \subset \text{Aut}(E^c)$ be the associated lift defined by (3.1.9). Denote $\zeta_t = f_t^*(Iy_t) \in \text{Lie } \tilde{\mathcal{G}}$ as in Proposition 3.1.3. Then, applying (3.4.2) and Lemma 3.3.1 we have

$$(3.4.4) \quad \begin{aligned} \frac{d^2}{dt^2} \mathcal{M}_\kappa(\kappa_t) &= \frac{d}{dt} \sigma_I(\dot{\kappa}_t) \\ &= \frac{d}{dt} \langle \mu_\alpha(f_t^{-1} \cdot p), \zeta_t \rangle \\ &= \omega_\alpha(Y_{\zeta_t}, \mathbf{I}Y_{\zeta_t}) + \langle \mu_\alpha(f_t^{-1} \cdot p), \frac{d}{dt} \zeta_t \rangle, \end{aligned}$$

where $\omega_\alpha(Y_{\zeta_t}, \mathbf{I}Y_{\zeta_t}) = \|Y_{\zeta_t}\|_\alpha^2$ is the squared norm of the infinitesimal action of ζ_t in the Kähler manifold $(\mathcal{P}, \mathbf{I}, \omega_\alpha)$ associated to (ω, H) . Therefore, \mathcal{M}_κ is convex along the solutions to $\frac{d}{dt} \zeta_t = 0$ and strictly convex if Y_{ζ_t} do not vanishes. Note that if $Y_{\zeta_t} = 0$, then f_t preserves the complex structure I for all t and so the curve $\kappa_t = (f_t)_* \kappa_0$ is linked by a smooth curve of holomorphic automorphisms on $\text{Aut}(E^c, I)$. We now check that the system (3.4.3) is equivalent to the vanishing of $\frac{d}{dt} \zeta_t$. Note first that the latter

condition is equivalent to the vanishing of $\zeta'_t = (f_t)_* \frac{d}{dt} \zeta_t$, that we can write as

$$\begin{aligned}\zeta'_t &= [y_t, Iy_t] + I\dot{y}_t \\ &= [-\frac{1}{2}\dot{H}_t + \theta_t^\perp(\check{y}_t), -\frac{\mathbf{i}}{2}\dot{H}_t + \theta_t^\perp(J\check{y}_t)] \\ &\quad - \frac{\mathbf{i}}{2}\ddot{H}_t + (\frac{d}{dt}\theta_t^\perp)(J\check{y}_t) + \theta_t^\perp(\frac{d}{dt}J\check{y}_t) \\ &= [\theta_t^\perp(\check{y}_t), \theta_t^\perp(J\check{y}_t)] - \frac{\mathbf{i}}{2}d_t\dot{H}_t(\check{y}_t) + \frac{1}{2}d_t\dot{H}_t(J\check{y}_t) \\ &\quad - \frac{\mathbf{i}}{2}\ddot{H}_t + (\frac{d}{dt}\theta_t^\perp)(J\check{y}_t) + \theta_t^\perp(\frac{d}{dt}J\check{y}_t),\end{aligned}$$

where θ_t^\perp denotes the horizontal lift of vector fields with respect to the Chern connection determined by H_t . Denote respectively by η_t and y'_t the Hamiltonian vector field associated to the smooth function $\dot{\phi}_t \circ \check{f}_t$ with respect to ω , and $\phi'_t = \ddot{\phi}_t - |d\dot{\phi}_t|_{\omega_t}^2$ with respect to ω_t . Then we have that

$$\frac{d}{dt}J\check{y}_t = \frac{d}{dt}(\check{f}_t)_*(\eta_t) = -[\check{y}_t, J\check{y}_t] + y'_t,$$

due to the equality $J\check{y}_t = y_{\dot{\phi}_t, \omega_t}$, that combined with

$$(\frac{d}{dt}\theta_t^\perp)(J\check{y}_t) = -\partial_t(\dot{H}_t)(J\check{y}_t),$$

gives

$$\begin{aligned}\zeta'_t &= \theta_t[\theta_t^\perp(\check{y}_t), \theta_t^\perp(J\check{y}_t)] + (\bar{\partial}_t - \partial_t)\dot{H}_t(J\check{y}_t) - \frac{\mathbf{i}}{2}\ddot{H}_t + \theta_t^\perp(y'_t) \\ &= F_t(Jy_{\dot{\phi}_t, \omega_t}, y_{\dot{\phi}_t, \omega_t}) + \mathbf{i}d_t\dot{H}_t(Jy_{\dot{\phi}_t, \omega_t}) - \frac{\mathbf{i}}{2}\ddot{H}_t + \theta_t^\perp(y'_t).\end{aligned}$$

The statement follows by splitting ζ'_t into its A_t vertical and horizontal part for each value of t . \square

The existence of smooth solutions to (3.4.3) has important consequences for the existence and uniqueness problem for the coupled equations that we describe in the following proposition. Recall from Example 3.1.5 that any real holomorphic vector field in $\text{Lie } \tilde{\mathcal{G}}_I$ gives a smooth solution to (3.4.3) defined for all time, that we call a *trivial solution*. As mentioned in §3.1, it is very likely that a short time existence result as the one for the geodesics in \mathcal{K} (see [48, Remark 3.3] holds for (3.4.3). This would imply that the first of the following obstructions extends the one coming from the character \mathcal{F}_α defined in the previous section. Long time existence of smooth solutions is not even true in general for the geodesic equation in \mathcal{K} (see [19]).

Proposition 3.4.2.

- (1) If there exists a solution $\kappa \in \tilde{\mathcal{K}}$ to the coupled equations then for any smooth solution $\kappa_t \subset \tilde{\mathcal{K}}$ to (3.4.3), with $t \in [0, s[$, $s \in \mathbb{R} \cup \{\infty\}$ and initial condition $\kappa_0 = \kappa$, the following inequality holds

$$\lim_{t \rightarrow s} \sigma_I(\dot{\kappa}_t) \geq 0,$$

with equality only if κ_t is a trivial solution.

- (2) If any two points in $\tilde{\mathcal{K}}$ are joined by a smooth solution to (3.4.3), then there exists at most one solution to the coupled equations modulo the action of $\text{Aut}(E^c, I)$. Under the same hypothesis, if there exists one solution, then the integral of the moment map is bounded from below.

PROOF. We use the same notation as in Proposition 3.4.1. To prove 1) let κ_t , $t \in [0, s]$ with $s \in \mathbb{R} \cup \infty$ be a smooth solution to (3.4.3). Then, it follows from (3.4.4) that

$$\frac{d}{dt}\sigma_I(\dot{\kappa}_t) = \|Y_{\zeta_0}\|_\gamma^2 \geq 0,$$

where $\zeta_0 = \zeta_t$ for all t , since 3.4.1 is satisfied. If $\kappa_0 = \kappa$, a solution to the coupled equations, then $\sigma_I(\dot{\kappa}_0) = 0$ and so $\sigma_I(\dot{\kappa}_t) \geq 0$. If $\sigma_I(\dot{\kappa}_t) = 0$ for all t , then ζ_0 is holomorphic and κ_t is therefore a trivial solution. Thus, we have an equality

$$\lim_{t \rightarrow s} \sigma_I(\dot{\kappa}_t) = \sigma_I(\dot{\kappa}_0) = \mathcal{F}_\alpha(\zeta_0).$$

To prove 2), fix a solution $\kappa_0 \in \tilde{\mathcal{K}}$ to the coupled equations. Then, if κ_1 is another solution, it is linked to κ_0 , by hypothesis, by a smooth solution κ_t to 3.4.1, with $t \in [0, 1]$. It follows then from (3.4.4) and 1) that κ_t is a trivial solution since $\sigma_I(\dot{\kappa}_1) = \sigma_I(\dot{\kappa}_0)$. The lower bound of \mathcal{M}_κ for any $\kappa \in \tilde{\mathcal{K}}$ follows from 1) and (3.4.2). \square

3.5. Complex surfaces

In this section we restrict to the case in which the base X is a compact complex surface. We use the integral of the moment map constructed in the previous section to prove some results about the uniqueness and existence problem for the coupled equations (3.0.7) with fixed positive coupling constants $0 < \alpha_0, \alpha_1 \in \mathbb{R}$. In this section we suppose that $\mathfrak{z} \ni z = 0$ in (3.0.7).

Let (E^c, I) be a holomorphic G^c -bundle over a compact complex surface (X, J) . We fix a Kähler class $[\omega]$ on X and suppose that $(\omega, H) \in \tilde{\mathcal{K}}$ is a solution to (3.0.7). Consider the functional on the space \mathcal{K} , of Kähler metrics with fixed Kähler class $[\omega]$, given by

$$\mathcal{M}_\kappa(\cdot, H) : \mathcal{K} \rightarrow \mathbb{R},$$

where \mathcal{M}_κ is the integral of the moment map defined in (3.4.2) for a choice of $\kappa \in \tilde{\mathcal{K}}$. Without loss of generality we suppose that $\kappa = (\omega, H)$ and denote by \mathcal{M}_H the functional above. Then, along a curve $\omega_t = \omega_{\phi_t}$ in \mathcal{K} , we can write

$$(3.5.1) \quad \begin{aligned} \mathcal{M}_H(\omega_t) &= - \int_0^t \int_X (\dot{\phi}_s(\alpha_0 S_{\omega_s} + \alpha_1 \Lambda_{\omega_s}^2(F_H \wedge F_H) - c)) \frac{\omega_s^2}{2!} \wedge ds \\ &= - \int_0^t \int_X \dot{\phi}_s(\alpha_0 S_{\omega_s} - c) \frac{\omega_s^2}{2!} - 2\alpha_1 \int_0^t \int_X \dot{\phi}_s(F_H \wedge F_H) \wedge ds. \end{aligned}$$

The crucial fact for our discussion concerning complex surfaces is the following: since (ω, H) is a solution to (3.0.7) with $\mathfrak{z} \ni z = 0$, the Chern connection A_H is Anti-Self-Dual and so

$$-2(F_H \wedge F_H) = |F_H|^2 \cdot \omega^2,$$

that is, a non-negative real 4-form, where the point-wise norm $|F_H|^2$ is taken with respect to ω . Straightforward computations as in Lemma 3.3.1 show then that

$$(3.5.2) \quad \frac{d^2}{dt^2} \mathcal{M}_H(\omega_t) = 2\alpha_0 \|Y_{\eta_t}\|^2 - \int_X (\ddot{\phi}_t - |d\dot{\phi}_t|_{\omega_t}^2)(\alpha_0 S_{\omega_t} - c) \frac{\omega_t^2}{2!} + \alpha_1 \int_X \ddot{\phi}_t |F_H|^2 \cdot \omega^2,$$

where $\|Y_{\eta_t}\|^2$ denotes the squared norm of the infinitesimal action of the time dependent Hamiltonian vector field $\eta_t = \eta_{\dot{\phi}_t \circ f_t} \in \text{Lie } \mathcal{H}$ in the space of complex structures \mathcal{J}^i compatible with ω , where f_t is the Moser curve of ω_t as in (1.3.15). As an immediate consequence we obtain

Lemma 3.5.1. *If any point in \mathcal{K} is linked to ω by a smooth geodesic, then \mathcal{M}_H is bounded from below. For any other metric $\omega' \in \mathcal{K}$ such that (ω', H) is a solution to (3.0.7), there exists a holomorphic transformation $f \in \text{Aut}(X, J)$ such that $f^*\omega' = \omega$. Moreover, if the connection A_H is not flat, then ω is the unique metric in \mathcal{K} with this property.*

PROOF. The first part of the statement follows from $d\mathcal{M}_{H|\omega} = 0$ and (3.5.2), due to the inequality

$$(3.5.3) \quad \frac{d^2}{dt^2} \mathcal{M}_H(\omega_t) = 2\alpha_0 \|Y_{\eta_t}\|^2 + \alpha_1 \int_X |d\dot{\phi}_t|_{\omega_t}^2 |F_H|^2 \cdot \omega^2 \geq 0$$

along smooth geodesics. For the second part, suppose that ω_{ϕ_t} joins ω to another solution ω' of (3.0.7). Then, $\eta_t = \eta_0$ since ϕ_t is a geodesic in the space of Kähler potentials and $d\mathcal{M}_{H|\omega'} = 0$. This implies that (3.5.3) is an equality and that η_0 is a holomorphic vector field with f_t being its flow. We conclude that $f_1^*\omega' = \omega$ with f_1 holomorphic and also that

$$|d\dot{\phi}_0|_{\omega_t}^2 \cdot |F_H|^2 = 0$$

everywhere on X for all values of t . If the connection A_H is not flat, then $\dot{\phi}_0$ is constant in a non empty open set of X , but since $\eta_0 = \eta_{\dot{\phi}_0}$ is holomorphic, then $\dot{\phi}_0$ is constant on X . Therefore f_1 is the identity transformation and $\omega' = \omega$ as claimed. \square

As a corollary of the previous lemma we obtain a result for toric complex surfaces, where smooth geodesics joining any two invariant metrics in \mathcal{K} are known to exist. Recall that a complex manifold is said to be toric if it is endowed with a holomorphic action of a complex torus with a dense orbit. We fix a maximal compact subgroup T of the complex torus and consider \mathcal{K}^T , the subspace of T -invariant Kähler metrics in \mathcal{K} .

Proposition 3.5.2. *Let (E^c, I) be a holomorphic G^c -bundle over a smooth toric compact complex surface (X, J) . If there exists a solution (ω, H) to (3.0.7) with $\omega \in \mathcal{K}^T$, then $\mathcal{M}_H: \mathcal{K}^T \rightarrow \mathbb{R}$ is bounded from below. For any other Kähler metric $\omega' \in \mathcal{K}^T$ such that (ω', H) is a solution to (3.0.7), there exists a holomorphic transformation $f \in \text{Aut}(X, J)$ such that $f^*\omega' = \omega$. Moreover, if the bundle is not topologically flat, then ω is the unique metric in \mathcal{K}^T with this property.*

We obtain a similar result for a complex surface with $c_1(X) \leq 0$. For the proof we adapt the argument of X. X. Chen in [17] concerning the Mabuchi K-energy.

Proposition 3.5.3. *Let (E^c, I) be a holomorphic G^c -bundle over a smooth compact complex surface (X, J) with $c_1(X) \leq 0$.*

- (1) *If there exists a solution $(\omega, H) \in \tilde{\mathcal{K}}$ to (3.0.7), then \mathcal{M}_H is bounded from below.*
- (2) *If the topological constant c vanishes, then ω is the unique metric in \mathcal{K} such that (ω, H) is a solution to (3.0.7).*

PROOF. Let ω' be the unique Kähler–Einstein metric in $[\omega]$, that it is known to exists due to the work of Aubin and Yau when $c_1(X) < 0$ and Yau’s solution to the Calabi Conjecture when $c_1(X) = 0$ (see [8]). Choose an arbitrary Kähler metric $\omega_\phi \in \mathcal{K}$. By [17, Lemma 7] for every $0 < \epsilon \ll 1$ there exists an ϵ -approximate smooth geodesic $\phi_t: [0, 1] \times X \rightarrow \mathbb{R}$ in the space of Kähler potentials joining ω with ω_ϕ , i.e. such that

$$(3.5.4) \quad (\ddot{\phi}_t - |d\dot{\phi}_t|_{\omega_t}^2) \cdot \omega_t^2 = \epsilon \cdot (\omega')^2.$$

Set $\varphi_t = (\ddot{\phi}_t - |d\dot{\phi}_t|_{\omega_t}^2) > 0$, $\omega_t = \omega_{\phi_t}$. Then, the functional \mathcal{M}_H along the curve ω_t satisfies

$$\begin{aligned} \frac{d^2}{dt^2} \mathcal{M}_H(\omega_t) &= 2\alpha_0 \|Y_{\eta_t}\|^2 - \int_X \varphi_t (\alpha_0 S_{\omega_t} - c) \frac{\omega_t^2}{2!} + \alpha_1 \int_X \ddot{\phi}_t |F_H|^2 \cdot \omega^2 \\ &= 2\alpha_0 \|Y_{\eta_t}\|^2 - \alpha_0 \int_X \varphi_t S_{\omega_t} \frac{\omega_t^2}{2!} + \epsilon \cdot c \int_X \frac{(\omega')^2}{2!} + \alpha_1 \int_X \ddot{\phi}_t |F_H|^2 \cdot \omega^2 \\ &\geq -\alpha_0 \int_X \varphi_t S_{\omega_t} \frac{\omega_t^2}{2!} + \epsilon \cdot c \cdot \text{Vol}(X), \end{aligned}$$

where we have used that $\ddot{\phi}_t \geq 0$ due to (3.5.4). Using now the identity (see [17, p. 223])

$$-\int_X \varphi_t S_{\omega_t} \frac{\omega_t^2}{2!} = \int_X |d\varphi_t|_{\omega_t}^2 \frac{1}{\varphi_t} \frac{\omega_t^2}{2!} - \int_X \varphi_t \Lambda_{\omega_t} \rho_{\omega'} \frac{\omega_t^2}{2!}$$

we conclude that

$$(3.5.5) \quad \frac{d^2}{dt^2} \mathcal{M}_H(\omega_t) \geq \epsilon \cdot c \cdot \text{Vol}(X),$$

since $\Lambda_{\omega_t} \rho_{\omega'} \leq 0$ due to the fact that ω' is a Kähler–Einstein metric and $c_1(X) \leq 0$. Therefore, from

$$\frac{d}{dt}|_{t=0} \mathcal{M}_H(\omega_t) = 0,$$

integrating on $[0, 1]$ it follows that

$$\mathcal{M}_H(\omega_\phi) - \mathcal{M}_H(\omega) \geq \epsilon \cdot c \cdot \text{Vol}(X) \cdot \frac{t^2}{2}|_0^1 = \epsilon \cdot \frac{c \cdot \text{Vol}(X)}{2}.$$

Since $c \in \mathbb{R}$ is a topological constant, making $\epsilon \rightarrow 0$, we conclude that \mathcal{M}_H is bounded in \mathcal{K} as claimed in 1). To prove 2) suppose that ω_ϕ is also a solution to (3.0.7) in \mathcal{K} . Note that if $c = 0$ then (3.5.5) together with $\frac{d}{dt}|_{t=1} \mathcal{M}_H(\omega_t) = 0$ gives

$$0 = 2\alpha_0 \|Y_{\eta_t}\|^2 + \alpha_0 \int_X |d\varphi_t|_{\omega_t}^2 \frac{1}{\varphi_t} \frac{\omega_t^2}{2!} + \alpha_1 \int_X \ddot{\phi}_t |F_H|^2 \cdot \omega^2,$$

from which η_t is a holomorphic vector field on (X, J) for all $t \in [0, 1]$. Since η_t is a complex Hamiltonian vector field, it vanishes somewhere on X but since $c_1(X) \leq 0$ then η_t is parallel and so it identically vanishes. Therefore $\omega = \omega_\phi$ as claimed. \square

Remark 3.5.4. Note that as $\mathfrak{z} \ni z = 0$ in (3.0.7), $c = 0$ implies that $c_1(X) = 0$ and $\hat{c} = 0$ (defined in (2.2.6)), and so the bundle is topologically trivial.

We finish this chapter proving that the lower bound on \mathcal{M}_H is an obstruction to the existence of solutions to (3.0.7) in any compact complex surfaces which admits a Kähler metric with nonnegative bisectional curvature. We use the result about weak solutions of the geodesic equation in \mathcal{K} found by X. X. Chen and G. Tian in [19] and adapt their argument about the Mabuchi K-energy. The condition on the bisectional curvature is imposed in order to have weak solutions to the geodesic equation in the space of Kähler potentials with bounded second derivatives (see [13]).

Proposition 3.5.5. *Let (E^c, I) be a holomorphic G^c -bundle over a smooth compact complex surface (X, J) which admits a Kähler metric ω_0 with nonnegative bisectional curvature. If there exists a solution $(\omega, H) \in \tilde{\mathcal{K}}$ to (3.0.7) with $[\omega] = [\omega_0]$, then \mathcal{M}_H is bounded from below.*

PROOF. Let $(\omega, H) \in \tilde{\mathcal{K}}$ be a solution to (3.0.7). Let $\omega_\phi = \omega + dd^c\phi$ be any other Kähler metric in \mathcal{K} for a given Kähler potential ϕ . We will prove that

$$(3.5.6) \quad \mathcal{M}_H(\omega_\phi) - \mathcal{M}_H(\omega) \geq 0.$$

To do this, we will use the following decomposition formula (cf. [19, formula 6.1])

$$(3.5.7) \quad \begin{aligned} \mathcal{M}_H(\omega_\phi) &= \alpha_0 \cdot \mathcal{M}_\omega(\omega_\phi) + c' \cdot I(\omega_\phi) + \alpha_1 \int_X \phi |F_H|^2 \cdot \omega^2, \quad \text{with} \\ I(\omega_\phi) &= \frac{1}{3!} \sum_{j=0}^2 \int_X \phi \cdot \omega_\phi^j \wedge \omega^{2-j}, \end{aligned}$$

where $c' = c - \alpha_0 \hat{S} \in \mathbb{R}$ and \mathcal{M}_ω is the Mabuchi K-energy on \mathcal{K} . Given any positive integer l , let $\Sigma_l = [-l, l] \times [0, 1]$ be the rectangle domain, regarded as Riemann surface with boundary, and consider the homogeneous complex Monge–Ampère equation

$$(3.5.8) \quad (\pi^* \omega + 2 \mathbf{i} \partial \bar{\partial} \phi)^{2+1} = 0,$$

where $\phi : \Sigma_l \times X \rightarrow \mathbb{R}$ is a real function and $\pi : \Sigma_l \times X \rightarrow X$ denotes the projection. Recall from [23] that a curve of Kähler potentials ϕ_t , $t \in [0, 1]$, satisfies the geodesic equation (1.3.12) if and only the function $\phi(s, t, x) = \phi_t(x)$ on $\Sigma_l \times X$ satisfies (3.5.8), and also that an arbitrary solution $\phi(s, t)$ of (3.5.8) on the square Σ_l satisfies

$$(3.5.9) \quad \Delta_{s,t} \phi^l := \left(\frac{d^2}{ds^2} + \frac{d^2}{dt^2} \right) \phi^l = |d(\frac{d}{dt} \phi^l) - d^c(\frac{d}{ds} \phi^l)|_{\omega_{\phi^l}}^2.$$

By [19, Theorem 1.1], there exists an almost smooth solution ϕ^l to (3.5.8) (see [19, Definition 1.3.3]) with boundary values

$$\phi^l(s, 0) = 0, \quad \phi^l(s, 1) = \phi, \quad \phi^l(\pm l, t) = (1 - t)\phi \quad (s, t) \in \Sigma_l.$$

To be more precise, one may modify slightly the boundary of Σ_l and the boundary data so that the domain is smooth without corners (see [19, Section 6.2]). Let us fix

a Kähler metric on $\Sigma_l \times X$ that does not depend on l . According to [19], there exists a positive constant $C > 0$ independent of l such that

$$(3.5.10) \quad 0 \leq |\partial\bar{\partial}\phi^l| \leq C,$$

where the norm is taken with respect to the fixed Kähler metric on $\Sigma_l \times X$ and

$$0 \leq (\omega_{\phi^l(s,t)})^2$$

for any $(s, t) \in \Sigma_l$. Note here that $\omega_{\phi^l(s,t)}$ can be less regular than C^∞ (a priori it is only a bounded two form on X). Moreover, the additional assumption on the existence of a Kähler metric on \mathcal{K} with nonnegative bisectional curvature, implies that all the partial derivatives of second order of ϕ^l are bounded. Due to the decomposition formula for the K-energy (see [19, formula 6.1]), it induces a well defined map

$$\mathcal{M}_\omega: \Sigma_l \rightarrow \mathbb{R}: (s, t) \mapsto \mathcal{M}_\omega(\omega_{\phi^l(s,t)}).$$

This map is C^1 continuous up to the boundary of Σ_l and weakly subharmonic in the interior (see [19, Theorem 6.1.1]). The previous bounds on ϕ^l and the decomposition formula (3.5.7) implies then that we have a well defined map

$$\mathcal{M}_H: \Sigma_l \rightarrow \mathbb{R}: (s, t) \mapsto \mathcal{M}_H(\omega_{\phi^l(s,t)}).$$

Let $\nu : \mathbb{R} \rightarrow \mathbb{R}$ be a smooth non-negative function such that $\nu \equiv 1$ on $[-\frac{1}{2}, \frac{1}{2}]$ and vanishes outside $[-\frac{3}{4}, \frac{3}{4}]$. Set

$$\nu^l(s) = \frac{\nu(\frac{s}{l})}{\int_{\mathbb{R}} \nu(s) ds}$$

and define

$$\begin{aligned} f^l(t) &= \int_{\mathbb{R}} \nu^l(s) \mathcal{M}_H(s, t) ds \\ &= \alpha_0 \cdot \int_{\mathbb{R}} \nu^l(s) \mathcal{M}_\omega(s, t) ds + c' \cdot \int_{\mathbb{R}} \nu^l(s) I(s, t) ds \\ &\quad + \alpha_1 \cdot \int_{\mathbb{R}} \nu^l(s) \int_X \phi^l(s, t) |F_H|^2 \omega^2 \wedge ds, \end{aligned}$$

where $I(s, t) = I(\omega_{\phi^l(s,t)})$. Then,

$$f^l(0) = \int_{\mathbb{R}} \nu^l(s) \mathcal{M}_H(\omega) ds = \mathcal{M}_H(\omega)$$

$$f^l(1) = \int_{\mathbb{R}} \nu^l(s) \mathcal{M}_H(\omega_\phi) ds = \mathcal{M}_H(\omega_\phi)$$

and

$$\frac{df^l}{dt} \Big|_{t=0} = 0,$$

since (ω, H) is a solution to (3.0.7). Then, we have that

$$\begin{aligned}
f^l(1) - f^l(0) &= \int_0^1 \int_0^r \frac{d^2 f^l}{dt^2} dt \wedge dr \\
&= \int_0^1 \int_0^r \int_{\mathbb{R}} \nu^l(s) \frac{d^2 \mathcal{M}_H(s, t)}{dt^2} ds \wedge dt \wedge dr \\
&= \int_0^1 \int_0^r \int_{\mathbb{R}} \nu^l(s) \Delta_{s,t} \mathcal{M}_H(s, t) ds \wedge dt \wedge dr \\
&\quad - \int_0^1 \int_0^r \int_{\mathbb{R}} \nu^l(s) \frac{d^2 \mathcal{M}_H(s, t)}{ds^2} ds \wedge dt \wedge dr \\
&= \alpha_0 \int_0^1 \int_0^r \int_{\mathbb{R}} \nu^l(s) \Delta_{s,t} \mathcal{M}_{\omega}(s, t) ds \wedge dt \wedge dr \\
&\quad + \alpha_1 \cdot \int_0^1 \int_0^r \int_{\mathbb{R}} \nu^l(s) \int_X \Delta_{s,t} \phi^l |F_H|^2 \cdot \omega^2 \wedge ds \wedge dt \wedge dr \\
&\quad - \int_0^1 \int_0^r \int_{\mathbb{R}} \nu^l(s) \frac{d^2 \mathcal{M}_H(s, t)}{ds^2} ds \wedge dt \wedge dr,
\end{aligned}$$

where we have used that ϕ^l has bounded second derivatives and that the I term makes no contribution to the second derivatives along solutions to (3.5.8) (see the comments before the proof of [19, Theorem 6.1.1]). Note also that

$$\Delta_{s,t} \phi^l = \frac{\partial^2 \phi^l}{\partial z \partial \bar{z}},$$

which is bounded by (3.5.10) and is non negative due to (3.5.9), where $z = t + \mathbf{i}s$ is the complex coordinate on Σ_l . The previous facts imply that we have a lower bound

$$\begin{aligned}
f^l(1) - f^l(0) &\geq - \int_0^1 \int_0^r \int_{\mathbb{R}} \nu^l(s) \frac{d^2 \mathcal{M}_H(s, t)}{ds^2} ds \wedge dt \wedge dr \\
&= - \int_0^1 \int_0^r \int_{\mathbb{R}} \frac{d^2 \nu^l(s)}{ds^2} \mathcal{M}_H(s, t) ds \wedge dt \wedge dr \\
&= - \frac{1}{l^2 v} \int_0^r \int_{\mathbb{R}} \frac{d^2 \nu^l}{ds^2} \Big|_{\frac{s}{l}} \mathcal{M}_H(s, t) ds \wedge dt \wedge dr,
\end{aligned}$$

where $v = \int_{\mathbb{R}} \nu(s) ds$. The bound (3.5.10) and the decomposition formulae (3.5.7) for \mathcal{M}_H and \mathcal{M}_{ω} provide a uniform bound $|\mathcal{M}_H(s, t)| \leq C'$ that, as in the proof of [19, Theorem 6.2.1], gives

$$\mathcal{M}_H(\omega_{\phi}) - \mathcal{M}_H(\omega) = f^l(1) - f^l(0) \geq - \frac{C'}{2l}.$$

This shows that (3.5.6) holds as $l \rightarrow \infty$. \square

CHAPTER 4

Algebraic stability

Let (X, L, E) be a triple consisting of a compact complex manifold X , an ample line bundle L and a holomorphic vector bundle E over X . We know from Section §3 that the pairs (ω, H) consisting of a Kähler metric $\omega \in c_1(L)$ and a Hermitian metric H on E correspond to points in the “complexified orbit” of the holomorphic structure of the bundle of frames of E with respect to the action of the extended gauge group. This fact, together with the moment map interpretation of the coupled equations (3.0.7) leads to formulate a conjecture in the vain of the usual identification of algebraic and symplectic quotients of projective manifolds provided by the Kempf-Ness Theorem, namely: *the existence of solutions (ω, H) to the coupled equations (3.0.7), with $\omega \in c_1(L)$, is equivalent to a suitable polystability of the triple (X, L, E) in the sense of GIT.* This chapter is devoted to the formulation of a stability notion adapted to this problem and to give evidence that the above conjecture holds.

4.1. K-stability

In this section we propose a notion of stability for a triple (X, L, E) consisting of a projective scheme X , an ample invertible sheaf L and a coherent sheaf E over X . Our stability notion depends on a positive real parameter α and is modelled on Donaldson’s definition of K -stability for polarized varieties (see §1.3.6). As in [24], it is formulated in terms of flat degenerations of the triple over \mathbb{C} . The central fiber of each degeneration is endowed with a \mathbb{C}^* -action and the stability of the triple, that we call α -K-stability, is controlled by a suitable numerical invariant F_α (see (4.1.2)) associated to this action. We are mainly interested in the case in which E is a holomorphic vector bundle over a smooth projective variety X , but even in this case we need to consider singular degenerations and hence our geometric objects in this section will be arbitrary algebraic schemes and coherent sheaves, rather than vector bundles and manifolds. When the degeneration considered is smooth, F_α will coincide, up to a multiplicative constant factor, with the value of the character \mathcal{F}_α given by (3.3.3) at the generator of the $S^1 \subset \mathbb{C}^*$ action (see (4.2)).

Let (X, L, E) be a triple consisting of a projective scheme X , an ample invertible sheaf L and a coherent sheaf E over X . We start defining the degenerations of (X, L, E) that we will consider. In what follows, a \mathbb{C}^* -*action* on such a triple means a \mathbb{C}^* -action on X with \mathbb{C}^* -linearisations on L and E .

Definition 4.1.1. A *test configuration* for (X, L, E) is a triple $(\mathcal{X}, \mathcal{L}, \mathcal{E})$ consisting of a scheme \mathcal{X} , an invertible sheaf \mathcal{L} and a coherent sheaf \mathcal{E} over \mathcal{X} , together with a \mathbb{C}^* -action on $(\mathcal{X}, \mathcal{L}, \mathcal{E})$ and a flat morphism $\pi: \mathcal{X} \rightarrow \mathbb{C}$, such that

- (1) the invertible sheaf \mathcal{L} is ample and the sheaf \mathcal{E} is flat over \mathbb{C} ,

- (2) $\pi : \mathcal{X} \rightarrow \mathbb{C}$ is \mathbb{C}^* -equivariant, where \mathbb{C}^* acts on \mathbb{C} by multiplication in the standard way,
- (3) the fibre (X_t, L_t, E_t) is isomorphic to (X, L, E) for all $t \in \mathbb{C} \setminus \{0\}$, where $X_t = \pi^{-1}(t)$, $L_t = \mathcal{L}|_{X_t}$ and $E_t = \mathcal{E}|_{X_t}$.

Any \mathbb{C}^* -action on (X, L, E) determines a special type of test configuration, called a *product configuration*, where $\mathcal{X} \cong X \times \mathbb{C}$ with the induced \mathbb{C}^* -action, \mathcal{L} and \mathcal{E} are obtained by pulling the \mathbb{C}^* -linearised sheaves L and E , and π is the canonical projection. This product configuration is called a *trivial configuration* if the given \mathbb{C}^* -action on (X, L, E) is trivial, i.e. if it consists of the trivial \mathbb{C}^* -action on X and the \mathbb{C}^* -action defined by scalar multiplication on L and E .

Let (X, L, E) be as in Definition 4.1.1 and such that E has n -dimensional support, where $n = \dim X$. Let $(\mathcal{X}, \mathcal{L}, \mathcal{E})$ be a test configuration for (X, L, E) . Then, the fibres (X_t, L_t, E_t) are all necessarily isomorphic for $t \neq 0$, due to the \mathbb{C}^* -action. Since π and \mathcal{E} are flat the Hilbert polynomial of E_t is independent of $t \in \mathbb{C}$ [34, Ch. III, Thm 9.9], so the central fibre (X_0, L_0, E_0) satisfies that X_0 is an n -dimensional projective scheme, L_0 is ample and E_0 has n -dimensional support. Recall that the degree of the Hilbert polynomial of E_t equals the dimension of its support (see [34, Ch. III, Ex. 5.7(a)] for the fact that L_t is ample). As $0 \in \mathbb{C}$ is fixed, the central fibre (X_0, L_0, E_0) has an induced \mathbb{C}^* -action. For each integer k , we define the integer $w(k)$ as the weight of the induced \mathbb{C}^* -action on the determinant of the cohomology of $E_0 \otimes L_0^k$, i.e. on the line (cf. e.g. [37, §§1.2 and 2.1])

$$\det H^*(X_0, E_0 \otimes L_0^k) := \bigotimes_{p=0}^n \left(\det H^p(X_0, E_0 \otimes L_0^k) \right)^{(-1)^p}.$$

Then, by the equivariant Riemann–Roch theorem, $w(k)$ is a polynomial of degree at most $n+1$ in k (cf. [24, §2.1]). Thus the following quotient has an expansion

$$(4.1.1) \quad F(E_0, k) := \frac{w(k)}{k P_{L_0}(E_0, k)} = F_0(E_0) + k^{-1} F_1(E_0) + k^{-2} F_2(E_0) + O(k^{-3})$$

with rational coefficients $F_i(E_0)$, where $P_{L_0}(E_0, k)$ is the Hilbert polynomial of E_0 , which has degree n in k . Now, given a real number $\alpha > 0$, the α -*invariant* of the test configuration $(\mathcal{X}, \mathcal{L}, \mathcal{E})$ is

$$(4.1.2) \quad F_\alpha(\mathcal{X}, \mathcal{L}, \mathcal{E}) := -F_1(\mathcal{O}_{X_0}) - \alpha (F_2(E_0) - F_2(\mathcal{O}_{X_0})),$$

where $F_j(\mathcal{O}_{X_0})$, $0 \leq j \in \mathbb{Z}$, are obtained from formula (4.1.1) setting $E_0 = \mathcal{O}_{X_0}$. When there is no confusion we will denote it just by F_α .

We already have all the ingredients to define the notion of α -K-stability. As usual, we define the *degree* of a coherent sheaf over a projective scheme with respect to a given ample invertible sheaf in terms of the first two leading coefficients of its Hilbert polynomial (see [37, Definition 1.2.11] for the precise definition).

Definition 4.1.2. Let (X, L, E) be a triple consisting of a projective scheme of dimension n , an ample invertible sheaf and a coherent sheaf of degree zero which is pure of dimension n . Let α be a positive real number. Then the triple (X, L, E) is said to

be α -K-semistable if

$$(4.1.3) \quad F_{k\alpha}(\mathcal{X}, \mathcal{L}, \mathcal{E}) \geq 0$$

for all integers $k > 0$ and all test configurations $(\mathcal{X}, \mathcal{L}, \mathcal{E})$ for (X, L^k, E) . This triple is α -K-stable if the inequality (4.1.3) is strict for all non-trivial $(\mathcal{X}, \mathcal{L}, \mathcal{E})$. It is called α -K-polystable if it is α -K-semistable and the inequality (4.1.3) is always strict unless $(\mathcal{X}, \mathcal{L}, \mathcal{E})$ is a product configuration.

The assumption in Definition 4.1.2 that E has degree zero is used to make the formulation of α -K-stability simpler. When the sheaf E has non-zero degree, the notion of α -K-stability involves a formula for the definition of the invariant F_α which is more complicated than (4.1.2). This is an interesting issue to which we wish to return in future work.

As usual, any triple as in Definition 4.1.2 satisfies the following chain of implications: (X, L, E) is α -K-stable \Rightarrow (X, L, E) is α -K-polystable \Rightarrow (X, L, E) is α -K-semistable.

An α -K-unstable triple is a triple which is not α -K-semistable, and an α -destabilizing test configuration for (X, L, E) is a test configuration with $F_\alpha < 0$. Note that α -K-unstability is an open condition in α , i.e. if a triple is not α -K-semistable then it admits a α -destabilizing test configuration $(\mathcal{X}, \mathcal{L}, \mathcal{E})$ and, by definition, it also destabilizes the triple for close values of the constant α .

Remark 4.1.3. Note that if (X, L, E) , as in Definition 4.1.2, is α -K-(semi/poly)stable then for any integer $k > 0$ the triple (X, L^k, E) is $(k\alpha)$ -K-(semi/poly)stable. Note also that

$$F_{k\alpha}(\mathcal{X}, \mathcal{L}^k, \mathcal{E}) = F_\alpha(\mathcal{X}, \mathcal{L}, \mathcal{E}),$$

for any test configuration $(\mathcal{X}, \mathcal{L}, \mathcal{E})$. This follows simply because the invariants F_i defined in (4.1.1) satisfy $F_i(L_0^k, E_0) = k^{1-i} F_i(L_0, E_0)$, as $F(L_0^k, E_0, h) = kF(L_0, E_0, hk)$.

4.2. A differential-geometric formula for the α -invariant

In this section we relate the character \mathcal{F}_α defined in (3.3.3) and the invariant F_α (4.1.2) of a test configuration when it has smooth central fiber. As a consequence we prove that the existence of solutions to the coupled equations (3.0.7) implies the vanishing of the α -invariant on product test configurations.

Let (X, L, E) be a triple consisting of a smooth compact polarized variety (X, L) and a holomorphic vector bundle E over X with rank r and degree zero. Given a positive integer $k > 0$, let $(\mathcal{X}, \mathcal{L}, \mathcal{E})$ be a test configuration for the triple (X, L^k, E) with smooth central fiber (X_0, L_0, E_0) . Associated to E_0 we have the holomorphic $\mathrm{GL}(r, \mathbb{C})$ -principal bundle of frames (E_0^c, I) and the character

$$\mathcal{F}_\alpha : \mathrm{Lie} \mathrm{Aut}(E_0^c, I) \rightarrow \mathbb{C}$$

defined in (3.3.3), corresponding to the holomorphic structure I , the Kähler class $c_1(L_0)$, the biinvariant pairing $(\cdot, \cdot) = -\mathrm{tr}$ in the Lie algebra of $\mathrm{GL}(r, \mathbb{C})$ and a pair of real constants $\alpha_0, k \cdot \alpha_1 \in \mathbb{R}$. Let $\zeta \in \mathrm{Lie} \mathrm{Aut}(E_0^c, I)$ be the holomorphic vector field generating the induced S^1 -action on E_0 and define the real constant $\alpha = \frac{\alpha_1 \pi^2 r}{\alpha_0}$.

Proposition 4.2.1.

$$(4.2.1) \quad F_{k\alpha}(\mathcal{X}, \mathcal{L}, \mathcal{E}) = \frac{-1}{4 \operatorname{vol}(X)\alpha_0} \mathcal{F}_\alpha(\zeta).$$

PROOF. To simplify the notation, we will drop the zero subscript in the smooth triple corresponding to the central fiber of the test configuration. We have to compute the coefficients of the normalized weight

$$F(E, m) = \frac{w_m}{m P_L(E, m)} = F_0(E) + m^{-1} F_1(E) + m^{-2} F_2(E) + O(m^{-3})$$

of the \mathbb{C}^* -action on the sheaf cohomology of $E(m)$ though since the action is algebraic it is enough to compute the weight of the induced S^1 -action. We will use the Riemann-Roch formula and its equivariant version for what we choose metrics h on L and H on E such that $\omega = \frac{i}{2\pi} F_h$ is a Kähler metric on X and $\operatorname{tr} F_H = 0$ (we are allowed to suppose this since the degree of E is zero). Averaging over S^1 if necessary we can suppose in addition that h and H are fixed by the S^1 -action. Recall that we have expansions

$$\begin{aligned} \frac{P_L(E, m)}{m^n} &= d_0 + m^{-1} d_1 + m^{-2} d_2 + O(m^{-3}) \\ \frac{w_m}{m^{n+1}} &= c_0 + m^{-1} c_1 + m^{-2} c_2 + O(m^{-3}). \end{aligned}$$

A direct computation shows then that

$$(4.2.2) \quad \begin{aligned} F_0(L, E) &= \frac{c_0}{d_0} \\ F_1(L, E) &= \frac{c_1 d_0 - c_0 d_1}{d_0^2} \\ F_2(L, E) &= \frac{-d_1 F_1(E)}{d_0} + \frac{c_2 d_0 - c_0 d_2}{d_0^2}. \end{aligned}$$

By Riemann-Roch formula we have

$$P_L(E, m) = \int_X ch(E \otimes L^m) Td(X) = \int_X \operatorname{tr} e^{m\omega \operatorname{Id} + \frac{i}{2\pi} F_H} Td(X).$$

Taking representatives $\rho_i \in \Omega^i(X)$ of the Todd class of X , i.e.

$$Td(X) = [1 + \rho_1 + \rho_2 + \dots] \in H^*(X, \mathbb{Z}),$$

we obtain coefficients

$$\begin{aligned} d_0 &= r \cdot \operatorname{vol}(X) \\ d_1 &= r \cdot \int_X \operatorname{tr} \Lambda \rho_1 \frac{\omega^n}{n!} \\ d_2 &= \frac{1}{2} \int_X \Lambda_\omega^2 \left[\frac{-1}{8\pi^2} \operatorname{tr} F_H \wedge F_H + r \cdot \rho_2 \right] \cdot \frac{\omega^n}{n!}. \end{aligned}$$

To compute the weight we use, as in [22], the equivariant Riemann-Roch formula. By this formula, w_m is given by the coefficient of t in

$$\int_X ch(E \otimes L^m)(t) \cdot Td(X)(t),$$

where now the Chern and Todd classes are regarded as equivariant cohomology classes. Let ζ be holomorphic vector field generating the S^1 -action on E and η the holomorphic vector field on X that it covers. To make the calculation we use the de Rham model of equivariant cohomology with the complex $(\Omega_X^*)^{S^1}[t]$ and differential $d + ti_\eta$. Now, writing

$$\zeta = \theta_{A_H}^\perp(\eta) + a$$

where $a \in \text{Lie } \mathcal{G}_H$, we have relations (see (2.2.9) and cf. [22, Section 2.2])

$$d_{A_H} a + i_y F_H = 0; \quad df = i_\eta \omega$$

for some $f \in C^\infty(X, \mathbb{R})$, provided that ζ and η are real holomorphic vector fields, η lifts to a holomorphic vector field on L and A_H is the Chern connection of H . Then the equivariant Chern character $ch(E \otimes L^m)(t) \in H_{S^1}^2(X)$ is represented by

$$ch(t) = \text{tr } e^{m[\omega - tf] + \frac{i}{2\pi}[F_H + ta]}$$

in this model. Let $Td(t) := 1 + (\rho_1 + tR_1) + (\rho_2 + tR_2) + \dots$ be the representative for $Td(X)$, where R_1 is a function on X and R_2 is a 2-form. Then

$$\begin{aligned} \int_X ch(t) \cdot Td(X)(t) &= \int_X \text{tr } e^{m(\omega \text{Id} - tf) + \frac{i}{2\pi}(F_H + ta)} \cdot Td(t) \\ &= \int_X \text{tr} \sum_{l=0}^{n+1} \frac{1}{l!} (m(\omega \text{Id} - tf) + \frac{i}{2\pi}(F_H + ta))^l \cdot \\ &\quad \cdot [1 + (\rho_1 + tR_1) + (\rho_2 + tR_2)] + O(t^2) \\ &= r \cdot \int_X (m^n \frac{\omega^n}{n!} + m^{n-1} \frac{\omega^{n-1}}{(n-1)!}) \wedge (1 + \rho_1 + \rho_2) \\ &\quad + t \cdot r \cdot \int_X (m^n \frac{\omega^n}{n!} + m^{n-1} \frac{\omega^{n-1}}{(n-1)!}) \wedge (R_1 + R_2) \\ &\quad + t \cdot \int_X \text{tr} [-fm^{n+1} \frac{\omega^n}{n!} \text{Id} + m^n (\frac{\mathbf{i}}{2\pi} a \frac{\omega^n}{n!} - f \frac{\omega^{n-1}}{(n-1)!} \text{Id}) \\ &\quad + m^{n-1} (\frac{-a\Lambda F_H}{4\pi^2} + \frac{\mathbf{i} a \omega^{n-1}}{2\pi(n-1)!} + \frac{f\Lambda^2 F_H^2 \omega^n}{16\pi^2 n!} \\ &\quad - \frac{f\omega^{n-2}}{(n-2)!} \text{Id})] \wedge (1 + \rho_1 + \rho_2) + m^{n+1} O(t^2, m^{-3}). \end{aligned}$$

Hence

$$\begin{aligned}
w_m &= \frac{d}{dt}_{|t=0} \int_X ch(L^m \otimes E)(t) \cdot Td(X)(t) \\
&= m^{n+1} \left(-r \cdot \int_X f \frac{\omega^n}{n!} \right) + m^n \left(\int_X (\text{tr} \frac{\mathbf{i}}{2\pi} a + r \cdot (R_1 - f \Lambda \rho_1)) \frac{\omega^n}{n!} \right) + \\
&\quad + m^{n-1} \left(\int_X \text{tr} \frac{\mathbf{i} a}{2\pi} \cdot \Lambda \left(\frac{\mathbf{i} F_H}{2\pi} + \rho_1 \text{Id} \right) \frac{\omega^n}{n!} + r \cdot \int_X \Lambda R_2 \frac{\omega^n}{n!} + \right. \\
&\quad \left. - \frac{1}{2} \int_X f \text{tr} \Lambda^2 \left(\frac{-1}{8\pi^2} F_H^2 + \rho_2 \text{Id} \right) \frac{\omega^n}{n!} \right) + m^{n+1} O(m^{-3})
\end{aligned}$$

and the coefficients of the weight are

$$\begin{aligned}
(4.2.3) \quad c_0 &= -r \cdot \int_X f \frac{\omega^n}{n!}, \\
c_1 &= \int_X (\text{tr} \frac{\mathbf{i}}{2\pi} a + r \cdot (R_1 - f \Lambda \rho_1)) \frac{\omega^n}{n!}, \\
c_2 &= \int_X \text{tr} \frac{\mathbf{i} a}{2\pi} \cdot \Lambda \left(\frac{\mathbf{i} F_H}{2\pi} + \rho_1 \text{Id} \right) \frac{\omega^n}{n!} + r \cdot \int_X \Lambda R_2 \frac{\omega^n}{n!}, \\
&\quad - \frac{1}{2} \int_X f \text{tr} \Lambda^2 \left(\frac{-1}{8\pi^2} F_H^2 + \rho_2 \text{Id} \right) \frac{\omega^n}{n!}.
\end{aligned}$$

Note that taking traces in $d_{A_H} a + i_y F_H = 0$ we obtain that $\text{tr } a$ is constant. Now, if for any function $f \in C^\infty(X)$ we define its average as

$$\hat{f} = \frac{\int_X f \frac{\omega^n}{n!}}{\text{vol}(X)}$$

it follows from (4.2.2) that

$$\begin{aligned}
(4.2.4) \quad F_0(E) &= -\hat{f} \\
F_1(E) &= \frac{\mathbf{i} \text{tr } a}{2\pi r} + \widehat{R}_1 - \frac{1}{\text{vol}(X)} \int_X f (\Lambda \rho_1 - \widehat{\Lambda \rho_1}) \frac{\omega^n}{n!}, \\
F_2(E) &= \frac{-1}{4\pi^2 \text{vol}(X)r} \int_X \text{tr} \left(a \Lambda F_H - \frac{1}{4} f (\Lambda^2 F_H^2 - \widehat{\Lambda^2 F_H^2}) \right) \frac{\omega^n}{n!}, \\
&\quad - \widehat{\Lambda \rho_1} \cdot F_1(\mathcal{O}_X) + \widehat{\Lambda R_2} - \frac{1}{2 \text{Vol}(X)} \int_X f (\Lambda^2 \rho_2 - \widehat{\Lambda^2 \rho_2}) \frac{\omega^n}{n!},
\end{aligned}$$

where $F_1(\mathcal{O}_X)$ has been obtained from the second formula setting $E = \mathcal{O}_X$. A similar manipulation with the F_2 term shows that

$$F_2(E) - F_2(\mathcal{O}_X) = \frac{- \int_X \text{tr} \left(a \Lambda F_H - \frac{1}{4} f (\Lambda^2 F_H^2 - \widehat{\Lambda^2 F_H^2}) \right) \frac{\omega^n}{n!}}{4\pi^2 \text{vol}(X)r}.$$

Note that $\rho_\omega/2$ also represents the cohomology class of ρ_1 . Since the previous computation does not depend on the chosen representative we can set $\rho_\omega/2 = \rho_1$. Then, we now from Lemma 2.2.3 in [23] that $\widehat{R}_1 = 0$ so

$$F_1(\mathcal{O}_X) = \frac{-1}{4 \text{vol}(X)} \int_X f (S_\omega - \widehat{S}) \frac{\omega^n}{n!}$$

where we have used the identity $\Lambda_\omega \rho_\omega = S_\omega/2$. Taking coupling constants $\alpha_0, k\alpha_1 \in \mathbb{R}$ for the definition of \mathcal{F}_α and setting $\alpha = \frac{r\pi^2\alpha_1}{\alpha_0}$ we get

$$F_{k\alpha}(\mathcal{X}, \mathcal{L}, \mathcal{E}) = \frac{-1}{4 \operatorname{vol}(X)\alpha_0} \mathcal{F}_\alpha(\zeta). \quad \square$$

As a corollary of the previous result we obtain the following relation between the existence of solutions to the coupled equations (3.0.7) and product test configurations for the triple.

Corollary 4.2.2. *Let (X, L, E) be a triple consisting of a smooth compact polarized variety (X, L) and a holomorphic vector bundle E over X with rank r and degree zero. Suppose that there exists a solution (ω, H) to the coupled equations (3.0.7) with $\omega \in c_1(L)$ and positive coupling constants $\alpha_0, \alpha_1 \in \mathbb{R}$. Let $\alpha = \frac{r\pi^2\alpha_1}{\alpha_0}$. Then, given a positive integer $k > 0$, and a product test configuration $(\mathcal{X}, \mathcal{L}, \mathcal{E})$ for the triple (X, L^k, E) , we have*

$$F_{k\alpha}(\mathcal{X}, \mathcal{L}, \mathcal{E}) = 0.$$

PROOF. Note that given $k > 0$, the initial solution (ω, H) to (3.0.7) provides another solution $(k \cdot \omega, H)$ on (X, L^k, E) with coupling constants $(\alpha_0, k\alpha_1)$. Note also that since $(\mathcal{X}, \mathcal{L}, \mathcal{E})$ is a product test configuration, the central fiber is isomorphic to (X, L^k, E) . The result follows now from Proposition 4.2.1 and Proposition 3.3.3. \square

Proposition 4.2.1 shows that $F_{k\alpha}$ is, up to multiplicative constant factor, equal to $\mathcal{F}_\alpha(\zeta)$ and hence to the moment map (2.3.3) whose zero locus corresponds to the solutions of the coupled equations (3.0.7). Once more this fits in the general picture of the Kempf-Ness Theorem explained in §1.1 if we think of $F_{k\alpha}$ as the Mumford weight (see §1.1.3 and cf. (1.1.9)) of a suitable 1-parameter subgroup associated to each test configuration. This interpretation of a test configuration is the topic of the next section.

4.3. Test configurations as 1-parameter subgroups

Test configurations for triples (X, L, E) can be interpreted as 1-parameter subgroups of a certain group acting on a relative Quot-scheme over a Hilbert scheme. This generalizes a similar description of test configurations for polarized varieties (X, L) in terms of 1-parameter subgroups of a certain group acting on a Hilbert scheme, first given by Donaldson [24] and studied further by Ross & Thomas [64]. When $\dim_{\mathbb{C}} X = 1$ and L is a power of the canonical bundle, Pandharipande [60] used Geometric Invariant Theory, for the same group action that we consider on a relative Quot-scheme to compactify the moduli space of pairs consisting of a smooth curve of genus greater than one and a semistable vector bundle over the curve.

Let H and P be polynomials of the same degree n with positive leading coefficients. Let V_k and W_k be complex vector spaces of dimensions $H(k)$ and $P(k)$, respectively, for each k large enough such that $H(k)$ and $P(k)$ are positive. Let

$$\operatorname{Hilb} = \operatorname{Hilb}_{\mathbb{P}}^{H'}$$

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be the Hilbert scheme parametrising closed subschemes of the projective space $\mathbb{P} = \mathbb{P}V_k$ with Hilbert polynomial $H'(h) := H(kh)$. Let $\mathcal{X} \subset \mathbb{P} \times \text{Hilb}$ be the universal family of subschemes of \mathbb{P} parametrised by Hilb and $\pi_H: \mathcal{X} \rightarrow \text{Hilb}$ the composition of the closed embedding $\mathcal{X} \hookrightarrow \mathbb{P} \times \text{Hilb}$ and the canonical projection $\mathbb{P} \times \text{Hilb} \rightarrow \text{Hilb}$. Then \mathcal{X} is a relative projective scheme over Hilb via the map π_H , and $\mathcal{L} = p^*\mathcal{O}_{\mathbb{P}}(1)$ is relatively very ample, where $p: \mathcal{X} \rightarrow \mathbb{P}$ is the composition of $\mathcal{X} \hookrightarrow \mathbb{P} \times \text{Hilb}$ and the projection $\mathbb{P} \times \text{Hilb} \rightarrow \mathbb{P}$. Furthermore, the fibre $X_x = \pi_H^{-1}(x)$ over a point $x \in \text{Hilb}$ is the closed subscheme of \mathbb{P} parametrised by x , and the restriction $L_x = \mathcal{L}|_{X_x}$ is very ample. Let

$$\text{Quot} = \text{Quot}_{\mathcal{X}/\text{Hilb}}^{P'}(W_k \otimes \mathcal{L}^{-1})$$

be the relative Quot-scheme which parametrises quotient sheaves of the fibres of $W_k \otimes \mathcal{L}^{-1}$ over the fibres of π_H with Hilbert polynomial $P'(h) := P(kh)$ with respect to \mathcal{L} . Let

$$\varphi: W_k \otimes q^*\mathcal{L} \twoheadrightarrow \mathcal{E}$$

be the universal family of quotient sheaves parametrised by Quot , where q is the projection $\mathcal{X} \times \text{Quot} \rightarrow \mathcal{X}$. Note that there is a natural projective map $\kappa: \text{Quot} \rightarrow \text{Hilb}$. Then \mathcal{E} is a coherent sheaf over $\mathcal{X} \times \text{Quot}$, flat over Quot , and the fibre of φ over a point $e \in \text{Quot}$ is the quotient sheaf $\varphi_e: W_k \otimes L_x \twoheadrightarrow E_e$ parametrised by e , where $x = \kappa(e)$.

The group $\text{GL}(V_k)$ acts canonically on \mathbb{P} , inducing an action on Hilb and hence on $\mathbb{P} \times \text{Hilb}$. As the latter preserves \mathcal{X} , $\text{GL}(V_k)$ also acts on Quot and π_H is clearly $\text{GL}(V_k)$ -equivariant. As Quot also has a canonical $\text{GL}(W_k)$ -action and both actions commute, the product $G_k = \text{GL}(V_k) \times \text{GL}(W_k)$ also acts on Quot .

Let (X, L, E) be a triple consisting of a projective scheme, an ample invertible sheaf and a coherent sheaf, such that the Hilbert polynomials of (X, L) and E with respect to L are H and P , respectively, that is,

$$H(k) = \chi(X, L^k) \text{ and } P(k) = \chi(X, E \otimes L^k) \text{ for all } k.$$

Let k be large enough that L^k is very ample,

$$H^i(X, L^k) = 0 \text{ and } H^i(X, E \otimes L^k) = 0 \text{ for all } i > 0.$$

and $E \otimes L^k$ is globally generated. Then there are isomorphisms

$$(4.3.1) \quad H^0(X, L^k) \cong V_k^* \text{ and } H^0(X, E \otimes L^k) \cong W_k^*.$$

Proposition 4.3.1. *Under these assumptions, there is a bijection between the set of 1-parameter subgroups of G_k , up to the choice of a pair of isomorphisms (4.3.1), and the set of test configurations $(\mathcal{X}, \mathcal{L}, \mathcal{E})$ for (X, L^k, E) with \mathcal{L} very ample, $\mathcal{E} \otimes \mathcal{L}$ globally generated,*

$$R^i \pi_* \mathcal{L} = R^i \pi_*(\mathcal{E} \otimes \mathcal{L}) = 0$$

for $i > 0$ and surjective evaluation map $\pi_*(\mathcal{E} \otimes \mathcal{L}) \otimes \mathcal{L}^{-1} \rightarrow \mathcal{E}$, up to the choice of isomorphisms $H^0(X_0, L_0) \cong V_k^*$ and $H^0(X_0, E_0 \otimes L_0) \cong W_k^*$ (with π as in Definition 4.1.1).

PROOF. To see this, fix the isomorphisms (4.3.1). Then the left-hand isomorphism in (4.3.1) and the linear system $|L^k|$ determine a closed embedding

$$(4.3.2) \quad i_x: X \hookrightarrow \mathbb{P}.$$

Since $E \otimes L^k$ is globally generated, the evaluation map $H^0(E \otimes L^k) \otimes L^{-k} \rightarrow E$ is an epimorphism which, together with the right-hand isomorphism in (4.3.1), determine another one

$$(4.3.3) \quad \varphi_e: W_k \otimes L^{-k} \rightarrow E.$$

Since $\mathcal{O}_{\mathbb{P}}(1)|_X = L^k$, the embedding i_x of (4.3.2) and the epimorphism φ_e of (4.3.3) define points

$$(4.3.4) \quad x \in \text{Hilb}, \quad e \in \text{Quot}, \quad \text{with } \kappa(e) = x,$$

and so we have identifications $X = X_x$, $L^k = L_x$ and $E = E_e$. Note that each pair of isomorphisms (4.3.1) defines a point e as in (4.3.4) and the set of all points defined by the set of all the isomorphisms (4.3.1) is an orbit under the G_k -action on Quot.

Let λ be a 1-parameter subgroup of G_k , i.e. a morphism of groups

$$\lambda = (\lambda_0, \lambda_1): \mathbb{C}^* \rightarrow G_k,$$

where $\lambda_0: \mathbb{C}^* \rightarrow \text{GL}(V_k)$, $\lambda_1: \mathbb{C}^* \rightarrow \text{GL}(W_k)$. Let $f_t = \lambda_0(t) \cdot x$, $g_t = \lambda(t) \cdot e$, for $t \in \mathbb{C}^*$. As Hilb and Quot are proper over \mathbb{C} ,

$$(4.3.5) \quad f_0 = \lim_{t \rightarrow 0} f_t \in \text{Hilb}, \quad g_0 = \lim_{t \rightarrow 0} g_t \in \text{Quot}$$

exist, and f_t and g_t extend to morphisms $f: \mathbb{C} \rightarrow \text{Hilb}: t \mapsto f_t$ and $g: \mathbb{C} \rightarrow \text{Quot}: t \mapsto g_t$, which satisfy $f = \kappa \circ g$. Using now f to pull back the universal family $\mathcal{X} \subset \mathbb{P} \times \text{Hilb}$ and the invertible sheaf \mathcal{L} , and g to pull back the universal quotient sheaf \mathcal{E} , we obtain a test configuration

$$(\mathcal{X}, \mathcal{L}, \mathcal{E}) = (\mathcal{X} \times_{\text{Hilb}} \mathbb{C}, f^*\mathcal{L}, g^*\mathcal{E})$$

for (X, L^k, E) . By construction, \mathcal{L} is very ample over the canonical projection $\pi: \mathcal{X} \rightarrow \mathbb{C}$, $R^i\pi_*\mathcal{L} = R^i\pi_*(\mathcal{E} \otimes \mathcal{L}) = 0$ for $i > 0$ and the natural evaluation map $\pi_*(\mathcal{E} \otimes \mathcal{L}) \otimes \mathcal{L}^{-1} \rightarrow \mathcal{E}$ is surjective.

For the converse, let $(\mathcal{X}, \mathcal{L}, \mathcal{E})$ be a test configuration for (X, L^k, E) , with \mathcal{L} very ample over $\pi: \mathcal{X} \rightarrow \mathbb{C}$,

$$R^i\pi_*\mathcal{L} = R^i\pi_*(\mathcal{E} \otimes \mathcal{L}) = 0$$

for $i > 0$ and surjective evaluation map $\pi_*(\mathcal{E} \otimes \mathcal{L}) \otimes \mathcal{L}^{-1} \rightarrow \mathcal{E}$. These conditions clearly imply that $\pi_*\mathcal{L}$ and $\pi_*(\mathcal{E} \otimes \mathcal{L})$ are vector bundles of ranks $H(k)$ and $P(k)$, respectively. Since the central fibre (X_0, L_0, E_0) has an induced \mathbb{C}^* -action, there are also induced \mathbb{C}^* -actions on $H^0(X_0, L_0)$ and $H^0(X_0, E_0 \otimes L_0)$, so each choice of isomorphisms

$$(4.3.6) \quad H^0(X_0, L_0) \cong V_k^* \text{ and } H^0(X_0, E_0 \otimes L_0) \cong W_k^*$$

defines a 1-parameter subgroup $\lambda = (\lambda_0, \lambda_1): \mathbb{C}^* \rightarrow G_k$.

To see that $(\mathcal{X}, \mathcal{L}, \mathcal{E})$ is recovered from λ , note that the assumptions on \mathcal{L} provide a \mathbb{C}^* -equivariant closed embedding $\mathcal{X} \subset \mathbb{P}_{\mathbb{C}}(\pi_*\mathcal{L})$ of schemes over \mathbb{C} with $\mathcal{L} \cong \mathcal{O}_{\mathbb{P}_{\mathbb{C}}}(-1)|_{\mathcal{X}}$, where $\mathcal{O}_{\mathbb{P}_{\mathbb{C}}}(1)$ is the standard invertible sheaf on the projective space bundle $\mathbb{P}_{\mathbb{C}}(\pi_*\mathcal{L})$

over \mathbb{C} associated to $\pi_*\mathcal{L}$. Now, $\pi_*\mathcal{L}$ and $\pi_*(\mathcal{E} \otimes \mathcal{L})$ are trivial as vector bundles (as they are defined over \mathbb{C}), with possibly non-trivial G_k -actions and central fibres

$$(\pi_*\mathcal{L})_0 = H^0(X_0, L_0) \text{ and } (\pi_*(\mathcal{E} \otimes \mathcal{L}))_0 = H^0(X_0, E_0 \otimes L_0),$$

since \mathcal{L} and \mathcal{E} are flat over \mathbb{C} . Hence two trivialisations of these vector bundles together with the isomorphisms (4.3.6) define isomorphisms

$$(4.3.7) \quad \pi_*\mathcal{L} \cong V_k \times \mathbb{C}, \quad \pi_*(\mathcal{E} \otimes \mathcal{L}) \cong W_k \times \mathbb{C},$$

of \mathbb{C}^* -linearised vector bundles, where \mathbb{C}^* acts on V_k and W_k via λ_0 and λ_1 , and $V_k \times \mathbb{C}$ and $W_k \times \mathbb{C}$ are vector bundles over \mathbb{C} with their diagonal \mathbb{C}^* -actions. Then the left-hand isomorphism (4.3.7) induces an isomorphism of \mathbb{C}^* -linearised varieties $(\mathbb{P}_{\mathbb{C}}(\pi_*\mathcal{L}), \mathcal{O}_{\mathbb{P}_{\mathbb{C}}}(1)) \cong (\mathbb{P} \times \mathbb{C}, p^*\mathcal{O}_{\mathbb{P}}(1))$ over \mathbb{C} , where $p: \mathbb{P} \times \mathbb{C} \rightarrow \mathbb{P}$ is the canonical projection. This isomorphism together with the closed embedding $\mathcal{X} \subset \mathbb{P}_{\mathbb{C}}(\pi_*\mathcal{L})$ provide a \mathbb{C}^* -equivariant map $f: \mathbb{C} \rightarrow \text{Hilb}: t \mapsto f_t$, such that

$$(\mathcal{X}, \mathcal{L}) \cong f^*(\mathcal{X}, \mathcal{L}) \text{ as } \mathbb{C}^*\text{-linearised schemes.}$$

Furthermore, the right-hand isomorphism (4.3.7) together with the surjective evaluation map $\pi_*(\mathcal{E} \otimes \mathcal{L}) \otimes \mathcal{L}^{-1} \twoheadrightarrow \mathcal{E}$ provide a \mathbb{C}^* -equivariant surjective morphism $W_k \otimes \mathcal{L}^{-1} \twoheadrightarrow \mathcal{E}$, which induces a \mathbb{C}^* -equivariant map $g: \mathbb{C} \rightarrow \text{Quot}: t \mapsto g_t$ such that $f = \kappa \circ g$ and

$$\mathcal{E} \cong g^*\mathcal{E} \text{ as } \mathbb{C}^*\text{-linearised sheaves.}$$

In particular, $(X, L^k, E) \cong (X_x, L_x, E_e)$ where $x = f_1$ and $e = g_1$, so

$$f_t = \lambda_0(t) \cdot x \text{ and } g_t = \lambda_1(t) \cdot e \text{ for all } t \in \mathbb{C}^*,$$

because f and g are \mathbb{C}^* -equivariant, and therefore $(\mathcal{X}, \mathcal{L}, \mathcal{E})$ is recovered from λ because Hilb and Quot are projective and hence separated over \mathbb{C} , so the limits (4.3.5) are unique. This concludes the construction of the bijection between 1-parameter subgroups of G_k and test configurations. \square

Remark 4.3.2. As already mentioned, the above correspondence between test configurations for a triple (X, L^k, E) and 1-parameter subgroups of G_k generalizes a similar correspondence between test configurations for a pair (X, L^k) and 1-parameter subgroups of $\text{GL}(V_k)$ (see [24, 64]). Let (X, L, E) be as in Definition 4.1.1 and k be large enough that L^k is very ample, $H^i(X, L^k) = H^i(X, E \otimes L^k) = 0$ for all $i > 0$ and $E \otimes L^k$ is globally generated. Let $(\mathcal{X}, \mathcal{L})$ be a test configuration for (X, L^k) . Then we can lift $(\mathcal{X}, \mathcal{L})$ to a test configuration $(\mathcal{X}, \mathcal{L}, \mathcal{E})$ for (X, L^k, E) . Furthermore, if $(\mathcal{X}, \mathcal{L})$ is a product or a trivial configuration, then it can be respectively lifted to a product or a trivial configuration for (X, L^k, E) . This follows simply because a test configuration for (X, L^k) corresponds to a 1-parameter subgroup $\lambda_0: \mathbb{C}^* \rightarrow \text{GL}(V_k)$ and so a possible lift is the one associated to the 1-parameter subgroup $\lambda = (\lambda_0, \text{Id}_{W_k}): \mathbb{C}^* \rightarrow G_k$. When $(\mathcal{X}, \mathcal{L})$ is a product test configuration then λ_0 fixes the point $x \in \text{Hilb}$, defined by (X, L^k) as before. Then λ fixes $\epsilon \in \text{Quot}$, defined by (X, L^k, E) , and so the associated test configuration is a product. The statement about trivial configurations can now be readily checked.

We will now see that α -K-stability reduces to K-stability of the underlying polarised variety in the limit case $\alpha = 0$ (note that Definition 4.1.2 also makes sense for $\alpha = 0$). We will also prove that α -K-semistability for all $0 < \alpha \ll 1$ implies K-semistability

of the polarised variety (X, L) . In the following proposition, K-(semi/poly)stability for a polarised variety (see §1.3.6) is defined as first introduced by Donaldson [24] (after Tian [74]), but in the slightly different formulation of Ross&Thomas [63, Definition 2.4].

Proposition 4.3.3. *A triple (X, L, E) is α -K-(semi/poly)stable, with $\alpha = 0$, if and only if the polarised variety (X, L) is algebraically K-(semi/poly)stable. If (X, L, E) is α -K-semistable for all $0 < \alpha \ll 1$, then (X, L) is algebraically K-semistable.*

PROOF. Let $(\mathcal{X}, \mathcal{L}, \mathcal{E})$ be a test configuration for (X, L^k, E) . Note that, in the limit case $\alpha = 0$, $F_\alpha(\mathcal{X}, \mathcal{L}, \mathcal{E})$ given by (4.1.2) is the Futaki invariant $-F_1(\mathcal{X}, \mathcal{L})$, as defined by Donaldson [24, §2.1] but with opposite sign. The “if” part follows immediately from this observation. To prove the “only if” part for K-semistability let $(\mathcal{X}, \mathcal{L})$ be a test configuration for (X, L^k) . Taking $l \gg 0$ large enough that Remark 4.3.2 applies $(\mathcal{X}, \mathcal{L}^l)$ lifts to a test configuration $(\mathcal{X}, \mathcal{L}^l, \mathcal{E})$ for (X, L^{lk}, E) . Then, in the limit case $\alpha = 0$,

$$F_{lk\alpha}(\mathcal{X}, \mathcal{L}^l, \mathcal{E}) = -F_1(\mathcal{X}, \mathcal{L}^l) = -F_1(\mathcal{X}, \mathcal{L}),$$

and so the “only if” part holds for K-semistability. The “only if” parts for K-stability and K-polystability now follow from the results in Remark 4.3.2 about lifts of product and trivial configurations.

To prove the second part of the statement suppose that (X, L) is not K-semistable. Then there exists a destabilizing test configuration $(\mathcal{X}, \mathcal{L})$ for (X, L^k) and taking $l \gg 0$ we can lift $(\mathcal{X}, \mathcal{L}^l)$ to a test configuration for (X, L^{lk}, E) . Therefore, by definition, the corresponding α -invariant is negative for $0 < \alpha \ll 1$ and so (X, L, E) is α -K-unstable for these values of α . \square

4.4. Base-preserving configurations and stable vector bundles

In this section, we show that α -K-stability of a triple implies Mumford–Takemoto stability of the coherent sheaf. First, we use a special type of test configurations to reformulate the usual definition of Mumford stability for a coherent sheaf (see §1.2.4) in terms of the invariant F_2 defined in (4.1.1).

Let X be an n -dimensional projective scheme with an ample invertible sheaf L . Let E be a coherent sheaf on X which is pure of dimension n . We will use the following special type of test configurations. A *base-preserving test configuration* for (X, L, E) is a test configuration $(\mathcal{X}, \mathcal{L}, \mathcal{E})$ such that $(\mathcal{X}, \mathcal{L})$ is a trivial configuration for (X, L) , i.e. the product of (X, L) and \mathbb{C} , where \mathbb{C}^* acts trivially on (X, L) and by multiplication on \mathbb{C} (see Remark 4.3.2 for details).

Proposition 4.4.1. *Let E be a coherent sheaf over an n -dimensional projective scheme X with an ample invertible sheaf L . Suppose that E is pure of dimension n . Then E is Mumford semistable if and only if*

$$(4.4.1) \quad F_2(E_0) \leq 0,$$

for all base-preserving test configurations $(\mathcal{X}, \mathcal{L}, \mathcal{E})$ for (X, L^k, E) , for all integers $k > 0$. The sheaf E is Mumford stable if and only if the inequality (4.4.1) is strict for all non-trivial base-preserving configurations $(\mathcal{X}, \mathcal{L}, \mathcal{E})$, and Mumford polystable if

and only if the inequality (4.4.1) is strict unless $(\mathcal{X}, \mathcal{L}, \mathcal{E})$ is a product base-preserving configuration.

PROOF. It is helpful to observe that a base-preserving test configuration is equivalent to the data of a \mathbb{C}^* -linearised coherent sheaf \mathcal{E} over $X \times \mathbb{C}$ (where \mathbb{C}^* acts trivially on X and by multiplication on \mathbb{C}), which is flat over \mathbb{C} and has fibres $E_t \cong E$ for $t \in \mathbb{C} \setminus \{0\}$. This is a product configuration precisely when $\mathcal{E} = p^*E$, where $p: X \times \mathbb{C} \rightarrow X$ is the canonical projection, for a \mathbb{C}^* -action on E covering the trivial \mathbb{C}^* -action on X , and it is a trivial configuration if \mathbb{C}^* acts on E by scalar multiplication.

For the “only if” part, let $F \subset E$ be a subsheaf and $P_L(F, k) = \chi(F \otimes L^k)$ its Hilbert polynomial. Then there exists a base-preserving configurations for (X, L, E) whose central fibre is $E_0 = F \oplus E/F$ where $t \in \mathbb{C}^*$ acts as $\text{diag}(t, 1)$, i.e. with weight 1 on F and weight 0 on E/F (see e.g [71, p 43]). Recall that, since E is pure of dimension $n = \dim X$, we have

$$P_L(F, k) = k^n d_0(F) + k^{n-1} d_1(F) + \dots,$$

for all integer k with $d_0(F) > 0$, and that the slope $\mu_L(F)$ of F is defined as the quotient $\frac{d_1(F)}{d_0(F)}$. Computing now the normalized weight $F(E_0, k)$ for this degeneration, we obtain

$$F(E_0, k) = \frac{P_L(F, k)}{k P_L(E, k)} = \frac{d_0(F)}{d_0(E)} \left(k^{-1} + k^{-2} (\mu_L(F) - \mu_L(E)) + O(k^{-3}) \right).$$

In other words, $F_2(E_0)$ equals the difference of slopes $\mu_L(F) - \mu_L(E)$, up to a multiplicative positive constant, so the condition that the weak/strict inequality holds in (4.4.1) for all non-trivial base-preserving test configurations implies that E is Mumford semistable/stable.

For the converse, let \mathcal{E} be a \mathbb{C} -flat \mathbb{C}^* -linearised coherent sheaf over $X \times \mathbb{C}$. Choosing $k \gg 0$, we obtain a surjective map

$$(4.4.2) \quad \epsilon: W_k \otimes L^{-k} \twoheadrightarrow E$$

as in §4.3, where $W_k^* \cong H^0(X, E \otimes L^k)$. We can choose k such that (4.4.2) is surjective for any fiber \mathcal{E}_t of \mathcal{E} , and so the 1-parameter subgroup $\lambda_k: \mathbb{C}^* \rightarrow GL(W_k)$ determined by the action on the central fiber E_0 splits $W_k^* \cong H^0(X, E_0 \otimes L^k)$ into eigenspaces

$$W_k = \bigoplus_{n \in \mathbb{Z}} V_n,$$

where \mathbb{C}^* acts with weight n on V_n and only a finite number of the V_n are non trivial. The eigenspaces V_n give a weight filtration $V_{\leq n} = \bigoplus_{j \leq n} V_j$ of W_k . Their images $F_{\leq n}$ under (4.4.2) give a filtration of E such that E_0 is the graded sheaf (see [37, Lemma 4.4.3])

$$E_0 \cong \bigoplus_{n \in \mathbb{Z}} F_{\leq n} / F_{\leq n-1}.$$

The weight w_l of the \mathbb{C}^* -action on $\det H^0(X, E_0 \otimes L^l)$ for $l \gg k$ is then given by

$$\begin{aligned} w_l &= \sum_{n \in \mathbb{Z}} (-n) P_L(F_{\leq n}/F_{\leq n-1}, l) \\ &= \sum_{n \in \mathbb{Z}} (-n) (P_L(F_{\leq n}, l) - P_L(F_{\leq n-1}, l)) \\ &= \sum_{n \in \mathbb{Z}} P_L(F_{\leq n}, l) \end{aligned}$$

(see [37, Lemma 4.4.4]). Thus, the normalized weight $F(E_0, l)$ equals

$$\begin{aligned} F(E_0, l) &= \sum_{n \in \mathbb{Z}} \frac{P_L(F_{\leq n}, l)}{l \cdot P_L(E, l)} \\ &= \sum_{n \in \mathbb{Z}} \frac{d_0(F_{\leq n})}{d_0(E)} (l^{-1} + l^{-2} (\mu_L(F_{\leq n}) - \mu_L(E)) + O(l^{-3})). \end{aligned}$$

This equality together with equation (4.1.1) imply that if E is Mumford–Takemoto semistable/stable, then the weak/strict inequality holds in (4.4.1) for all non-trivial base-preserving test configurations $(\mathcal{X}, \mathcal{L}, \mathcal{E})$.

Finally, we will prove the result about polystability. If the equality in (4.4.1) holds only for product base-preserving configurations, then we have seen that E must be Mumford–Takemoto semistable. If E is strictly semistable, there exists a proper subsheaf $E' \subset E$ with $\mu_L(F) = \mu_L(E)$. Since the degeneration associated to F has zero weight, $E \cong F \oplus E/F$ and therefore we conclude by iteration that E must be Mumford polystable. For the converse, suppose that E is Mumford polystable, i.e.

$$E \cong \bigoplus_i F_i$$

where F_i are Mumford stable sheaves with $\mu(F_i) = \mu(E)$. Given a base-preserving configuration for which we have an equality in (4.4.1), the corresponding central fibre must be equal to

$$E_0 \cong \bigoplus_j F_j$$

i.e. $E_0 \cong E$, so this configuration must be a product configuration. \square

Corollary 4.4.2. *Let (X, L, E) be a triple consisting of a projective scheme of dimension n , an ample invertible sheaf and a coherent sheaf, pure of dimension n . Then we have the following implications for all $\alpha > 0$: if (X, L, E) is α -K-(semi/poly)stable, then E is Mumford (semi/poly)stable with respect to L .*

PROOF. Immediate from Proposition 4.4.1 and the fact that for any base-preserving configuration $(\mathcal{X}, \mathcal{L}, \mathcal{E})$ for (X, L^k, E) , the central fibre of $(\mathcal{X}, \mathcal{L})$ is $(X_0, L_0) = (X, L^k)$ with the trivial \mathbb{C}^* -action, so $F_i(\mathcal{O}_{X_0}) = 0$ for all $i > 0$ and hence (4.1.2) reduces to

$$F_\alpha(\mathcal{X}, \mathcal{L}, \mathcal{E}) = -\alpha F_2(E_0),$$

whose sign does not depend on the actual value of $\alpha > 0$. \square

Another consequence of Proposition 4.4.1 is the following relation between the existence of solutions to the coupled equations (3.0.7) and base-preserving configurations for the triple.

Corollary 4.4.3. *Let (X, L, E) be a triple consisting of a smooth compact polarized variety (X, L) and a holomorphic vector bundle E over X with rank r and degree zero. Suppose that there exists a solution (ω, H) to the coupled equations (3.0.7) with $\omega \in c_1(L)$. Then, given a positive integer $k > 0$, any $\alpha > 0$ and a base-preserving test configuration $(\mathcal{X}, \mathcal{L}, \mathcal{E})$ for the triple (X, L^k, E) , we have*

$$F_{k\alpha}(\mathcal{X}, \mathcal{L}, \mathcal{E}) \geq 0,$$

with equality only if the test configuration is a product.

PROOF. If there exists a solution to (3.0.7) then E is Mumford–Takemoto polystable by the Hitchin–Kobayashi correspondence for vector bundles and so the statement holds from Proposition 4.4.1 as in Corollary 4.4.2. \square

4.5. A conjecture for degree zero holomorphic vector bundles

In this section we formulate a conjecture relating the notion of α -K-stability for triples (X, L, E) defined in section §4.1 with the existence of solutions to the coupled equations (3.0.7). We will suppose that (X, L) is a smooth compact polarized variety and E is a holomorphic vector bundle over X with rank r and degree zero. Recall from Proposition 4.2.1 that the numerical invariant F_α associated to any test configuration with smooth central fiber is, up to multiplicative constant factor, equal to the moment map (2.3.3) at the generator of the $S^1 \subset \mathbb{C}^*$ action and that the zero locus of this moment map corresponds to the solutions of the coupled equations. Once more this fits in the general picture of the Kempf–Ness Theorem explained in §1.1 if we view a test configuration as a 1-parameter subgroup (see §4.3) and the associated α -invariant as the corresponding Mumford weight (see §1.1.3 and cf. (1.1.9)). This fact motivates Definition 4.1.2 and leads to think that the notion of α -K-stability should play an important role in the existence problem for the coupled equations (3.0.7).

Concerning direct relations between the equations (3.0.7) and the notion of α -K-stability, we have proved in Corollary 4.2.2 that if there exists a solution to the coupled equations (3.0.7) with positive coupling constants $\alpha_0, \alpha_1 \in \mathbb{R}$, then $F_\alpha = 0$ for any product test configuration with $\alpha = \frac{r\pi^2\alpha_1}{\alpha_0}$. Under the same hypothesis, we have proved in Corollary 4.4.3 that $F_\alpha \geq 0$ for any base preserving test configuration, with equality only if the test configuration is a product. We have also proved in Proposition 4.3.3 that in the limit case $\alpha \rightarrow 0$ the notion of α -K-stability is equivalent to the notion of algebraic K-stability of the underlying polarized manifold.

These facts provide supporting evidence for the following conjecture.

Conjecture 4.5.1. *If there exists a Kähler metric $\omega \in c_1(L)$ and a Hermitian metric H on E satisfying the coupled equations (3.0.7) with coupling constants α_0, α_1 then the triple (X, L, E) is α -K-polystable.*

Remark 4.5.2. The previous conjecture for the limit case $\alpha \rightarrow 0$, i.e. for the cscK problem, has been verified mainly thanks to the work of Tian [74], Donaldson [25], Mabuchi [49] and Stoppa [67].

This conjecture is a difficult problem even in complex dimension 1. Note that in this case the coupled equation are a system in separated variables independent of the coupling constants. In fact, they reduce to the HYM equation for the bundle due to existence of cscK metrics in any compact Riemann surface. When $\dim_{\mathbb{C}} X = 1$, Conjecture 4.5.1 reduces to prove that Mumford–Takemoto stability of the bundle implies α -K-stability of the triple for any $\alpha > 0$. The algebro-geometric approach to the proof of this fact seems to be very difficult in the light of previous work of R. Thomas and J. Ross in [64]. A more likely approach, valid for arbitrary dimension of X , seems to be provided by Donaldson’s methods in [25] or Mabuchi’s methods in [49]. The main issue to adapt the arguments in [25] or [49] will be to prove that α -K-stability arises as a suitable asymptotic version of a Chow-type stability notion in the relative quot scheme considered in Section §4.3. The verification of this stability issue will provide a strong evidence of the fact that (3.0.7) are gauge theoretic equations related to the algebro-geometric moduli problem for triples (X, L, E) . As an immediate future project we will try to prove that the existence of solutions to (3.0.7) implies the α -K-semistability of the triple adapting Donaldson’s methods in [25]. This conjecture was first verified for the limit case $\alpha = 0$, i.e. for the cscK problem, after the work of S. Zhang, H. Luo and S.K. Donaldson.

Concerning the converse of the statement in Conjecture 4.5.1, as remarked by J. Ross and R. Thomas [63], a recent example in [2] suggests that the algebraic notion of K-stability defined in [24] for a polarised variety may not be sufficient to guarantee the existence of a cscK metric and so a stronger notion may be required. Since the α -K-stability of a triple (X, L, E) for $\alpha = 0$ coincides with the K -stability of (X, L) , it seems equally necessary to strengthen the α -K-stability condition introduced in Definition 4.1.2, so that the converse of the statement in Conjecture 4.5.1 may actually hold. This can be done exactly as in [63] for polarised varieties, by requiring the inequality (4.1.3) in Definition 4.1.2 for configurations $(\mathcal{X}, \mathcal{L}, \mathcal{E})$, where \mathcal{L} may now be an ample \mathbb{R} -divisor.

To finish this chapter, let us analyze two particularly simple situations, when X is a Riemann surface and when E is the trivial bundle of rank r over a toric complex surface, where the α -K-polystability of the triple implies the existence of solutions to the equations (3.0.7). Recall from §4.4, Corollary 4.4.2, that α -K-(semi/poly)stability implies Mumford (semi/poly)stability for any $\alpha > 0$. Due to the Theorem of Donaldson–Uhlenbeck–Yau, the latter stability condition is equivalent to the existence of solutions to the Hermitian–Yang–Mills equation (1.2.5) for each Kähler metric in $c_1(L)$.

Suppose first that X is a Riemann surface. Given a positive real constant $\alpha > 0$, if the triple (X, L, E) is α -K-polystable in the sense of Definition 4.1.2 then there exists a solution to the HYM equation for any Kähler metric $\omega \in c_1(L)$. Choose ω to be the unique cscK metric in this Kähler class and H the associated Einstein–Hermitian metric on E . Then (ω, H) is a solution to (3.0.7) since the term $\Lambda_{\omega}^2(F_H \wedge F_H)$ in (3.0.7) vanishes. We conclude that if $\dim_{\mathbb{C}} X = 1$ and the triple (X, L, E) is α -K-polystable

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with $\alpha > 0$, then there exists a solution to (3.0.7), and so the converse of the statement in Conjecture 4.5.1 is verified.

Suppose now that $E \cong \oplus^r \mathcal{O}_X$ is the trivial bundle and (X, L, E) is α -K-polystable for some $\alpha \geq 0$. We claim that (X, L) is algebraically K-polystable. Let $(\mathcal{X}, \mathcal{L})$ be a test configuration for (X, L^k) . Then, by its very definition, this test configuration defines one for the triple (X, L^k, \mathcal{O}_X) and so one for $(X, L, \oplus^r \mathcal{O}_X)$ that we denote by $(\mathcal{X}, \mathcal{L}, \mathcal{E})$. Then, if (X_0, L_0, E_0) is its central fiber we have that $E_0 \cong \oplus^r \mathcal{O}_{X_0}$ as \mathbb{C}^* -linearized sheaves over X_0 . Then it is immediate to check that the normalized weight $F(E_0, k)$ satisfies

$$F(E_0, k) = \frac{w_k(E_0)}{k P_{L_0}(E_0, k)} = \frac{rw_k(\mathcal{O}_{X_0})}{kr P_{L_0}(\mathcal{O}_{X_0}, k)} = F(\mathcal{O}_{X_0}, k)$$

and so we have an equality

$$0 \leq F_{lk\alpha}(\mathcal{X}, \mathcal{L}^l, \mathcal{E}) = -F_1(\mathcal{X}, \mathcal{L}^l) = -F_1(\mathcal{X}, \mathcal{L}).$$

Thus, (X, L) is algebraically K-semistable. If the α -invariant vanishes then $(\mathcal{X}, \mathcal{L}, \mathcal{E})$ is a product and so is $(\mathcal{X}, \mathcal{L})$. We conclude that (X, L) is K-polystable as claimed. Suppose in addition that the Yau–Tian–Donaldson conjecture is satisfied for (X, L) (e.g. (X, L) is a toric polarized surface or $L^k \cong K_X$, $0 \geq k \in \mathbb{Z}$ [8, 26]). Then, by the previous reasoning, α -K-polystability of $(X, L, \oplus^r \mathcal{O}_X)$ implies the existence of a cscK metric $\omega \in c_1(L)$. Thus, the pair (ω, A) , with A any flat connection on $\oplus^r \mathcal{O}_X$, satisfies the coupled equations (3.0.7) for any α and we conclude the following. If the triple $(X, L, \oplus^r \mathcal{O}_X)$ is α -K-polystable and (X, L) satisfies the Yau–Tian–Donaldson conjecture, then there exists a solution to the coupled equations (3.0.7). If X is a toric complex surface, Donaldson has proved the Yau–Tian–Donaldson conjecture [26] and so α -K-polystability of $(X, L, \oplus^r \mathcal{O}_X)$ implies the existence of solutions to (3.0.7).

CHAPTER 5

Examples and cscK metrics on ruled manifolds

This chapter contains some examples of solutions to (0.0.3). In §5.4 we also discuss how the existence of solutions to the limit case $\alpha_0 = 0$ can be applied, using results of Y. J. Hong in [36], to obtain cscK metrics on ruled manifolds.

In §5.1 we discuss the case of vector bundles over Riemann surfaces and projectively flat bundles over Kähler manifolds satisfying a topological constraint. These are the situations considered in the moduli constructions for pairs in [60] and [66]. In both cases, the gauge theoretical equations equivalent to the moduli problem provide particular solutions to the coupled equations (0.0.3). In this examples the equations (0.0.3) reduce to the limit case $\alpha_1 = 0$ (cscK equation and HYM equation) and we give a geometric explanation of this fact.

In §5.2 we consider homogeneous Hermitian holomorphic bundles E over homogeneous Kähler manifolds and show the existence of invariant solutions to (0.0.3) in any Kähler class provided that the representation defining E is irreducible. The invariant solutions arise as (simultaneous) solutions to the limit systems $\alpha_0 = 0$ and $\alpha_1 = 0$ in (0.0.3). When the complex dimension of the base is 2 the invariant solutions satisfy in addition the Kähler analogue of the Einstein–Yang–Mills equation (see equation (2.4.6)) considered in §2.4.

In §5.3 we discuss some (well known) examples of stable bundles over Kähler–Einstein manifolds without infinitesimal automorphisms where Theorem 0.0.1 applies. This provide our first examples of solutions in which the Kähler metric is not cscK, showing that the coupled equations (0.0.3) are more than the cscK equation and the HYM equation together.

In §5.4 we discuss the relation of the limit $\alpha_0 = 0$ in (0.0.3) with the cscK equation on ruled manifolds as mentioned above.

5.1. Riemann surfaces and projectively flat bundles

Let E^c be a holomorphic principal G^c bundle over a polarized compact Riemann surface (X, L) . We fix a maximal compact subgroup $G \subset G^c$. In this case the coupled equations (3.0.7), for a G -reduction H and a Kähler metric ω , split into the system in separated variables

$$(5.1.1) \quad \left. \begin{array}{l} F_H = z \cdot \omega \\ S_\omega = \hat{S} \end{array} \right\},$$

since $\dim_{\mathbb{C}} X = 1$ and the term $(F_H \wedge F_H)$ vanishes. Recall that F_H denotes the curvature of the Chern connection A of H and that $\hat{S} \in \mathbb{R}$ and $z \in \mathfrak{z}$ are constant over X . Then, the solutions to the coupled equations (0.0.3) are given by pairs (ω, H) , where

ω is the unique cscK metric in $c_1(L)$ modulo automorphisms and H is a reduction such that A is a HYM connection (1.2.5). Due to the Hitchin-Kobayashi correspondence, examples of solutions satisfying (5.1.3) are given by polystable G^c -bundles over X . When L is a multiple of the canonical bundle of X , Pandharipande [60] used Geometric Invariant Theory to compactify the moduli space of pairs (X, E) consisting of a smooth curve of genus greater than one and a semistable vector bundle over the curve. In [60, Proposition 8.2.1] the author proves that such a pair is GIT stable if and only if E is Mumford stable. Therefore, the gauge theoretical equation corresponding to this moduli construction is simply the HYM equation. The latter is equivalent to (5.1.1) since there is always a cscK metric in $c_1(L)$.

Let now (X, ω) be a compact Kähler manifold of arbitrary dimension. Suppose that E is a projectively flat vector bundle over X satisfying the topological constraint

$$(5.1.2) \quad c_1(E) = -\frac{r\lambda}{2\pi}[\omega],$$

where $\lambda \in \mathbb{R}$ is determined by the first Chern class of the bundle and the Kähler class $[\omega]$. Then, doing a conformal change if necessary, there exists a Hermitian metric H on E which is a solution to (see [40, Corollary 2.7])

$$(5.1.3) \quad F_H = z \cdot \omega,$$

where $z = i\lambda \text{Id} \in \mathfrak{z}$. Therefore, as can be readily checked from the equations, the pair (ω, H) is a solution to (3.0.7) if and only if ω is a cscK metric. We conclude that, when E is projectively flat and (5.1.2) is satisfied, there exists a solution to (3.0.7) if $[\omega]$ admits a cscK metric. Suppose that

$$X \cong \mathbb{C}^n/\Gamma$$

is a complex torus given by a lattice Γ in \mathbb{Z}^{2n} . In this case, examples of holomorphic vector bundles E admitting a projectively flat Hermitian structure H are given by representations of a central extension of Γ into $U(r) \subset GL(r, \mathbb{C})$. It follows from [40, Theorem 7.54] that $\pm c_1(E)$ is a Kähler class admitting a cscK metric ω so E satisfies the constraint (5.1.2).

In [66], G. Schumacher and M. Toma build a moduli space of (non-uniruled) polarized Kähler manifolds equipped with isomorphism classes of stable vector bundles via the method of versal deformations. In this work the authors were able to endow the moduli space with a Kähler metric, under the cohomological constraint (5.1.2), when the base is Kähler Einstein and the bundle is projectively flat. The gauge theoretical equations corresponding to this moduli construction are therefore equivalent to (5.1.1) and so, by the discussion in the previous paragraph, they provide particular solutions to the coupled equations (3.0.7). Recall that the cscK equation and the KE equation are equivalent, using Hodge theory, if the class of the polarization is a multiple of the first Chern class of X .

In the two situations discussed the coupled equations (3.0.7) admit solutions coming from the system in separated variables (5.1.1). There is an underlying geometric reason for this fact provided by the extended gauge group $\tilde{\mathcal{G}}$ (2.1.1) associated to a solution (ω, H) to (5.1.3) (in the case of curves we assume stability of the bundle in order to

have such solution). The point is that the Chern connection A of H determines a Lie algebra splitting of the short exact sequence

$$0 \rightarrow \text{Lie } \mathcal{G} \rightarrow \text{Lie } \tilde{\mathcal{G}} \rightarrow \text{Lie } \mathcal{H} \rightarrow 0,$$

where $\text{Lie } \mathcal{G}$ and $\text{Lie } \mathcal{H}$ denote, respectively, the Lie algebras of the gauge of the reduction H and the group of Hamiltonian symplectomorphisms in (X, ω) . The splitting is given by the Lie algebra homomorphism

$$(5.1.4) \quad \Gamma: \text{Lie } \mathcal{H} \cong C^\infty(X)/\mathbb{R} \rightarrow \text{Lie } \tilde{\mathcal{G}}: \phi \rightarrow \theta_A^\perp(\eta_\phi) - z\phi,$$

where $i_{\eta_\phi}\omega = d\phi$ and θ_A^\perp denotes the horizontal lift with respect to the connection A . To see this note that

$$\begin{aligned} [\Gamma(\phi_1), \Gamma(\phi_2)] &= [\theta_A^\perp(\eta_{\phi_1}) - z\phi_1, \theta_A^\perp(\eta_{\phi_2}) - z\phi_2] \\ &= \theta_A^\perp([\eta_{\phi_1}, \eta_{\phi_2}]) - z\{\phi_1, \phi_2\} + (F_A - z\omega)(\eta_{\phi_1}, \eta_{\phi_2}) \\ &= \Gamma(\{\phi_1, \phi_2\}) + (F_A - z\omega)(\eta_{\phi_1}, \eta_{\phi_2}), \end{aligned}$$

where $\{\phi_1, \phi_2\}$ is the Poisson bracket in $C^\infty(X)/\mathbb{R}$ given by ω . Note that this homomorphism does not extend in general to the Lie algebra of the group of diffeomorphisms of X . Therefore, when $\dim_{\mathbb{C}} X = 1$ or E is projectively flat, the coupled system (3.0.7) is able to admit ‘decoupled’ solutions due to the fact that $\text{Lie } \tilde{\mathcal{G}}$ is a semidirect product of $\text{Lie } \mathcal{G}$ and $\text{Lie } \mathcal{H}$.

5.2. Homogeneous bundles over homogeneous Kähler manifolds

For the basic material on this topic we refer to [11] and [40]. Let X be a compact homogeneous Kählerian manifold (i.e. admitting a Kähler metric), of a compact group G . In other words $X = G/G_o$, for a closed subgroup $G_o \subset G$, equipped with the canonical G -invariant complex structure (see Remark 8.99 in [11]). In this situation, homogeneous holomorphic vector bundles E of rank r over X are in one to one correspondence with representations of G_o on $GL(r, \mathbb{C})$. For any invariant Kähler metric ω on X there exists a unique G -invariant Hermitian–Yang–Mills unitary connection A provided that the representation inducing E is irreducible (see [40, Proposition 6.1]). Moreover, for any such choice of invariant metric and connection, the function $\Lambda_\omega^2 \text{tr } F_A \wedge F_A$ on X is G -invariant and hence constant. Therefore, it turns that A satisfies the stronger equation

$$(5.2.1) \quad \left. \begin{array}{l} \Lambda_\omega F_A = \mathbf{i} \lambda \text{Id} \\ \Lambda_\omega^2 \text{tr } F_A \wedge F_A = -\frac{4\hat{c}}{(n-1)!} \end{array} \right\},$$

where $\hat{c} \in \mathbb{R}$ is as in (2.2.6) and $\lambda \in \mathbb{R}$ is determined by the first Chern class of the bundle and the Kähler class $[\omega]$. The system (5.2.1) corresponds to the limit

$$\alpha_0 \rightarrow 0$$

in (3.0.7). Let us fix a pair of arbitrary coupling constants $\alpha_0, \alpha_1 > 0$ and a homogeneous holomorphic vector bundle E over X associated to an irreducible representation. Then, any Kähler class on X determines a unique G -invariant solution (ω, A) to the coupled equations with coupling constants α_0 and α_1 . To see this note that every de Rham class on X (in particular, every Kähler class) contains a unique G -invariant

representative, obtained from an arbitrary representative by averaging. Trivially, the scalar curvature of any G -invariant Kähler metric is constant. Therefore, the unique G -invariant solution of (3.0.7) arises as a simultaneous solution of the cscK equation and (5.2.1), corresponding to the limit cases $\alpha_0 = 0$, and $\alpha_1 = 0$.

We will now give some examples of solutions to the Einstein–Yang–Mills type equations (2.4.6) in the case of complex surfaces. Let X be a compact complex surface and E a homogeneous holomorphic vector bundle over X of rank r and zero degree. Then, the condition to be Hermitian–Yang–Mills for a unitary connections A in E is equivalent to the Anti-Self-Dual (ASD) equation and so (2.4.6) can be restated as

$$(5.2.2) \quad \left. \begin{aligned} F_A^+ &= 0 \\ \alpha_0(\rho_\omega - c'\omega) &= \alpha_1(2(\Lambda_\omega F_A, F_A) - \Lambda_\omega(F_A \wedge F_A) - c''\omega) \end{aligned} \right\},$$

for real constants c' and $c'' \in \mathbb{R}$. Since $\dim_{\mathbb{C}} X = 2$, if A is ASD with respect to ω we have that (see (2.4.3))

$$\begin{aligned} 2(\Lambda_\omega F_A, F_A) - \Lambda_\omega(F_A \wedge F_A) &= -\Lambda_\omega(F_A \wedge F_A) \\ &= \Lambda_\omega(|F_A|^2 \cdot \omega^2) \\ &= |F_A|^2 \cdot \omega, \end{aligned}$$

where the pointwise norm of the curvature is taken with respect to ω . Suppose now that (X, ω) is a compact homogeneous Kähler-Einstein surface, i.e. it is a complex torus or is simply-connected (see Corollary 8.98 in [11]). We suppose also that the vector bundle E is homogeneous and associated to an irreducible representations of G_o on $GL(r, \mathbb{C})$. Consider the Hermitian metric H on E whose Chern connection A is HYM. Since the function $|F_A|^2$ is constant over X by invariance, the pair (ω, A) satisfies the equation (5.2.2). Relying in the discussion of §2.4 the invariant Riemannian metric associated to (ω, A) in the total space of the $U(r)$ -bundle of frames of (E, H) over X is Einstein.

5.3. Stable bundles over Kähler-Einstein manifolds

We see now some examples where Theorem 3.2.4 can be applied. The following will provide our first examples of solutions to the coupled equations (3.0.7) in which the Kähler metric is not cscK.

Let X be a high degree hypersurface of \mathbb{P}^3 . By theorems of T. Aubin and S. T. Yau (see e.g. Theorem 11.7 in [11]), there exists a unique Kähler-Einstein metric $\omega \in c_1(X)$ with negative (constant) scalar curvature. Moreover, $c_1(X) < 0$ implies that the group of automorphisms of the complex manifold X is discrete (see [11, Proposition 2.138]). Let E be a smooth $SU(2)$ -principal bundle over X with second Chern number

$$k = \frac{1}{8\pi^2} \int_X \text{tr } F_A \wedge F_A \in \mathbb{Z},$$

where A is a connection on E . When k is sufficiently large, the moduli space M_k of Anti-Self-Dual (ASD) connections A on E with respect to ω is non-empty (see [27, Section 10.1.14]). Moreover, taking k large enough, M_k is non-compact but admits a compactification. Let A be a connection that determines a point in M_k . Then, A is

irreducible and so we can apply Theorem 3.2.4 obtaining solutions $(\omega_\alpha, A_\alpha)$ to (0.0.3) for nonzero values of the coupling constants α_0, α_1 and small ratio

$$\alpha = \frac{\alpha_1}{\alpha_0}.$$

We claim that, if the pointwise norm

$$(5.3.1) \quad |F_{A_0}|_{\omega_0}^2 : X \rightarrow \mathbb{R}$$

of the initial HYM connection $A_0 = A$ with respect to the KE metric $\omega_0 = \omega$ is not constant, then ω_α is not cscK for $0 < \alpha \ll 1$. To see this note that $(\omega_\alpha, A_\alpha)$ approaches uniformly to (ω_0, A_0) (see Lemma 3.2.2) and so

$$\lim_{\alpha \rightarrow 0} \| |F_{A_\alpha}|_{\omega_\alpha}^2 - |F_{A_0}|_{\omega_0}^2 \|_{L^\infty} \rightarrow 0.$$

Therefore, if (5.3.1) is not constant then $|F_{A_\alpha}|_{\omega_\alpha}^2$ is not constant for small α and our claim follows from

$$S_{\omega_\alpha} = \frac{c}{\alpha_0} - \alpha \cdot \Lambda_{\omega_\alpha}^2 (F_{A_\alpha} \wedge F_{A_\alpha}) = \frac{c}{\alpha_0} + \alpha \cdot |F_{A_\alpha}|_{\omega_\alpha}^2,$$

where $c \in \mathbb{R}$. This last equation is satisfied since $(\omega_\alpha, A_\alpha)$ is a solution to (0.0.3). To choose an ASD connection for which (5.3.1) is not a constant, we consider a sequence of ASD connections $\{A^l\}_{l=0}^\infty$ on M_k approaching to a point in the boundary of the compactification. When $l \gg 0$ the connections A_l start bubbling. This bubbling is reflected in the fact that the function (5.3.1) becomes more and more concentrated in a finite number of points of the manifold. Therefore, eventually, we obtain an ASD irreducible connection for which (5.3.1) is not a constant. To be more precise, recall that any point in the boundary of the compactification of M_k is given by an ideal connection (see [27, Definition 4.4.1]), i.e. an unordered d -tuple (p_1, \dots, p_d) of points on X and a connection A_∞ in M_{k-d} , the moduli space of ASD connections on suitable smooth $SU(2)$ -bundle E_{k-d} with second Chern number $k-d$. If $[A_l] \rightarrow [A_\infty]$ then, for any continuous function f on X (see [27, Theorem 4.4.4])

$$(5.3.2) \quad \lim_{l \rightarrow \infty} \int_X f \operatorname{tr} F_{A_l} \wedge F_{A_l} = \int_X f \operatorname{tr} F_{A_\infty} \wedge F_{A_\infty} + 8\pi^2 \sum_{m=1}^d f(p_m)$$

We take A_∞ in M_{k-d} with $d > 0$. If $|F_{A_l}|_\omega^2$ is constant for every l , using the previous equation and the equality

$$|F_{A_l}|_\omega^2 \cdot \omega^2 = \operatorname{tr} F_{A_l} \wedge F_{A_l}$$

we obtain that $d = 0$, and so a contradiction (e.g. take in (5.3.2) a sequence of test functions f_n approaching the characteristic function of X^0 on X).

The hypothesis of Theorem 3.2.4 hold in much more generality. A family of examples generalizing the previous one is provided by stable bundles over Kähler–Einstein manifolds. Let (X, L) be a compact polarised manifold whose first Chern class $c_1(X)$ is positive negative or zero and

$$c_1(L) = \lambda c_1(X), \quad \lambda \in \mathbb{Z}.$$

When $c_1(X) < 0$ (e.g. if X is a high degree hypersurface of \mathbb{P}^m), X has finite group of automorphisms and there exists, by the already mentioned results of Aubin and Yau, a unique Kähler–Einstein metric $\omega \in c_1(L)$. If $c_1(X) = 0$ there exists, by S. T.

Yau solution to Calabi's Conjecture (see e.g. [11, Theorem 11.7]), a unique Ricci-flat metric on $c_1(L)$. The dimension of the group of automorphisms of such manifolds is equal to its first Betti number (see [11, Remark 11.22]), hence simply connected ones (e.g. $K3$ surfaces) are examples of complex Ricci-flat manifolds with finite group of automorphisms. If $c_1(X) > 0$ it is still unknown in general if X carries or not a Kähler-Einstein metric. Let us restrict to the case

$$X = \mathbb{P}^2 \# m\overline{\mathbb{P}}^2,$$

the complex surface obtained by blowing up \mathbb{P}^2 at m generic points (see [78]). If we take m such that $3 < m < 8$ then $c_1(X) > 0$, X has finite automorphism group (see Remark 3.12 in [76]) and it was proved in [78] that X admits a Kähler-Einstein metric.

On the other hand, given a compact polarised manifold (X, L) (without any assumption on $c_1(X)$), an asymptotic result of Maruyama [52] states that L -stable vector bundles E over X of rank r exist for $r > \dim X > 2$ and

$$(5.3.3) \quad c_2(E) \cdot c_1(L)^{n-2} \gg 0.$$

If X has finite group of automorphisms and it is endowed with a Kähler-Einstein metric $\omega \in c_1(L)$ as before, we can apply Theorem 3.2.4.

5.4. CscK metrics on ruled manifolds

We now briefly discuss the relation between equation (5.2.1), given by the limit

$$\alpha_0 \rightarrow 0$$

in (3.0.7), and the existence of solutions to the cscK equation on ruled manifolds. We use existence results of Y. J. Hong [35, 36].

Let E be a holomorphic stable vector bundle with degree zero over a compact Kähler manifold (X, ω) with constant scalar curvature. Let H be the (unique) Hermitian metric on E whose Chern connection A is HYM, given by Donaldson–Uhlenbeck–Yau Theorem ([21][79]). Let $\mathbb{P}(E)$ be the projectivized bundle of E and L the tautological bundle over $\mathbb{P}(E)$. We consider the induced connection in L^* by A and denote its curvature by $F_{A_{L^*}}$. The 2-form

$$\frac{i}{2\pi} F_{A_{L^*}}$$

is nondegenerate on the fibres since it induces the Fubini-Study metric. Then, for k large enough

$$\widehat{\omega}_k = \frac{i}{2\pi} F_{A_{L^*}} + k\pi^*\omega$$

is a Kähler metric in $\mathbb{P}(E)$. When X has finite automorphisms group Y.J. Hong proved in [35] that $[\widehat{\omega}_k]$ has a cscK metric for $k \gg 0$, using a deformation argument. Let denote by $\tilde{\mathcal{G}}$ the extended gauge group of the $PU(r)$ bundle of frames associated to the reduction H and the symplectic form ω . We denote by $\tilde{\mathcal{G}}_I \subset \tilde{\mathcal{G}}$ the subgroup preserving the connection A . In [36] he was able to remove the assumption on $\text{Aut } X$ under the additional assumptions (see [36, Definition I.A]):

$$\tilde{\mathcal{G}}_I \subset \text{Aut } \mathbb{P}(E)$$

finite and

$$(5.4.1) \quad \Lambda_\omega^2 (\operatorname{tr} F_A \wedge \operatorname{tr} F_A + \operatorname{tr} F_A \wedge \rho_\omega + F_A \wedge F_A) = \text{const.},$$

that, by our assumption on the degree of E , reduces to

$$\Lambda_\omega^2 \operatorname{tr} F_A \wedge F_A = -\frac{4\hat{c}}{(n-1)!} \in \mathbb{R}.$$

The condition (5.4.1) appears when one splits the linearization of the cscK equation on $\mathbb{P}(E)$ into vertical and horizontal parts with respect to the connection A .

Then, we conclude that the existence of a solution to (5.2.1) when $c_1(E) = 0$ and \mathcal{G}_I is finite is a sufficient condition for the existence of a cscK metric in $[\widehat{\omega}_k]$ for $k \gg 0$ (see [36, Theorem III.A]). We would like to study further this relation trying to prove that the existence of solutions to the coupled equations for small $\frac{\alpha_1}{\alpha_0} > 0$ implies the existence of constant scalar curvature Kähler metrics on $\mathbb{P}(E)$ with Kähler class $c_1(L_k)$ for large k . This would provide a generalization of Hong's results in [36].

CHAPTER 6

Solutions on \mathbb{C}^2

The idea of this chapter was suggested to the author by Professor Olivier Biquard and developed during a visit to the Institut de Mathématiques de Jussie under his supervision, for which the author is very grateful. Let us consider \mathbb{C}^2 with the standard complex structure endowed with the trivial smooth principal bundle $E = \mathbb{C}^2 \times SU(2)$. Given a real coupling constant $\alpha \in \mathbb{R}$ and any integer $k \in \mathbb{Z}$ we will look for pairs (ω, A) , where ω is a Kähler metric on \mathbb{C}^2 and A is an $SU(2)$ -connection on E , satisfying the coupled equations (0.0.3), that we can rewrite in this context using the Hodge star operator $*_{\omega} : \Omega^k \rightarrow \Omega^{4-k}$ as

$$(6.0.2) \quad \left. \begin{aligned} F_A^{+\omega} &= 0 \\ S_{\omega} + \alpha *_{\omega} \operatorname{tr} F_A \wedge F_A &= \alpha 8\pi^2 k \end{aligned} \right\}.$$

The connections and metrics that we will consider have special asymptotic properties at infinity. In fact, any connection A will be required to satisfy the topological constraint

$$(6.0.3) \quad \int_{\mathbb{C}^2} \operatorname{tr} F_A \wedge F_A = 8\pi^2 k, \quad k \in \mathbb{Z}.$$

The fixed asymptotic behavior at infinity is analogue in this non-compact case of fixing the Kähler class and the topological type of the bundle. We will show that in \mathbb{C}^2 there is no obstruction for finding solutions for arbitrary coupling constant α . Solutions with $\alpha = 0$ are provided by pairs (ω_0, A) , where ω_0 is the euclidean metric $\frac{i}{2} \sum_j dz_j \wedge d\bar{z}_j$ and A is any k -instanton, i.e. any solution to the Anti-Self-Duality equation

$$(6.0.4) \quad F_A^{+\omega_0} = 0$$

with fixed topological invariant (6.0.3). In this section we will show that for each $\alpha \in \mathbb{R}$ and $k \in \mathbb{Z}$ there exists a solution $(\omega_{\alpha}, A_{\alpha})$ of the coupled equations (6.0.2) with coupling constant α and fixed topological invariant $8\pi^2 k$. The pair $(\omega_{\alpha}, A_{\alpha})$ converges uniformly at infinity to a pair $(\omega_0, A_{0,\alpha})$, where ω_0 is standard Euclidean metric and $A_{0,\alpha}$ is a k -instanton. Moreover, the will give solutions for which the Kähler metric is not cscK, making explicit the ‘bubbling’ argument in §5.3. The method we will follow to find solutions to (6.0.2) has two steps: first, we deform a solution for zero coupling constant γ via the application of the implicit function theorem in weighted Hölder spaces. For this it will be crucial the previous work of C. Arezzo and F. Pacard in [4] concerning the mapping properties of the scalar curvature in weighted spaces. Second, we apply a rescaling procedure via real homotheties. The uniform convergence at infinity will be provided by the fact that $A_{\alpha} - A_{0,\alpha}$ has coefficients in a weighted Hölder space. For simplicity in the first part we develop the arguments starting with the basic 1-instanton (whose definition we recall below) and explain how the argument generalizes for a suitable choice of k -instanton.

6.1. Preliminaries: Instantons on \mathbb{C}^2 .

We recall now some well known facts about instantons on $\mathbb{C}^2 \cong \mathbb{R}^4$ (see [5], [7] and [27]). Following [5] we identify \mathbb{R}^4 with the algebra of quaternions \mathbb{H} via the map

$$x = (x_1, x_2, x_3, x_4) \rightarrow x_1 + x_2 \cdot \mathbf{i} + x_3 \cdot \mathbf{j} + x_4 \cdot \mathbf{k} \in \mathbb{H}$$

where $\mathbf{i}, \mathbf{j}, \mathbf{k}$ satisfy the usual relations. The group $Sp(1)$ of quaternions with unit norm

$$x\bar{x} = |x|^2 = 1$$

gets identified with $SU(2)$ and its lie algebra can be viewed as the pure imaginary quaternions with basis i, j, k . The basic 1-instanton is defined as the $SU(2)$ connection

$$A = \text{Im} \frac{\bar{x}dx}{1+|x|^2} = \frac{1}{2} \cdot \frac{\bar{x}dx - d\bar{x}}{1+|x|^2}.$$

The curvature of A can be written as

$$F_A = \frac{d\bar{x} \wedge dx}{(1+|x|^2)^2},$$

where

$$\begin{aligned} d\bar{x} \wedge dx &= 2(dx_1 \wedge dx_2 - dx_3 \wedge dx_4) \cdot \mathbf{i} + 2(dx_1 \wedge dx_3 + dx_2 \wedge dx_4) \cdot \mathbf{j} \\ &\quad + 2(dx_1 \wedge dx_4 - dx_2 \wedge dx_3) \cdot \mathbf{k}. \end{aligned}$$

The components in \mathbf{i}, \mathbf{j} and \mathbf{k} form a basis of ASD 2-forms and with respect to the flat metric ω_0 so $F_A^+ = 0$.

Any other 1-instanton A' with respect to ω_0 is produced by a conformal transformation

$$f_{\beta,b} : x \rightarrow \beta(x - b),$$

where β is a positive real scalar and $b \in \mathbb{H}$. Multi-instantons or k -instantons, are given by pulling back the direct sum of the basic instanton A on \mathbb{H}^k via maps $u : \mathbb{H} \rightarrow \mathbb{H}^k$ of the form

$$(6.1.1) \quad u(x) = (x \text{Id} - B)\lambda,$$

where B is a symmetric $k \times k$ matrix of quaternions and λ is column vector of quaternions in \mathbb{H}^m satisfying some compatibility conditions (see [5, p. 25]). Hence, any k -instanton can be written as

$$A_u = \text{Im} \frac{u^*du}{1+|u|^2},$$

where the norm in \mathbb{H}^k is taken with respect to the direct sum of the euclidean norm.

6.2. Mapping properties of the connection-laplacian

Let A be the basic 1-instanton in \mathbb{C}^2 and consider its covariant derivative d_A acting on sections $a \in C^\infty(\mathbb{R}^4, \mathfrak{su}(2))$. In this section we study the mapping properties of the Hodge laplacian Δ_A as a map between weighted Hölder spaces. This is a self-adjoint elliptic operator of order 2 with respect to the Hermitian product provided by the trace tr and the Hermitian product in forms given by ω_0 . We shall discuss for which values of the weight Δ_A is an isomorphism. First we define weighted Hölder spaces of our

interest. The general theory about weighted spaces that we will use can be found in [47], [9], [59], [45] and [16].

Let $k \in \mathbb{N}$, $\beta \in (0, 1) \subseteq \mathbb{R}$, $\delta, r \in \mathbb{R}$, $r > 0$ be numbers and consider

$$B_r = \{z \in \mathbb{C}^2 : |z| < r\}$$

and \overline{B}_r its closure. For some fixed r_0 the space $C_\delta^{k,\beta}(\mathbb{C}^2 \setminus B_{r_0})$ is defined to be the space $C_{loc}^{k,\beta}(\mathbb{C}^2 \setminus B_{r_0})$ of locally Hölder functions ϕ for which the following norm

$$\|\phi\|_{C_\delta^{k,\beta}(\mathbb{C}^2 \setminus B_{r_0})} := \sup_{r \geq r_0} r^{-\delta} \|\phi(r \cdot)\|_{C^{k,\beta}(\overline{B}_2 \setminus B_1)}$$

is finite. From now on we fix $r_0 > 0$ and consider the following.

Definition 6.2.1. Given k, β and δ as before, the space $C_\delta^{k,\beta}$ is defined to be the space of functions $\phi \in C_{loc}^{k,\beta}(\mathbb{C}^2)$ for which the following norm is finite

$$\|\phi\|_{C_\delta^{k,\beta}} := \|\phi\|_{C^{k,\beta}(\overline{B}_{2r_0})} + \|\phi\|_{C_\delta^{k,\beta}(\mathbb{C}^2 \setminus B_{r_0})}.$$

In a similar way we define the spaces $D_\delta^{k,\beta}$ of sections of $\mathbb{C}^2 \times \mathfrak{su}(2)$.

Definition 6.2.2. Given k, β and δ as before, the space $D_\delta^{k,\beta}$ is defined to be the space of functions $a : \mathbb{C}^2 \rightarrow \mathfrak{su}(2)$, $a = a_1 \cdot \mathbf{i} + a_2 \cdot \mathbf{j} + a_3 \cdot \mathbf{k}$, such that $a_l \in C_\delta^{k,\beta}$ for $l = 1, 2, 3$.

Remark 6.2.3. The sign notation we use for the weight δ agrees with that in [4], [45] and [9] but has opposite sign with respect to the one in [47]. With our notation the weight reflects order of growth: functions in $C_\delta^{k,\beta}$ grow at most like $|x|^\delta$ when $|x| \rightarrow \infty$.

In the next proposition we summarize the properties of the spaces $C_\delta^{k,\beta}$ that we will need for our discussion. Analogue properties are satisfied by the spaces $D_\delta^{k,\beta}$.

Proposition 6.2.4 ([16],[59],[45]).

- (1) There is a continuous embedding $C_\delta^{k,\beta} \subseteq C_{\delta'}^{k',\beta'}$ if $\alpha > \alpha'$, $k > l$ and $\delta' > \delta$.
- (2) If $\phi \in C_\delta^{k,\beta}$ then $\nabla \phi \in C_{\delta-1}^{k-1,\beta}$.
- (3) The product is a continuous map from $C_\delta^{k,\beta} \times C_{\delta'}^{k,\beta}$ to $C_{\delta+\delta'}^{k,\beta}$.

Let Δ_A be the Hodge Laplacian of the basic 1-instanton given by

$$\Delta_A = d_A^* d_A = - *_e d_A *_e d_A,$$

where $*_e$ denotes the Hodge operator of the euclidean metric ω_0 acting on forms. To apply the theory of elliptic operators on weighted spaces it is convenient to express the previous Hodge Laplacian with respect to the radial coordinate $0 < r < \infty$ in \mathbb{R}^4 . Let $*_{S^3}$ and d_{S^3} be, respectively, the Hodge operator on forms and the De Rham differential in the Riemannian three sphere $S^3 \subseteq \mathbb{R}^4$.

Lemma 6.2.5. *There exists a connection B in the trivial $SU(2)$ -bundle over the sphere S^3 such that*

$$(6.2.1) \quad \Delta_A = -\partial_{rr}^2 - \frac{3}{r}\partial_r - \frac{1}{r^2} *_S (d_{S^3} + \frac{r^2}{1+r^2} B) *_S (d_{S^3} + \frac{r^2}{1+r^2} B).$$

PROOF. To compute formula (6.2.1) let us take polar coordinates in \mathbb{R}^4

$$\begin{aligned}x_1 &= r \cos \phi_1 = rw_1, \\x_2 &= r \cos \phi_2 \sin \phi_1 = rw_2, \\x_3 &= r \sin \phi_3 \sin \phi_2 \sin \phi_1 = rw_3, \\x_4 &= r \cos \phi_3 \sin \phi_2 \sin \phi_1 = rw_4,\end{aligned}$$

with $0 < r < \infty$ and $\sum_{l=1}^4 w_l^2 = 1$. Note that the basic instanton has the following expression in polar coordinates $x = rw$

$$\begin{aligned}A &= \frac{1}{2(1+r^2)}(\bar{x} \cdot dx - d\bar{x} \cdot x) \\&= \frac{1}{2(1+r^2)}(r\bar{w} \cdot d(rw) - d(r\bar{w}) \cdot rw) \\&= \frac{r^2}{2(1+r^2)}(\bar{w} \cdot dw - d\bar{w} \cdot w) = \frac{r^2}{1+r^2}B,\end{aligned}$$

where B is the $\mathfrak{su}(2)$ -valued 1-form

$$(6.2.2) \quad B = \frac{1}{2}(\bar{w} \cdot dw - d\bar{w} \cdot w).$$

Given $w \in S^3$ and any tangent vector $v_w \in T_w S^3$ we have

$$B_w(v_w) = \lim_{r \rightarrow \infty} A_{rw}(r \cdot v_w),$$

where $r : \mathbb{R}^4 \rightarrow \mathbb{R}^4$ is the homothety given by multiplication by r . Thus, B has a coordinate free expression and so it extends to a true connection on $S^3 \times SU(2)$. It can be readily checked from (6.2.2) that B is $SU(2)$ -invariant and flat, although we will not need this properties for our discussion. The covariant derivative of A can be now written as

$$d_A = \partial_r + d_{S^3} + \frac{r^2}{1+r^2}B.$$

Let $a \in C^\infty(\mathbb{R}^4, \mathfrak{su}(2))$. If we denote by $*_e$ the euclidean Hodge operator in \mathbb{R}^4 and by vol_{S^3} the standard volume form in S^3 we have

$$\begin{aligned}\Delta_A a &= d_A^* d_A = - *_e d_A *_e d_A a \\&= - *_e (\partial_r + d_{S^3} + \frac{r^2}{1+r^2}B) *_e (\partial_r a \cdot dr + d_{S^3} a + \frac{r^2}{1+r^2}[B, a]) \\&= - *_e (\partial_r + d_{S^3} + \frac{r^2}{1+r^2}B)(\partial_r a \cdot r^3 \text{vol}_{S^3} - rdr \wedge (d_{S^3} a + \frac{r^2}{1+r^2}[B, a])) \\&= - *_e (\partial_r (r^3 \partial_r a) dr \wedge \text{vol}_{S^3} + rdr \wedge (d_{S^3} a + \frac{r^2}{1+r^2}[B, a])) \\&= - \partial_{rr}^2 a - \frac{3}{r} \partial_r a - \frac{1}{r^2} *_S^3 (d_{S^3} a + \frac{r^2}{1+r^2}[B, a]) *_S^3 (d_{S^3} a + \frac{r^2}{1+r^2}[B, a]).\end{aligned}$$

□

Let $k \geq 2$ be a fixed positive integer. It follows from the properties of the weighted Hölder spaces summarized in Proposition 6.2.4, that, for all $\beta \in [0, 1)$ and $\delta \in \mathbb{R}$, the differential operator Δ_A induces a well defined bounded linear operator

$$\Delta_{A,\delta} : D_\delta^{k,\beta} \rightarrow D_{\delta-2}^{k-2,\beta}.$$

Following the notation in [47] we consider the operator

$$(6.2.3) \quad \Delta_\infty = -r^2 \partial_{rr}^2 - 3r \partial_r + \Delta_B,$$

where Δ_B denotes the Hodge Laplacian of the connection (6.2.2). The operator (6.2.3) induces well defined bounded linear operator $\Delta_{\infty,\delta} : D_\delta^{k,\beta} \rightarrow D_{\delta}^{k-2,\beta}$ for each value of the weight $\delta \in \mathbb{R}$. Making the change of variable $t = \ln r$ we obtain a translation invariant operator (see [47, p. 414])

$$(6.2.4) \quad \Delta_{\infty,t} = -\partial_{tt}^2 - 2\partial_t + \Delta_B$$

with associated operator

$$(6.2.5) \quad \Delta_{\infty,t}(\delta) = -\delta^2 - 2\delta + \Delta_B, \quad \delta \in \mathbb{R},$$

acting on sections $b \in C^\infty(S^3, \mathfrak{su}(2))$. Let $0 \leq \lambda_1 \leq \lambda_2 \leq \dots$ be the eigenvalues, counted with multiplicity, of the operator Δ_B , the Hodge Laplacian of the connection (6.2.2). Then, the homogeneous problem $\Delta_{\infty,t}(\delta)b = 0$ has nontrivial solutions if and only if

$$(6.2.6) \quad \delta_l^\pm = -1 \pm \sqrt{1 + \lambda_l},$$

where $\delta_l^\pm \in \mathbb{R}$ are the indicial roots of Δ_A (see [47, p. 416]). Then, since the λ_l are all greater or equal than 0 we obtain the following result.

Lemma 6.2.6. *There is no indicial root of Δ_A in $] -2, 0 [\subseteq \mathbb{R}$.*

It follows from [47, Theorem 6.1] that the operator $\Delta_{A,\delta}$ is Fredholm if and only if δ is not an indicial root of Δ_A . From [47, Lemma 7.1] it also follows that if $\delta_1, \delta_2 \in \mathbb{R}$ are weights such that there is no indicial root in $[\delta_1, \delta_2]$ then the dimension of Kernel and coKernel do not change. The heuristic reason for this fact is that the solutions to the homogeneous problem corresponding to Δ_∞ (equivalently to $\Delta_{\infty,t}$) do determine at ∞ all the possible asymptotic behaviors of the solutions of the homogeneous problem corresponding to Δ_A . Since Δ_A is self adjoint, using duality arguments in weighted Sobolev spaces (see the comments after the proof of Proposition 6.2.7) we also obtain that

$$(6.2.7) \quad \text{CoKer}\Delta_{A,\delta} \cong \text{Ker}\Delta_{A,-\delta-2},$$

provided that δ is not an indicial root. We now apply this facts to prove the main result of this section.

Proposition 6.2.7. *The operator $\Delta_{A,\delta}$ is injective if $\delta < 0$. It is surjective if δ is not an indicial root and $\delta > -2$.*

PROOF. Suppose first that there exists $a \in D_\delta^{k,\beta}$ such that $\Delta_A a = 0$. By elliptic regularity it follows that a is smooth. Let $B_r \subseteq \mathbb{R}^4$ be the open ball of radius r . If

$\delta < -1$ then we can integrate by parts obtaining

$$\begin{aligned} 0 &= \int_{\mathbb{R}^4} \text{tr } a \Delta_A a \text{ vol}_e = \lim_{r \rightarrow \infty} \int_{B_r} \text{tr } a \Delta_A a \text{ vol}_e \\ &= \lim_{r \rightarrow \infty} \int_{B_r} \text{tr } d_A a \wedge *_e d_A a + \lim_{r \rightarrow \infty} \int_{\partial B_r} \text{tr } a \wedge *_e d_A a \\ &= \lim_{r \rightarrow \infty} \int_{B_r} \text{tr } d_A a \wedge *_e d_A a = \int_{\mathbb{R}^4} \text{tr } d_A a \wedge *_e d_A a = \|d_A a\|_{\mathbb{R}^4}^2, \end{aligned}$$

where vol_e denotes the usual volume form in \mathbb{R}^4 . For the previous computation we have used that $|d_A a|_{\mathbb{R}^4}^2 \in D_{2\delta-2}^{k-2,\beta} \subseteq L^1(\mathbb{R}^4)$ due to $2\delta - 2 + 3 = 2\delta + 1 < -1$ and that the term $|\int_{\partial B_r} \text{tr } a \wedge *_e d_A a|$ grows at most like $r^{2\delta+2}$ as $r \rightarrow \infty$. Since the connection A is irreducible $d_A a = 0$ implies $a = 0$ thus $\Delta_{A,\delta}$ is injective if $\delta < -1$. Now, since there is no indicial root of Δ_A in $] -2, 0 [$, we conclude that $\Delta_{A,\delta}$ is injective if $\delta < 0$. To prove (2) we recall from (6.2.7) that $\Delta_{A,\delta}$ is surjective if and only if $\Delta_{A,-2-\delta}$ is injective, provided that δ is not an indicial root. This holds if $\delta > -2$ due to (1). \square

Note that in [47] the authors work in weighted Sobolev spaces, though it follows from standard inclusion relations between weighted Hölder and Sobolev spaces (see e.g. [45]) that their results hold also in weighted Hölder spaces. The rest of this section is devoted to clarify this point and to give a proof of (6.2.7). The reader familiarized with the theory of weighted spaces can skip this part going directly to §6.3. Let us first define weighted Sobolev spaces: we denote by $W_\delta^{0,p}$ be the space of locally integrable functions $\phi : \mathbb{R}^4 \rightarrow \mathbb{R}$ such that the following norm is finite

$$\|\phi\|_{p,0,\delta} = \left(\int_{\mathbb{R}^4} |< x >^{-\delta} \phi|^p < x >^{-4} \text{vol}_e \right)^{1/p}$$

where $< x > = (1 + |x|^2)^{1/2}$ and $\text{vol}_e = \frac{\omega_0^2}{2!}$ denotes the usual volume form in \mathbb{R}^4 . For a non-negative integer k the weighted Sobolev space $W_\delta^{k,p}$ is the set of locally integrable functions ϕ such that $|\nabla^l \phi| \in W_{\delta-l}^{0,p}$ for $0 \leq l \leq k$ with norm

$$\|\phi\|_{p,k,\delta} = \sum_{l=0}^k \|\nabla^l \phi\|_{p,0,\delta-l}.$$

In a similar way we define weighted sobolev spaces for sections in $\mathbb{C}^2 \times \mathfrak{su}(2)$ that we denote by $U_\delta^{k,p}$. Then, the spaces $W_\delta^{k,p}$ satisfy the following properties (see [45])

- (1) There is a continuous embedding $W_\delta^{k,p} \subseteq W_{\delta'}^{k',p}$ if $k > l$ and $\delta' > \delta$.
- (2) If $\phi \in W_\delta^{k,p}$ then $\nabla \phi \in W_{\delta-1}^{k-1,p}$.
- (3) There is a continuous embedding $C_{\delta-\epsilon}^{k,\beta} \subseteq W_\delta^{k,p}$ for every $p \geq 1$ and $\epsilon > 0$.
- (4) Suppose $p > 1$ and $l - k - \beta > 4/p$, then there is a continuous embedding $W_\delta^{l,p} \subseteq C_\delta^{k,\beta}$.
- (5) The dual space of $W_\delta^{k,p}$ is $W_{-4-\delta}^{k,p'}$, with $\frac{1}{p'} + \frac{1}{p} = 1$.

Analogue properties are satisfied by the spaces $U_\delta^{k,p}$. Property 2) implies that Δ_A induces a well defined operator

$$\Delta_\delta^U : U_\delta^{k,2} \rightarrow U_{\delta-2}^{k-2,2}$$

for each $\delta \in \mathbb{R}$. We will prove from the previous inclusion properties that there exists isomorphisms $\text{Ker}\Delta_\delta^U \cong \text{Ker}\Delta_{A,\delta}$ and $\text{coKer}\Delta_\delta^U \cong \text{coKer}\Delta_{A,\delta}$. This proves that our previous assertions concerning Hölder spaces follow from the results in [47]. Note first that the dimension of the Kernel and coKernel in weighted Sobolev spaces of a self adjoint elliptic operator do not depend on k . For the Kernel it follows from elliptic regularity and for the coKernel it follows from the fact that, since our operator Δ_A is self-adjoint, any Kernel is isomorphic to some coKernel with dual weight. Choosing any $\epsilon > 0$ and large $l >> 0$ from proposition 6.2.4 we have the inclusions

$$(6.2.8) \quad U_\delta^{l,2} \subseteq D_\delta^{k,\beta} \subseteq U_{\delta+\epsilon}^{k,2}$$

inducing others $\text{Ker}\Delta_\delta^U \subseteq \text{Ker}\Delta_{A,\delta}^{k,\alpha} \subseteq \text{Ker}\Delta_{\delta+\epsilon}^U$. From this we conclude, choosing ϵ small, $\text{Ker}\Delta_\delta^U \cong \text{Ker}\Delta_{A,\delta}$ since the dimension of the Kernel in weighted Sobolev spaces do not change if we do not cross an indicial root (see [47, Lemma 7.1]). On the other hand, using (6.2.8) and elliptic regularity we obtain

$$\text{Im}(\Delta_\delta^U : U_\delta^{l,2} \rightarrow U_{\delta-2}^{l-2,2}) = D_{\delta-2}^{k-2,\beta} \cap \text{Im}(\Delta_A : D_\delta^{k,\beta} \rightarrow D_{\delta-2}^{k-2,\beta})$$

By the second theorem of isomorphism $\text{coKer}\Delta_\delta^U \leq \text{coKer}\Delta_{A,\delta}$. A similar argument using the right hand-side inclusion in (6.2.8) shows $\text{coKer}\Delta_{A,\delta} \leq \text{coKer}\Delta_{\delta+\epsilon}^U$, from where we obtain $\text{coKer}\Delta_\delta^U \cong \text{coKer}\Delta_{A,\delta}$ choosing ϵ small enough.

Finally, we give a proof of (6.2.7) using the previous arguments. Since Δ_A is self-adjoint and the dual space $(U_{\delta-2}^{k-2,2})'$ is isomorphic to $U_{-\delta-2}^{k-2,2}$, if δ is not an indicial root we have an isomorphism $\text{CoKer}\Delta_\delta^U \cong \text{Ker}\Delta_{-\delta-2}^U$ (see [47, Proposition 4.1]). Hence, the operator $\Delta_{A,\delta}$ is surjective if and only if the operator $\Delta_{A,-2-\delta}$ is injective, provided that δ is not an indicial root.

6.3. Solutions for arbitrary coupling constant

Now we prove the existence of solutions to (6.0.2) for arbitrary coupling constant. First we deform the solution with $\alpha = 0$ provided by the basic 1-instanton A in \mathbb{C}^2 and the standard flat Kähler metric ω_0 . Let $E^c = \mathbb{C}^2 \times SL(2, \mathbb{C})$ be the $SL(2, \mathbb{C})$ principal bundle endowed with the holomorphic structure provided by A . We fix the standard Hermitian metric H_0 in \mathbb{C}^2 . Given a smooth map $a \in C^\infty(\mathbb{C}^2, \mathfrak{su}(2))$ we can consider the induced hermitian metric $H_0 \cdot e^{ia}$ and its Chern connection associated to the holomorphic structure in E^c . We denote by $F_{e^{ia}}$ the curvature of this connection. Note that

$$\frac{d}{dt}_{t=0} F_{e^{ita}} = \mathbf{i} \bar{\partial}_A \partial_A a.$$

Let $\mathcal{K} \subseteq C^\infty(\mathbb{C}^2, \mathbb{R})$ be the space of Kähler potentials of the euclidean metric ω_0 . Let

$$L : \mathbb{R} \times \mathcal{K} \times C^\infty(\mathbb{C}^2, \mathfrak{su}(2)) \rightarrow C^\infty(\mathbb{C}^2, \mathbb{R}) \times C^\infty(\mathbb{C}^2, \mathfrak{su}(2))$$

be the differential operator

$$(6.3.1) \quad L(\alpha, \phi, a) = (S_{\omega_\phi} + \alpha *_{\omega_\phi} \text{tr} F_{e^{ia}}^2 - \alpha 8\pi^2, \Lambda_{\omega_\phi} F_{e^{ia}})$$

where $\omega_\phi = \omega_0 + dd^c\phi$. Due to the fact that A is an Hermitian–Yang–Mills connection the derivative of L with respect to the last two variables at $(0, 0, 0)$ is

$$(6.3.2) \quad d_{\phi,a}L|_{(0,0,0)} = (\Delta_{\omega_0}^2\phi, \frac{1}{2}\Delta_{A_0}a + \langle F_{A_0}, \partial\bar{\partial}\phi \rangle)$$

where Δ_{ω_0} is the Hodge laplacian of the euclidean metric.

Lemma 6.3.1. *Let $\delta_1, \delta_2 \in \mathbb{R}$ with $0 < \delta_1 < 1$ and $-2 < \delta_2 < 0$. We fix $k \geq 4$ a non-negative integer and $0 < \beta < 1$, $\beta \in \mathbb{R}$. Then the differential operator L defined in (6.3.1) induces a well defined C^1 map*

$$(6.3.3) \quad L_{\delta_1, \delta_2} : V_{\delta_1, \delta_2} \rightarrow C_{\delta_1-4}^{k-4, \beta} \times D_{\delta_2-2}^{k-4, \beta}$$

where V_{δ_1, δ_2} is a neighbourhood of $(0, 0, 0)$ in $\mathbb{R} \times C_{\delta_1}^{k, \beta} \times D_{\delta_2}^{k-2, \beta}$. For such a pair (δ_1, δ_2) there exists $\epsilon = \epsilon(\delta_1, \delta_2) > 0$ and a C^1 curve

$$\Phi :]-\epsilon, \epsilon[\rightarrow C_{\delta_1}^\infty \times D_{\delta_2}^\infty$$

of solutions to the system $L(\alpha, \Phi(\alpha)) = 0$.

PROOF. Let $\delta_1, \delta_2 \in \mathbb{R}$ with $0 < \delta_1 < 1$ and $-2 < \delta_2 < 0$. We fix the basic 1-instanton A with curvature F_A . Note that the components of F_A are in C_{-4}^∞ and the components of the euclidean metric ω_0 are in C_0^∞ . We can take a neighbourhood

$$(0, 0, 0) \in V_{\delta_1, \delta_2} \subseteq \mathbb{R} \times C_{\delta_1}^{k, \beta} \times D_{\delta_2}^{k-2, \beta}$$

such that if $(0, \phi, a) \in V_{\delta_1, \delta_2}$ then $\det(\omega_\phi) \neq 0$. We denote by \propto equality modulo multiplicative constants. Let ϕ such that $(0, \phi, 0) \in V_{\delta_1, \delta_2}$. If s_ϕ denotes the scalar curvature of ω_ϕ , then $s_\phi \in C_{\delta_1-4}^{k-4, \beta}$ by [4, Proposition 4.1], provided that $0 < \delta_1 < 1$. On the other hand, if $a \in D_{\delta_2}^{k-2, \beta}$ then

$$h = e^{ia} = \sum_{l=0}^{\infty} \frac{(\mathbf{i}a)^l}{l!} \in D_0^{k-2, \beta},$$

by properties of the weighted Hölder spaces summarized in Proposition 6.2.4. Recall that $F_{e^{ia}} = F_A + \bar{\partial}_A h^{-1} \partial_A h$. then, since the components of A are in C_{-1}^∞ , we have

$$\begin{aligned} \partial_A h \in D_{\delta_2-1}^{k-3, \beta} &\Rightarrow \bar{\partial}_A h^{-1} \partial_A h \in D_{\delta_2-2}^{k-4, \beta} \\ &\Rightarrow F_{e^{ia}} \in D_{\max\{\delta_2-2, -4\}}^{k-4, \beta} = D_{\delta_2-2}^{k-4, \beta} \\ &\Rightarrow F_{e^{ia}}^2 \in D_{2\delta_2-4}^{k-4, \beta}. \end{aligned}$$

Finally, denoting by \propto equality modulo multiplicative constant, we have

$$\begin{aligned} *_{\omega_\phi} \text{tr } F_{e^{ia}}^2 &\propto (\det \omega_\phi)^{-1} \Lambda^2 F_{e^{ia}}^2 \in C_{2\delta_2-4}^{k-4, \beta} \subset C_{\delta_1-4}^{k-4, \beta}, \\ \Lambda_{\omega_\phi} F_{e^{ia}} &\propto (\det \omega_\phi)^{-1} \sum_{j=0}^1 \Lambda^j F_{e^{ia}} \wedge (dd^c\phi)^j \in C_{\delta_2-2}^{k-2, \beta}, \end{aligned}$$

where Λ denotes the adjoint of the Lefschetz operator $\alpha \rightarrow \alpha \wedge \omega_0$ acting on forms and multiplicative constants on each summand has been omitted. For the last computation

we have used that (omitting multiplicative constant factors)

$$\det \omega_\phi \propto 1 + \sum_{l=1}^2 \Lambda^l (dd^c \phi)^l \in C_0^{k-2,\beta}.$$

We thus conclude that the differential operator L defined in (6.3.1) induces well defined C^1 operators $L_{\delta_1,\delta_2} : V_{\delta_1,\delta_2} \rightarrow C_{\delta_1-4}^{k-4,\beta} \times D_{\delta_2-2}^{k-4,\beta}$ for the chosen values of the weights. From [4, Proposition 5.5] we know that the double Laplacian

$$\Delta_{\omega_0}^2 : C_{\delta_1}^{k,\beta} \rightarrow C_{\delta_1-4}^{k-4,\beta}$$

is surjective and has Kernel equal to the constant functions if $0 < \delta_1 < 1$. We also know that the connection-laplacian Δ_{A,δ_2} is an isomorphism for $-2 < \delta_2 < 0$. From (6.3.2) we conclude that $d_{\phi,a} L_{\delta_1,\delta_2}|_{(0,0,0)}$ is an isomorphism modulo constant functions in the weighted space $C_{\delta_1}^{k,\beta}$. Applying the implicit function theorem we obtain a curve of solutions $\Phi :]-\epsilon, \epsilon[\rightarrow V_{\delta_1,\delta_2}$ to $L(\alpha, \Phi(\alpha)) = 0$ for small $\epsilon > 0$. The regularity of the solutions follows from local uniqueness choosing arbitrary large values of the integer k . \square

Remark 6.3.2. A crucial fact for the existence result is that the operator L defined in (6.3.1) induces well defined operators L_{δ_1,δ_2} for that values of the weights for which its linearization (see equation (6.3.2)) induces an isomorphism in weighted Hölder spaces. This is not obvious a priori and strongly depends on the shape of the equations (6.0.2).

The curve Φ in the statement of Lemma 6.3.1 provides a curve of pairs $(\omega_\alpha, A'_\alpha)$ where $\omega_\alpha = \omega_0 + dd^c \phi_\alpha$ is a Kähler metric in \mathbb{C}^2 and A'_α is the Chern connection of the Hermitian metric $e^{ia_\alpha} H_0$. This curve of pairs is a curve of solutions to (6.0.2) when we consider the system for $SL(2, \mathbb{C})$ connections. To obtain $SU(2)$ -connections solving (6.0.2) we must do a complex gauge transformation $h_\alpha = e^{ia_\alpha}$. The new curve of $SU(2)$ -connections is

$$(6.3.4) \quad A_\alpha = h_\alpha A^{1,0} h_\alpha^{-1} + h_\alpha \partial h_\alpha^{-1} + h_\alpha^{-1} A^{1,0} h_\alpha + h_\alpha^{-1} \bar{\partial} h_\alpha$$

that tends uniformly to A at infinity since $a_\alpha \in C_{\delta_2}^\infty$ with weight $-2 < \delta_2 < 0$. Note also that the 2-form $\omega_\alpha - \omega_0 = dd^c \phi_\alpha$ has coefficients in $C_{\delta_1-2}^\infty$ with $0 < \delta_1 < 1$ so ω_α is asymptotically euclidean at infinity. Finally, a similar discussion as in the proof of Proposition 6.2.7 shows that we can write

$$\int_{B_r} \text{tr } F_A \wedge F_A - \int_{B_r} \text{tr } F_{A_\gamma} \wedge F_{A_\gamma} = \int_{B_r} d\sigma = \int_{\partial B_r} \sigma$$

where $B_r \subseteq \mathbb{R}^4$ is the open ball of radius r and σ is a three form such that

$$\lim_{r \rightarrow \infty} \int_{\partial B_r} \sigma = 0,$$

due to the fact that its coefficients are in suitable C_δ^∞ spaces. This proves that the topological constraint (6.0.3) is satisfied for every α . An important remark is that the point-wise squared norm $|F_A|^2$ of the curvature of the basic instanton is non-constant with respect to the euclidean metric. In fact,

$$|F_A|^2 = \frac{24}{(1 + |x|^2)^4},$$

which is a bell-shaped function on \mathbb{C}^2 centered at the origin. Similar arguments as in §5.3, show that when we apply Lemma 6.3.1 to this initial data it produces solutions of the coupled system (6.0.4) in which the Kähler metric has non constant scalar curvature.

A similar discussion can be done for k -instantons, with arbitrary $k \in \mathbb{Z}$ choosing suitable functions u in (6.1.1). For example, one can take

$$u(x) = (x \operatorname{Id} - B)\lambda$$

with B a diagonal matrix with different quaternionic entries and λ equals $(\lambda_1, \dots, \lambda_k)$ with λ_j real positive scalars. The resulting k -instanton therefore looks like a superposition of k instantons with scales λ_j and centres b_j (see [5]). This correspond to the classical solutions discovered by t' Hooft. For this particular choice of u 's the coefficients of the corresponding instanton A_u are again in C_{-1}^∞ , the coefficients of the curvature F_{A_u} are in C_{-4}^∞ , the connection-Laplacian admits an expression similar to (6.2.1) and, therefore, similar results as that of Lemma 6.2.7 and Lemma 6.3.3 can be proved. We sum up this discussion in the following Lemma.

Lemma 6.3.3. *Let $k \in \mathbb{Z}$, $\delta_1 \in]0, 1[\subseteq \mathbb{R}$ and $\delta_2 \in]-2, 0[\subseteq \mathbb{R}$. For each triple (k, δ_1, δ_2) and suitable choice of a k -instanton A there exists a small interval $I \subseteq \mathbb{R}$ containing the origin such that for each $\alpha \in I$ there exists a solution $(\omega_\alpha, A_\alpha)$ of the coupled equations with coupling constant α and fixed topological invariant (6.0.3). The metric ω_α is an asymptotically euclidean Kähler metric with $\omega_\alpha - \omega_0 \in C_{\delta_1-2}^\infty$ and A_α converges uniformly to A with $A_\alpha - A \in D_{\delta_2-1}^\infty$.*

Remark 6.3.4. Note that the interval I in the previous proposition depends a priori on the concrete k -instanton we chose to apply the implicit function theorem (thus on the concrete function u in formula (6.1.1)). Note also that the connection A_α has bounded L^2 norm in \mathbb{R}^4 with respect to the Hermitian metric provided by ω_α due to the formula

$$\|F_{A_\alpha}\|_{\omega_\alpha}^2 = \|F_{A_\alpha}^-\|_{\omega_\alpha}^2 = \int_{\mathbb{R}^4} \operatorname{tr} F_{A_\alpha} \wedge F_{A_\alpha} = 8\pi^2 k.$$

We finish with the main result of this section that holds by the previous Lemma together with a suitable re-scaling of the solutions to (6.0.2) obtained for small coupling constant α .

Theorem 6.3.5. *Let $k \in \mathbb{Z}$, $\delta_1 \in]0, 1[\subseteq \mathbb{R}$ and $\delta_2 \in]-2, 0[\subseteq \mathbb{R}$. For each triple (k, δ_1, δ_2) and each $\alpha \in \mathbb{R}$ there exists a solution $(\omega_\alpha, A_\alpha)$ of the coupled equations with coupling constant α and fixed topological invariant (6.0.3). The metric ω_α is an asymptotically euclidean Kähler metric with $\omega_\alpha - \omega_0 \in C_{\delta_1-2}^\infty$. For each α there exists a k -instanton A'_α , such that $A_\alpha - A'_\alpha \in D_{\delta_2-1}^\infty$.*

PROOF. Let (ω, A) be a solution to (6.0.2) with non-vanishing coupling constant α and topological invariant (6.0.3) for some $k \in \mathbb{Z}$. We claim that for each $\beta \in \mathbb{R}$ such that $\alpha \cdot \beta > 0$ there exists a solution (ω, A_β) with coupling constant β and the same topology. Let $c = \beta/\alpha > 0$ and consider the Kähler form $c\omega$. Then the pair $(c\omega, A)$

satisfies

$$\begin{aligned}
 S_{c\omega} + \beta *_{c\omega} \text{tr } F_A \wedge F_A &= \frac{1}{c} S_\omega + \frac{\beta}{c^2} *_\omega \text{tr } F_A \wedge F_A \\
 &= \frac{1}{c} (S_\omega + \alpha *_\omega \text{tr } F_A \wedge F_A) \\
 &= \beta 8\pi^2 k
 \end{aligned}$$

so it is a solution with coupling constant β , due to the fact that the ASD equation (6.0.4) is conformally invariant. Consider now the homothety $h_{1/\sqrt{c}}(z) = z/\sqrt{c}$, that satisfies $h_{1/\sqrt{c}}^* c\omega = \omega$. Then it clear that (ω, A_β) is a solution with coupling constant β . The connection A_β has fixed topological invariant $8\pi^2 k$ in (6.0.3) since the instanton number k is preserved by conformal transformations (see [5] and [27]). It is clear from the proof that if the initial connection converges at infinity to some k -instanton A' , then the re-scaled connection converges to $h_{1/\sqrt{c}}^* A'$. \square

CHAPTER 7

Conclusions and future directions

In this thesis we have defined coupled equations for Kähler metrics and Hermite–Yang–Mills connections with symplectic and variational interpretations. We have seen that these equations are well suited for the moduli problem with Kähler structure for metrics and connections. We have constructed families of examples of solutions and we have related the coupled equations with the cscK problem on ruled manifolds. We have defined a new notion of stability for triples (X, L, E) , consisting of a projective polarised manifold (X, L) and a holomorphic vector bundle E . A conjecture relating this stability notion with the existence of solutions of the coupled system has been stated and supporting evidence has been given.

This thesis starts a programme, the one of studying coupled equations for pairs consisting of a Kähler metric and a connection, and sets many open questions that will be considered in future work. Possible future directions are: 1) construct moduli spaces of metrics and connections with Kähler structure, 2) proving a Hitchin–Kobayashi correspondence for the coupled equations, starting with the case of toric manifolds and toric vector bundles, 3) development of gluing methods for the coupled equations adapting those in [4] for asymptotically Euclidean Kähler metrics and the well known methods for instantons of H.C. Taubes and S.K. Donaldson (see [27]), 4) study the relation of the coupled equations with the cscK equation for a Kähler metric on the total space of an associated bundle.

As an immediate future project we will try to prove that the existence of solutions to (0.0.3) implies the α -K-semistability of the triple adapting Donaldson’s methods in [25]. This fact was first verified for the cscK problem, after the work of S. Zhang, H. Luo and S.K. Donaldson.

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